

REPORT ON THE  
SPECTRAL IP TEST SURVEY  
OVER THE  
GRUM DEPOSIT LINE 108W  
FARO, YUKON TERRITORY  
FOR  
CYRUS ANVIL MINES LTD.

VOL. I

-001967



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PHOENIX GEOPHYSICS LIMITED

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OVER THE  
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FARO, YUKON TERRITORY  
FOR  
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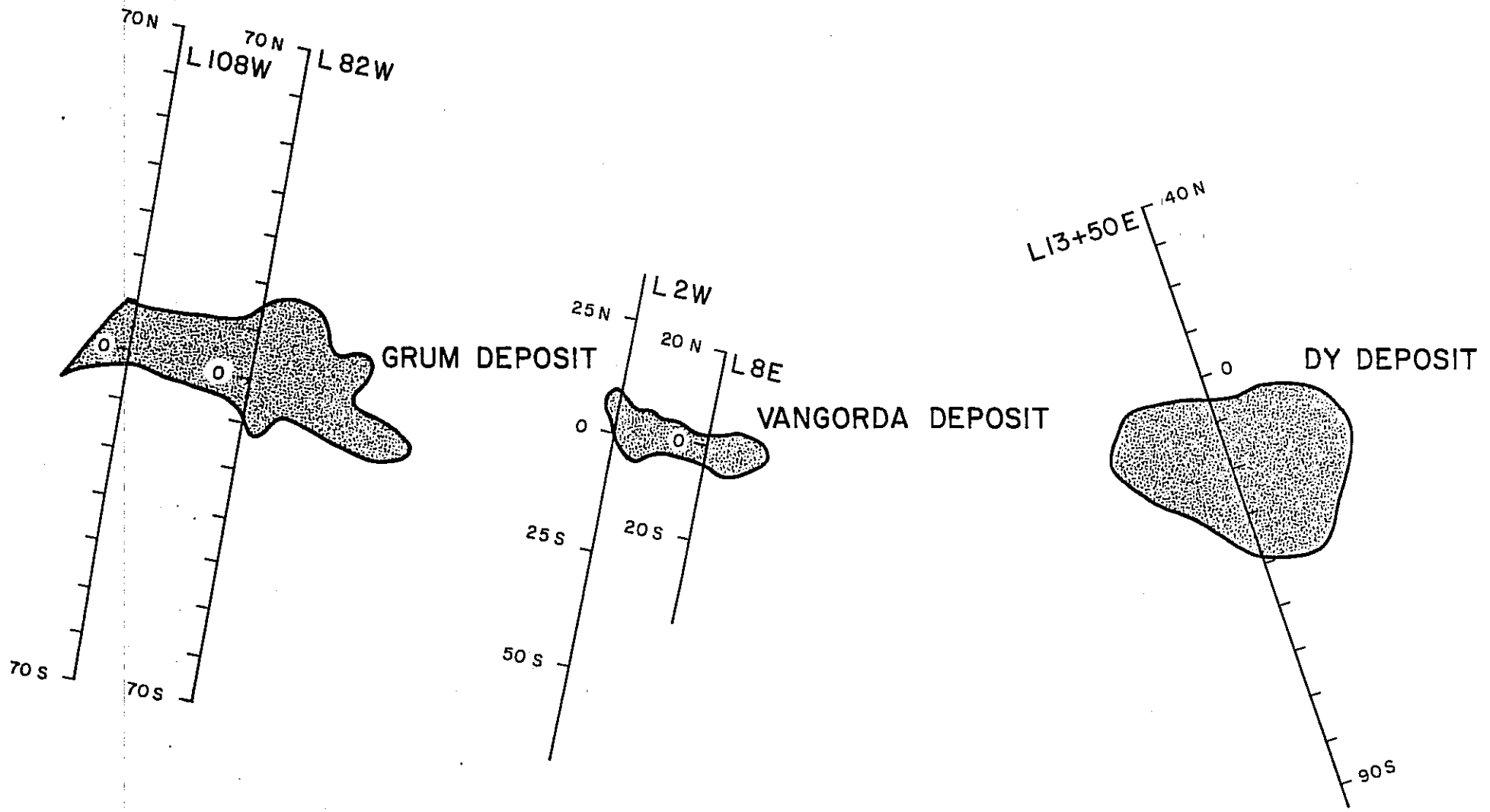
1. INTRODUCTION

At the request of Dr. David Jennings, Chief Geologist for Cyprus-Anvil Mines Ltd., we have completed a brief Spectral IP Test Survey over the Grum Deposit. The Grum Deposit is a known orebody on the company's property at Faro, Yukon Territory. Spectral IP measurements were made on Line 108W using  $X = 500'$ ;  $n = 1,2,3,4,5,6$  and  $X = 1000'$ ;  $n = 1,2,3,4,5,6$ .

As shown on the enclosed plan map, Line 108W crosses over the west end of the Grum Deposit. At this location, the depth to the top of the sulphide mineralization is almost three hundred meters. It is necessary to use large electrode intervals and electrode separations to detect a source at this depth. Inductive coupling effects can distort the measurements, but the use of the Spectral IP Method permits coupling effects to be removed.

2. PRESENTATION OF RESULTS

The apparent resistivity and the apparent phase IP results at 1.0 Hz



TEST LINE SURVEY POSITION  
VANGORDA

are enclosed with this report. The field results have been plotted using the pseudosection format. A few of the complete spectral plots, with computer inversions, are also enclosed with this report, in order to clarify the interpretation.

All of the pseudosection plots that output from the computer interpretation of the Spectral IP data, as well as all of the spectral plots are included in Volume II of this report.

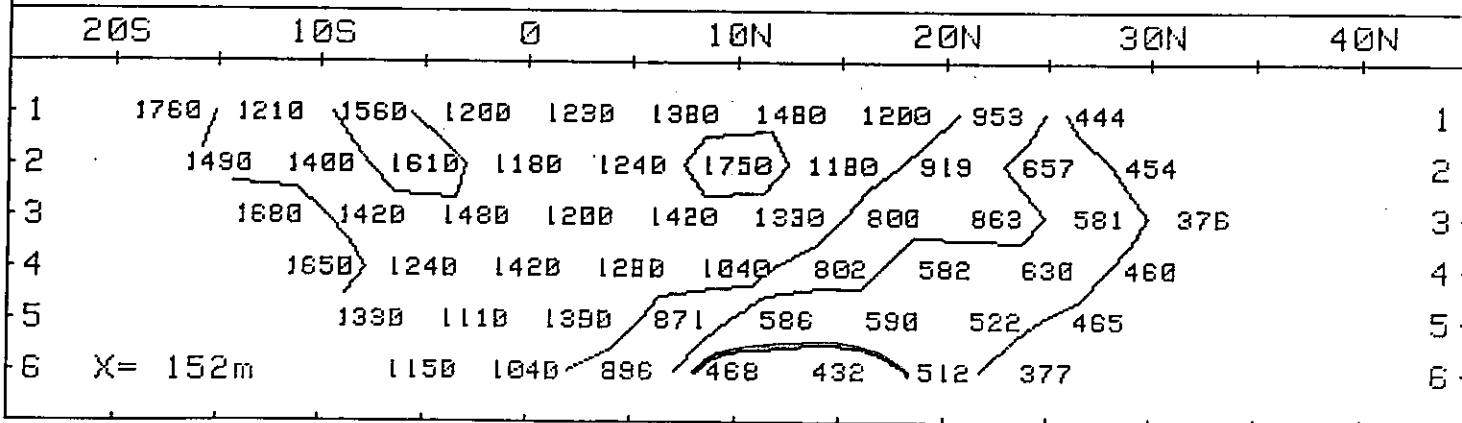
### 3. DISCUSSION OF RESULTS

For the measurements with both  $X = 500'$  and with  $X = 1000'$ , the most obvious feature on the apparent resistivity pseudosection is the definite contact at about 25+00N. There is a high resistivity rock unit to the south and a lower resistivity rock unit to the north. The  $X = 1000'$  measurements show that the southern edge of the high resistivity (non-porous and dense) rock unit is at 50+00S.

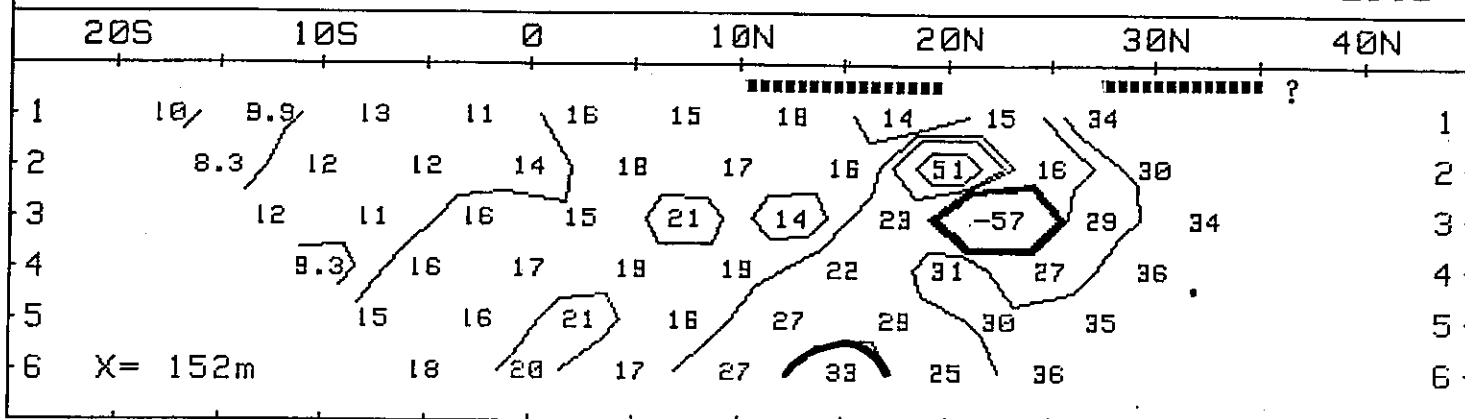
The  $X = 500'$  measurements suggest that the lower resistivity rocks are present, at depth beneath the high resistivity rock unit, to the south of 25+00N. The pseudosection for the  $X = 1000'$  measurements confirm that the lower resistivity rock unit is present, at depth, at least as far south as 0+00 to 10+00S. It is possible that the lower resistivity rocks extend completely beneath the 7500 meter surface extent of the high resistivity rock unit.

The  $X = 500'$  measurements show that there is an increased background IP effect in the lower resistivity rocks to the north of 25+00N. In addition, there are slightly anomalous apparent IP effects for  $n = 5$  and  $n = 6$  in the region from 5+00N to 25+00N; these would correlate with the upper edge of the Grum Deposit.

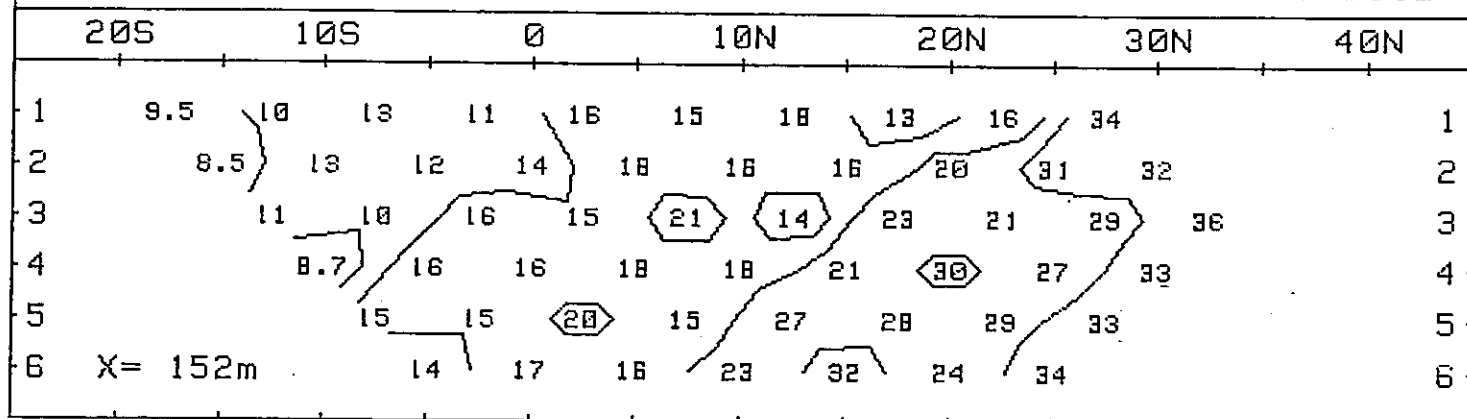
PHOENIX SPECTRAL IP OBS. RESISTIVITY (ohm-m) 1 Hz  
 CYPRUS ANVIL GRUM DEP. FARO YT 108W 152m 12/7/81 L312



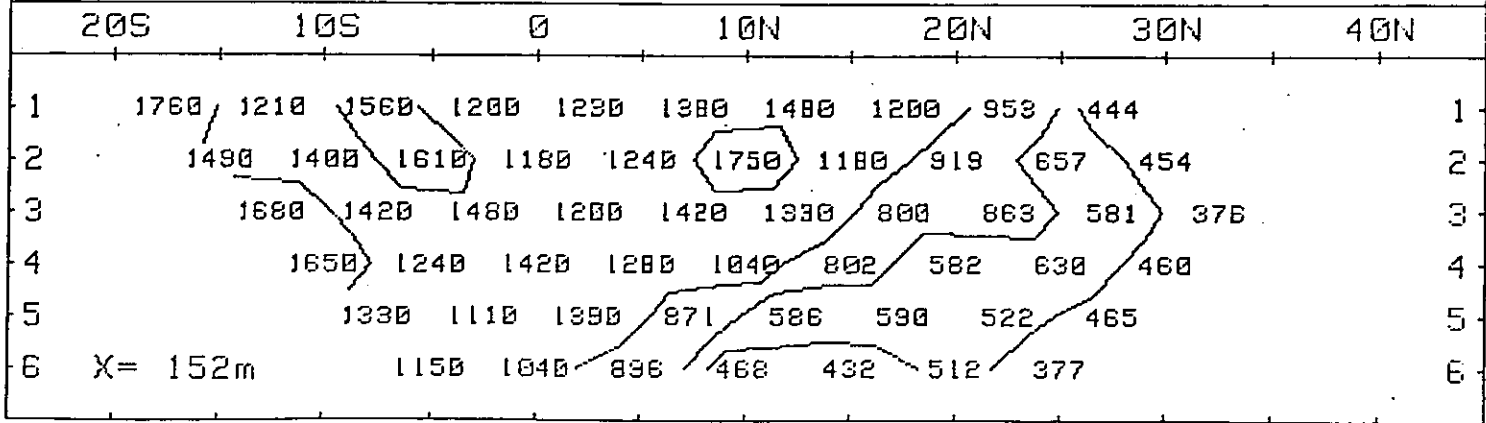
PHOENIX SPECTRAL IP PHASE (mrad) at 1 Hz  
 CYPRUS ANVIL GRUM DEP. FARO YT 108W 152m 12/7/81 L312



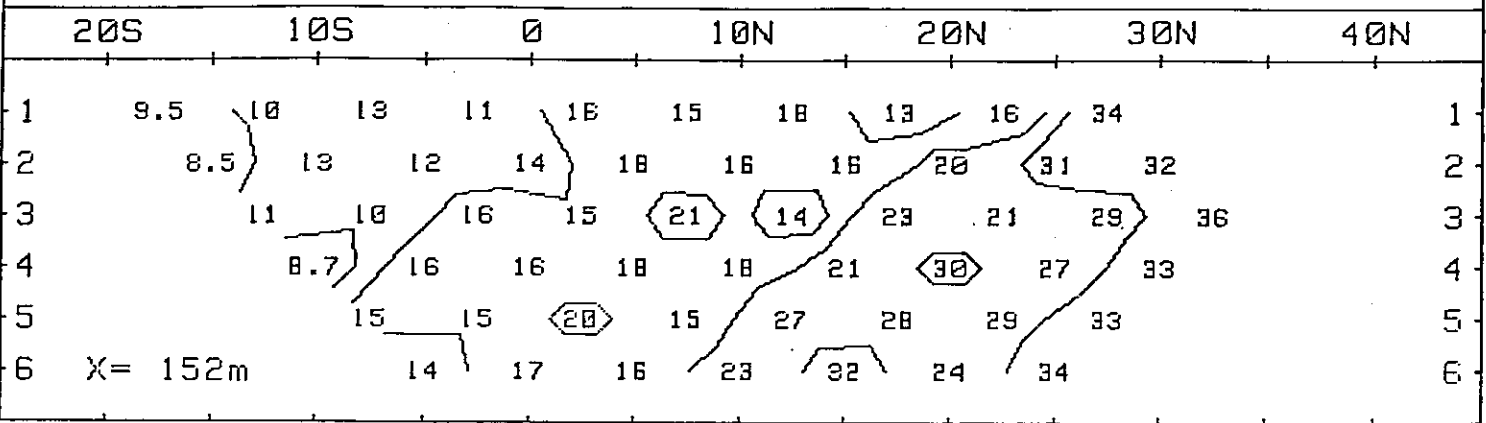
PHOENIX SPECTRAL IP DECOUPLED PHASE (mrad) 1 Hz  
 CYPRUS ANVIL GRUM DEP. FARO YT 108W 152m 12/7/81 L312



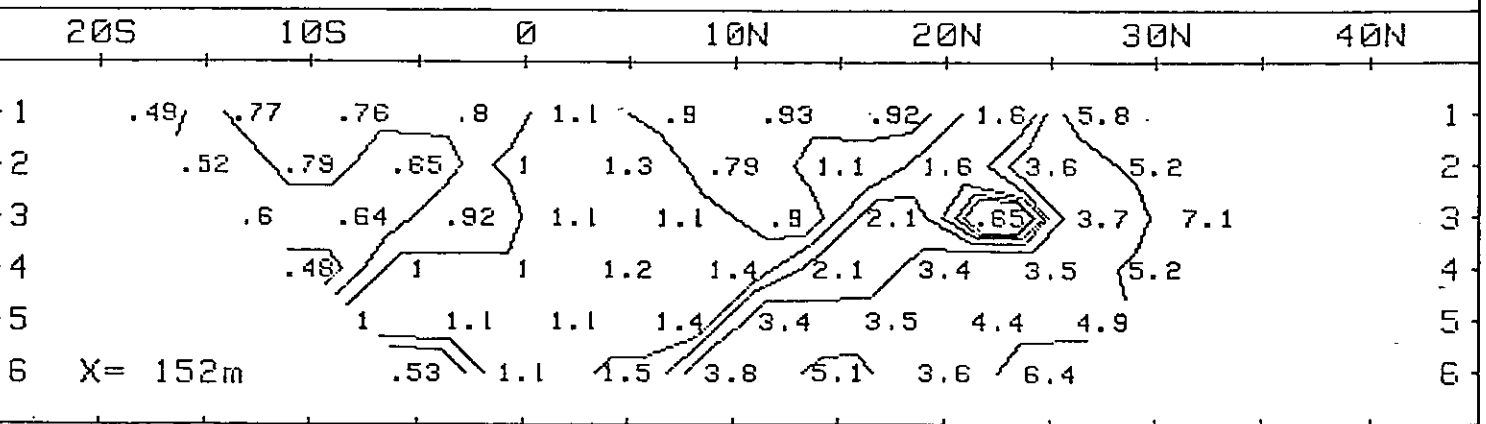
PHOENIX SPECTRAL IP OBS. RESISTIVITY (ohm-m) 1 Hz  
 CYPRUS ANVIL GRUM DEP. FARO YT 108W 152m 12/7/81 L312

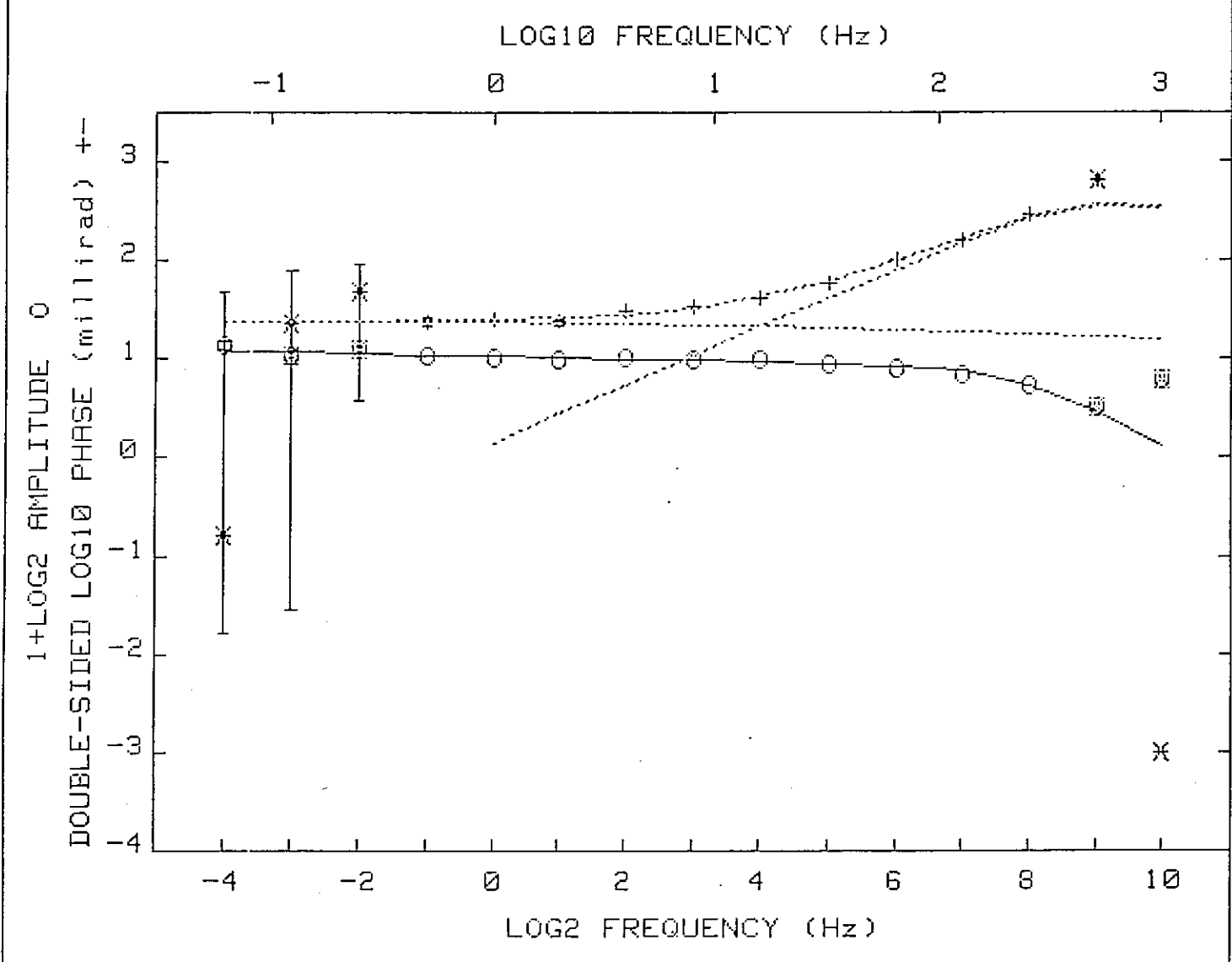


PHOENIX SPECTRAL IP DECOUPLED PHASE (mrad) 1 Hz  
 CYPRUS ANVIL GRUM DEP. FARO YT 108W 152m 12/7/81 L312



PHOENIX SPECTRAL IP DECOUPLED METAL FACTOR 1Hz  
 CYPRUS ANVIL GRUM DEP. FARO YT 108W 152m 12/7/81 L312





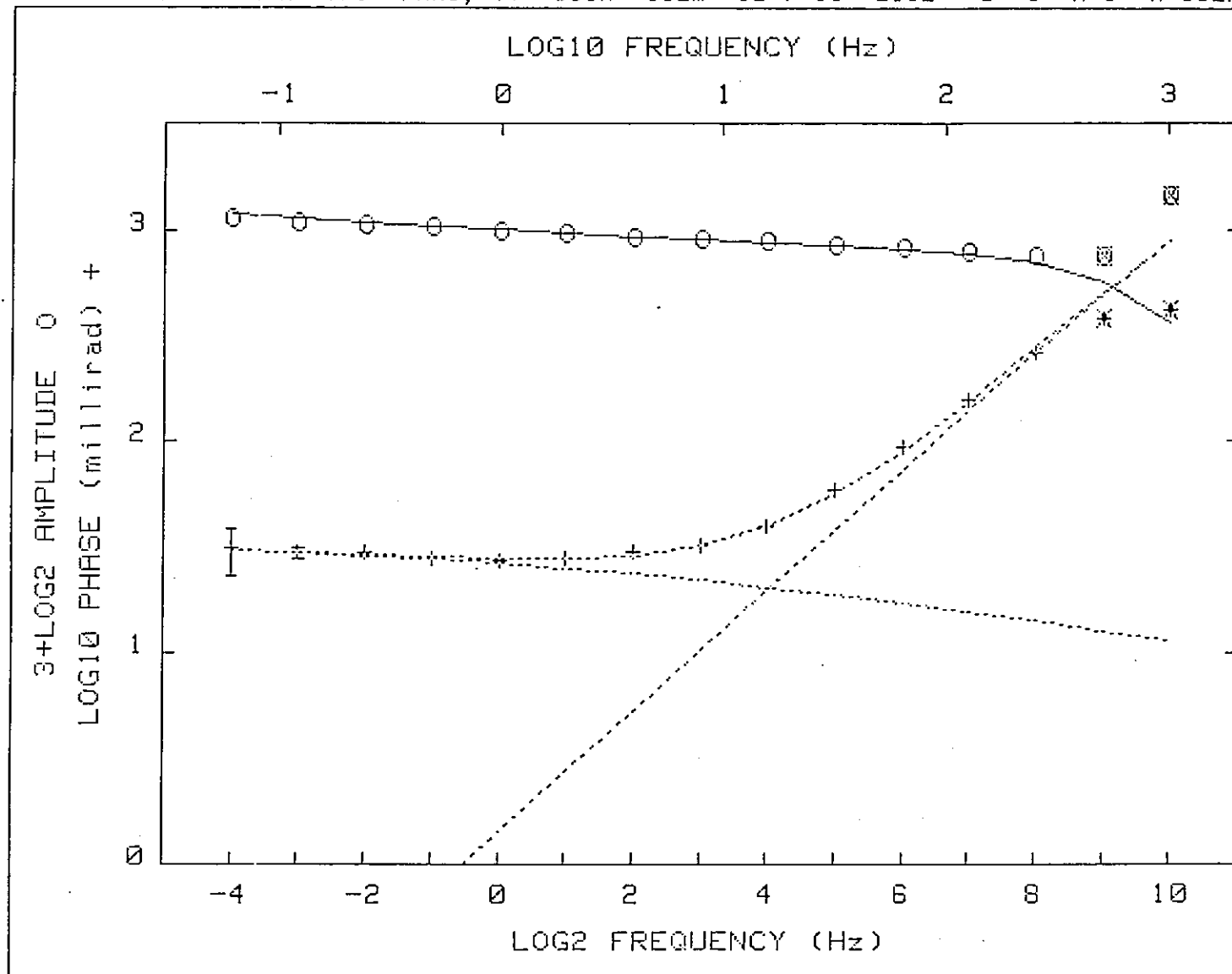
DRL: Number of dispersions= 2  
 C1=.15 M2=.39 T2=.00038 C2=.97 fixed

Iter	Lambda	Rchsq	R0	M1	T1
0	1.E-02	.00064	1.293	.336	3.9E+00
1	1.E-02	.00064	1.294	.336	3.9E+00

Pct Std Deviations            4.6    4.4    913.3

Correlation Matrix            1.000  
                                   .933    1.000  
                                   .984    .885    1.000

----- Apparent Resistivity Measured at 1-Hz is 511.7 -----  
 Apparent Resistivity Calculated from Inductive Coupling is 1709



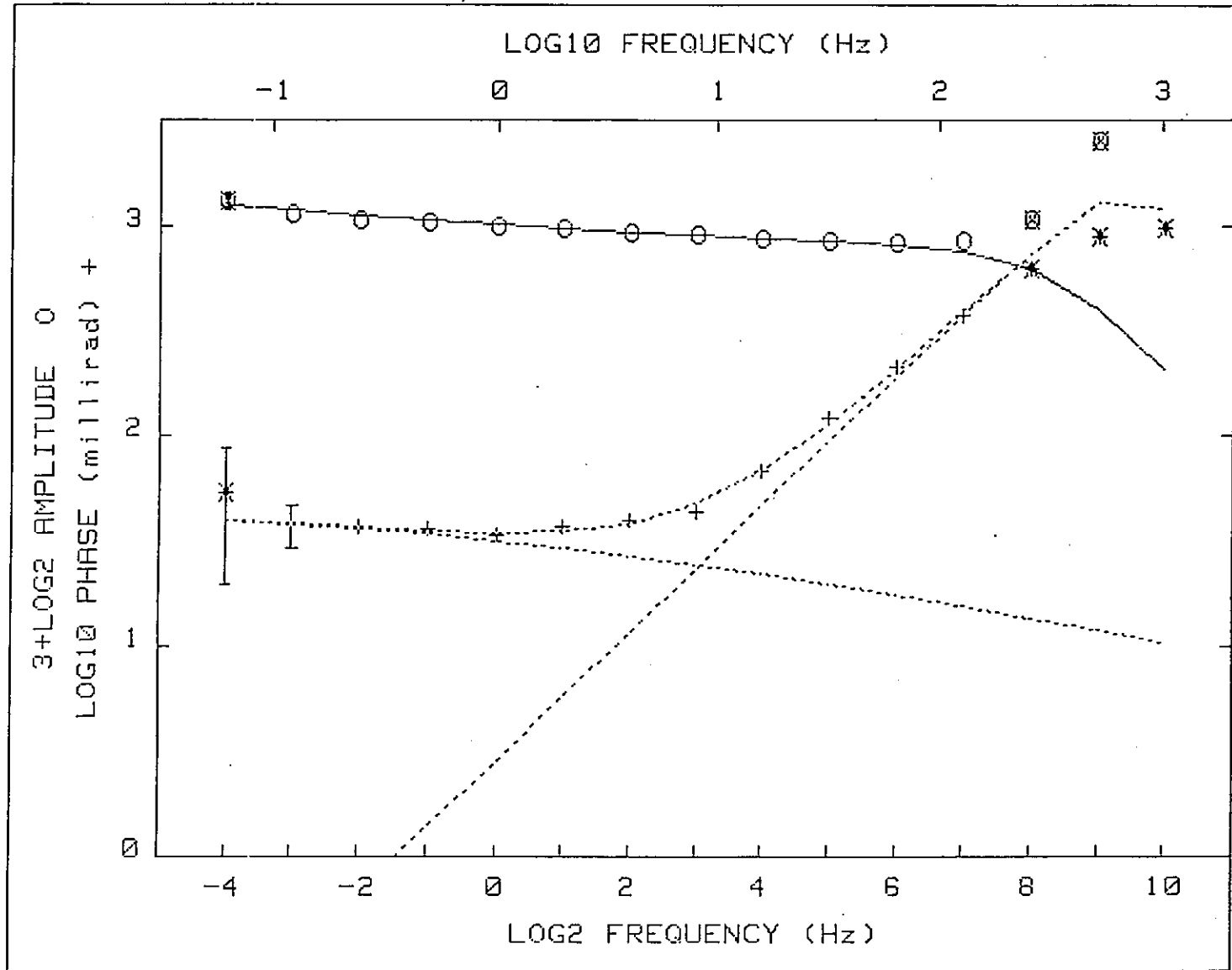
CRL: Number of dispersions= 2  
M2=1 T2=.00011 C2=.94 fixed

Iter	Lambda	Rchsq	R0	M1	T1	C1
0	1.E-02	.00018	1.311	.324	2.0E+01	.198
1	1.E-02	.00018	1.312	.324	2.0E+01	.198

Pct Std Deviations            1.8     4.1     70.7     4.3

Correlation Matrix		1.000			
		.943	1.000		
		.962	.880	1.000	
		-.637	-.833	-.535	1.000

Apparent Resistivity Measured at 1 Hz is 586.2  
Apparent Resistivity Calculated from Inductive Coupling is 1498



CRL: Number of dispersions= 2  
 M2=1 T2=.00031 C2=1 fixed

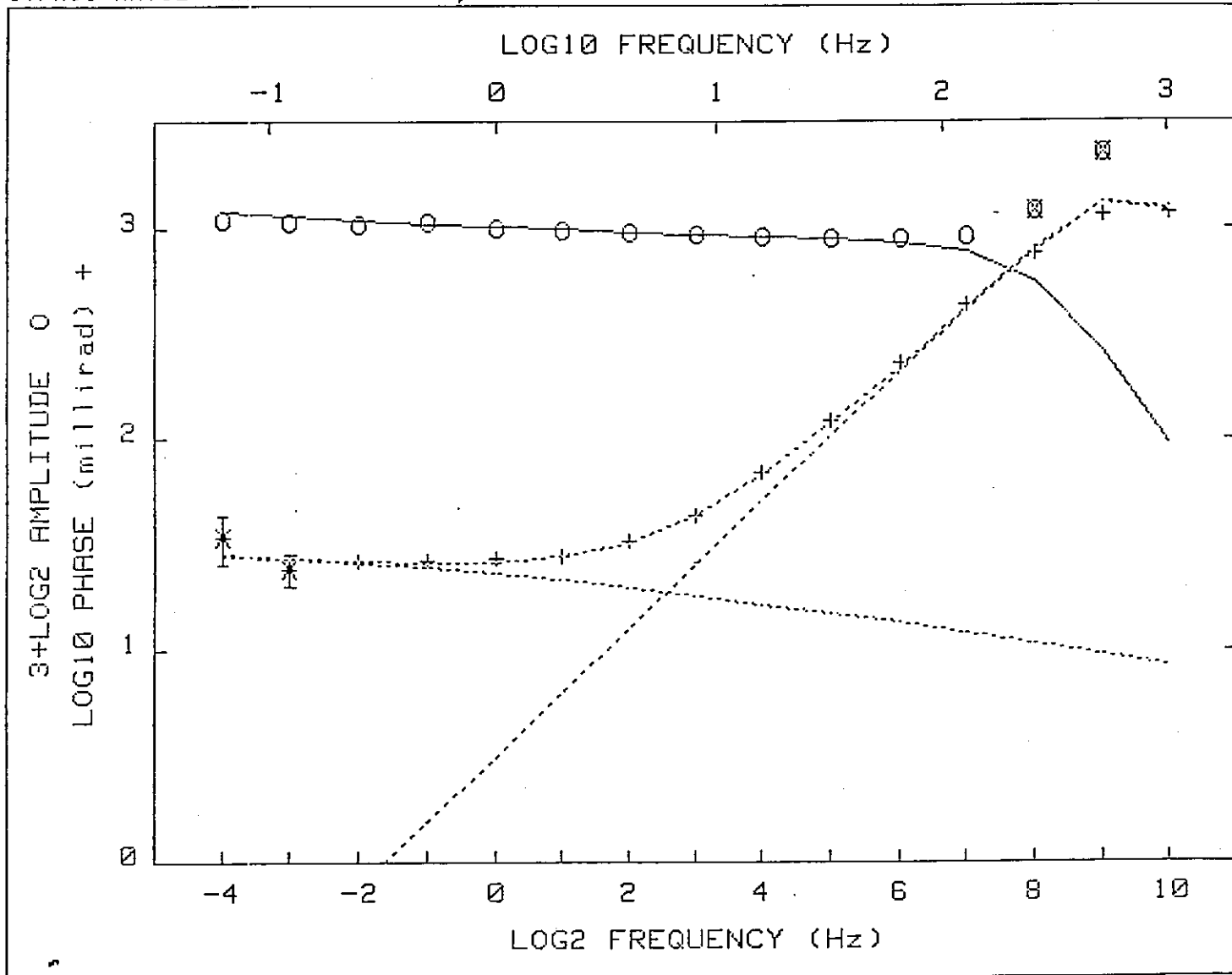
Iter	Lambda	Rchsq	R0	M1	T1	C1
0	1.E-02	.00043	1.431	.377	5.4E+01	.223
1	1.E-02	.00042	1.433	.378	5.6E+01	.223

Pct Std Deviations            3.1    5.0    96.0    6.7

Correlation Matrix

	1.000			
	.923	1.000		
	.972	.899	1.000	
	-.111	-.451	-.118	1.000

Apparent Resistivity Measured at 1 Hz is 492.0  
 Apparent Resistivity Calculated from Inductive Coupling is 879.1



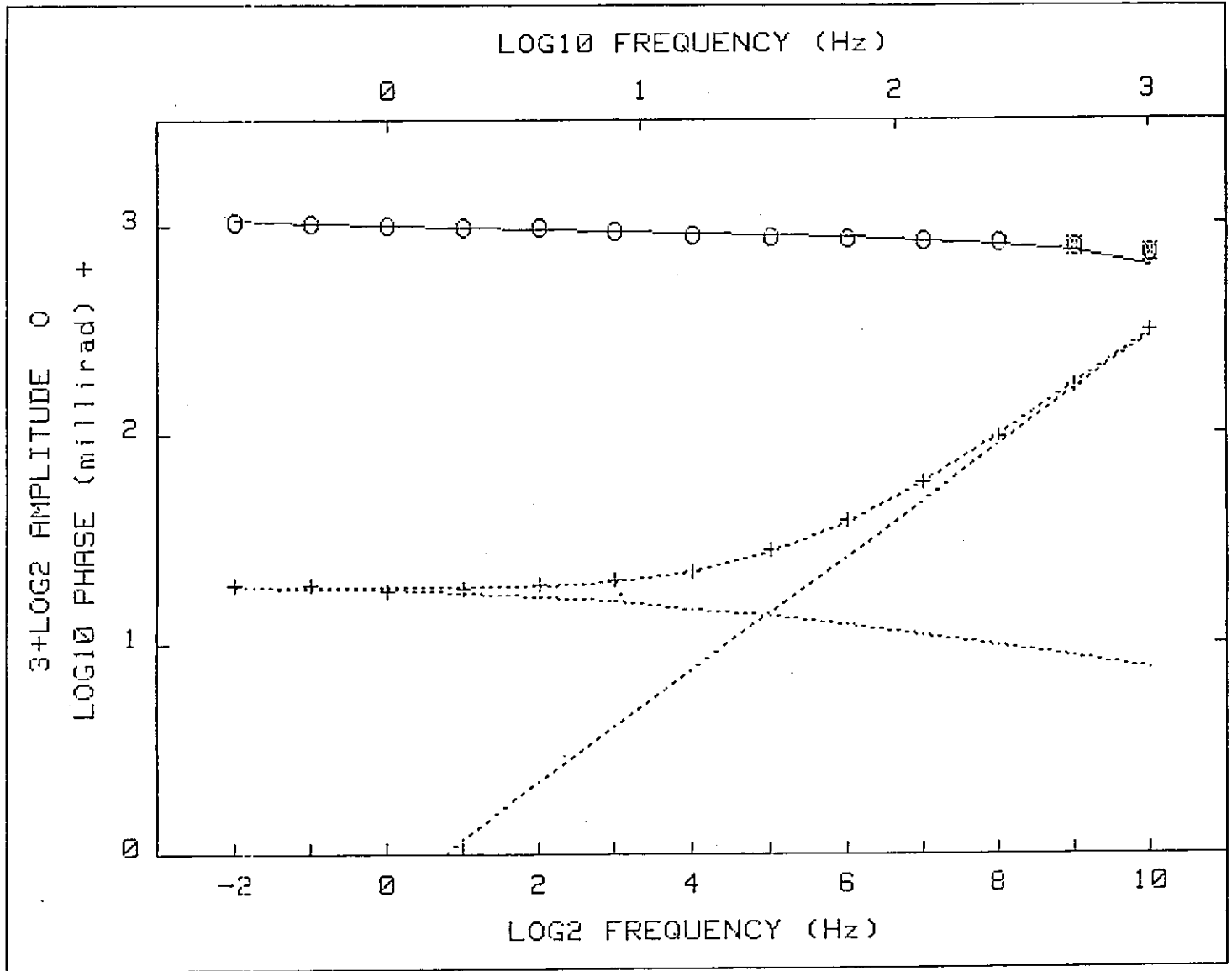
CRL: Number of dispersions= 2  
M2=1 T2=.00038 C2=1 fixed

Iter	Lambda	Rchsq	R0	M1	T1	C1
0	1.E-02	.00080	1.238	.266	8.4E+00	.219
1	1.E-02	.00076	1.251	.280	9.3E+00	.203
2	1.E-03	.00071	1.271	.288	1.7E+01	.201
3	1.E-04	.00069	1.296	.294	3.9E+01	.210
4	1.E-05	.00068	1.300	.296	4.0E+01	.212

Pct Std Deviations            2.6    5.9    126.0    13.1

Correlation Matrix		1.000			
		.748	1.000		
		.944	.637	1.000	
		.134	-.516	.233	1.000

Apparent Resistivity Measured at 1 Hz is 468.1  
Apparent Resistivity Calculated from Inductive Coupling is 719.4



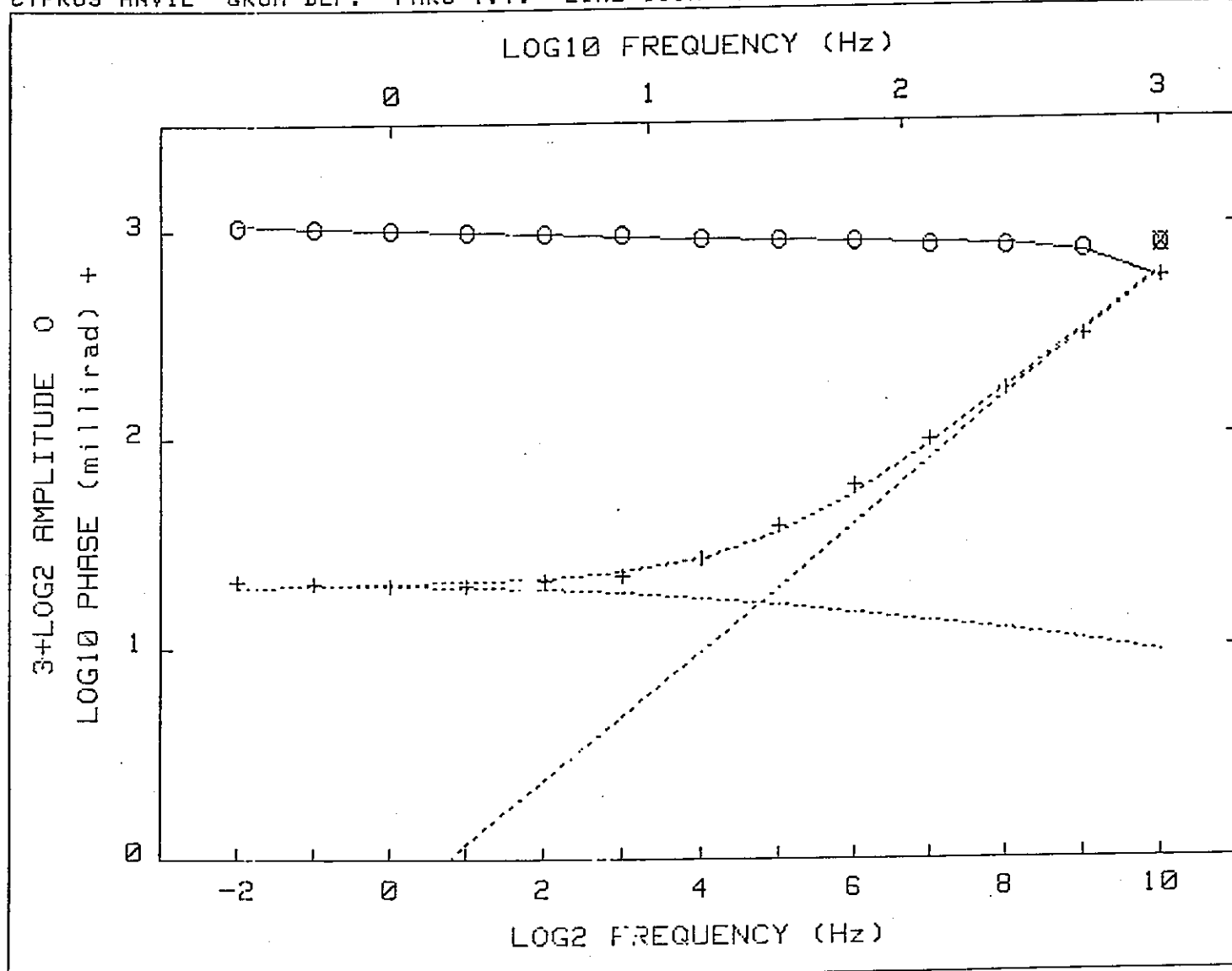
CRL: Number of dispersions= 2  
 C1=.25 M2=1 fixed

Iter	Lambda	Rchsq	R0	M1	T1	T2	C2
0	1.E-02	.03050	1.540	.367	4.8E+02	1.1E-05	.892
1	1.E-02	.00219	1.377	.310	4.0E+02	1.4E-05	.707
2	1.E-03	.00081	1.288	.269	1.7E+02	2.4E-05	.777
3	1.E-03	.00057	1.236	.237	6.3E+01	2.6E-05	.789
4	1.E-03	.00041	1.192	.211	2.1E+01	2.8E-05	.805
5	1.E-03	.00026	1.160	.192	7.7E+00	3.0E-05	.825
6	1.E-04	.00025	1.102	.162	8.2E-01	3.5E-05	.870
7	1.E-05	.00006	1.112	.171	7.2E-01	3.5E-05	.886
8	1.E-06	.00006	1.112	.171	7.1E-01	3.5E-05	.887

Pct Std Deviations .4 1.1 29.9 2.7 1.2

Correlation Matrix

	1.000				
	.757	1.000			
	.896	.674	1.000		
	-.482	-.212	-.557	1.000	
	-.573	-.198	-.692	.906	1.000



CRL: Number of dispersions= 2  
 C1=.25 M2=1 C2=1 fixed

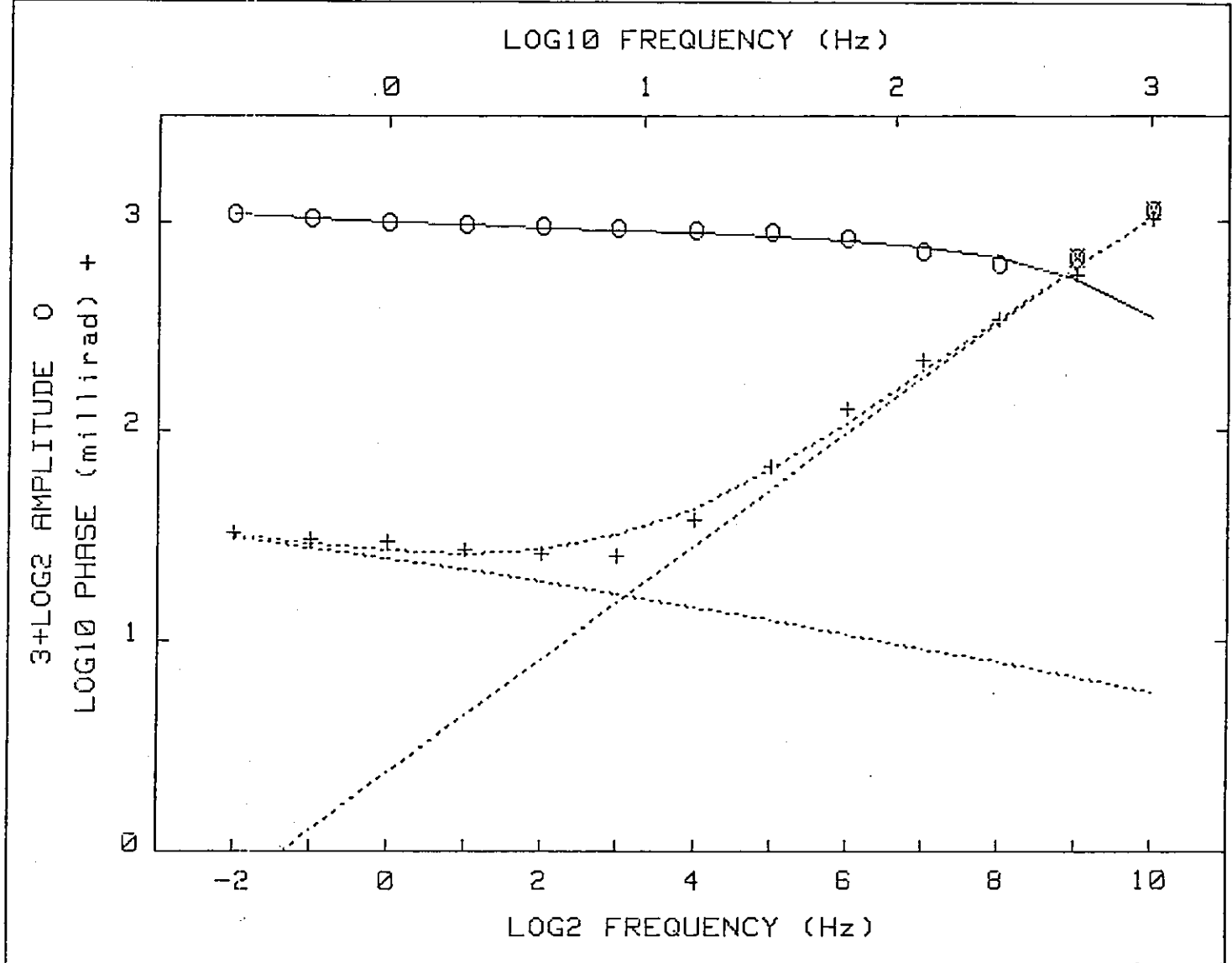
Iter	Lambda	Rchsq	R0	M1	T1	T2
0	1.E-02	.00047	1.129	.187	5.6E-01	7.9E-05
1	1.E-02	.00035	1.118	.185	5.2E-01	8.3E-05
2	1.E-03	.00033	1.113	.183	4.0E-01	8.2E-05
3	1.E-04	.00033	1.111	.183	3.4E-01	8.2E-05

Pct Std Deviations .8 1.9 53.7 2.3

Correlation Matrix

1.000				
.709	1.000			
.825	.565	1.000		
.084	-.189	.217	1.000	

Apparent Resistivity Measured at 1 Hz is 810.3  
 Apparent Resistivity Calculated from Inductive Coupling is 2216



CRL: Number of dispersions= 2  
 C1=.25 M2=1 fixed

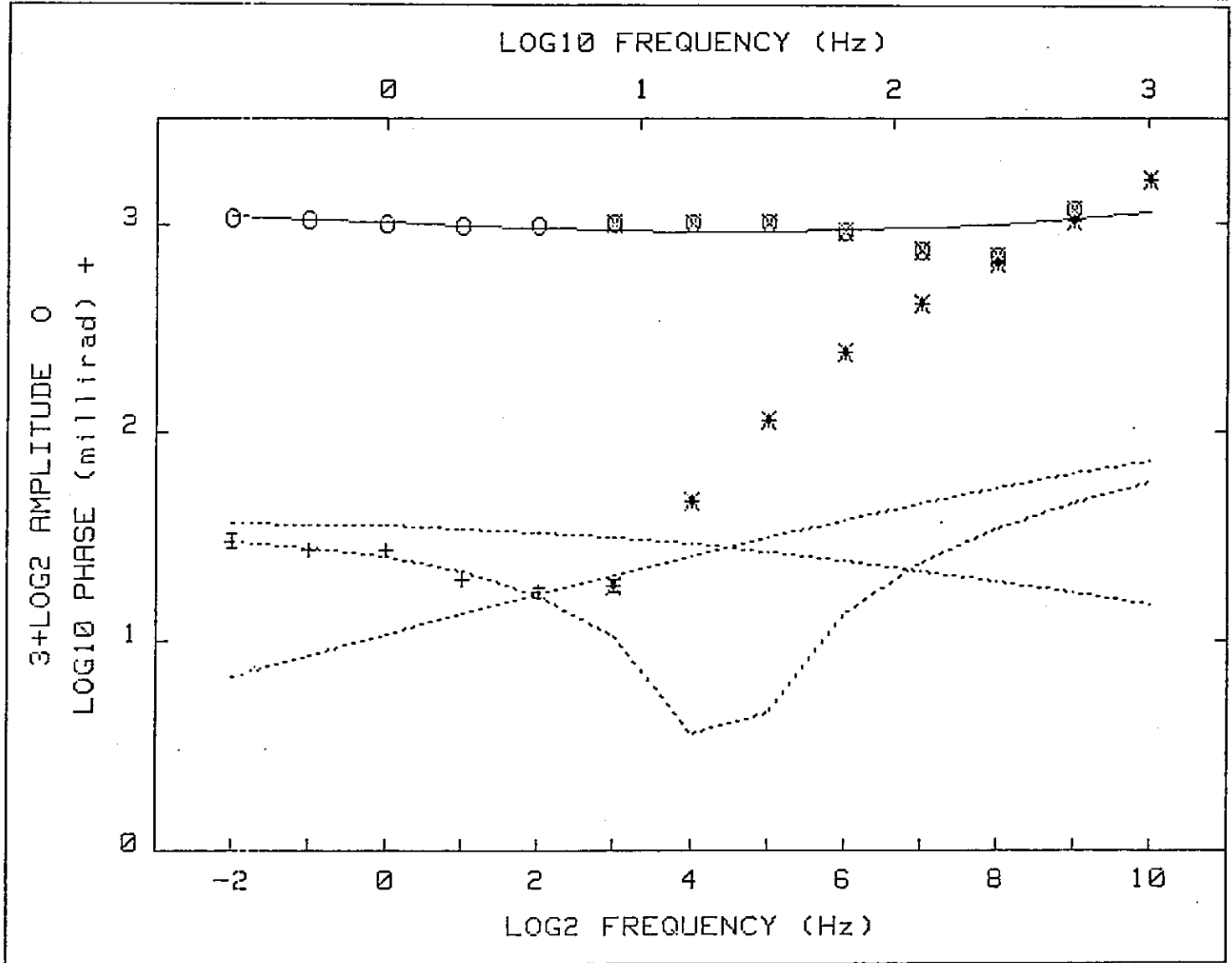
Iter	Lambda	Rchsq	R0	M1	T1	T2	C2
0	1.E-02	.00297	1.320	.283	4.2E+01	1.6E-04	.931
1	1.E-02	.00163	1.332	.309	4.5E+01	1.4E-04	.917
2	1.E-03	.00153	1.402	.341	1.3E+02	1.3E-04	.903
3	1.E-03	.00148	1.461	.366	2.3E+02	1.2E-04	.899
4	1.E-03	.00146	1.506	.384	3.6E+02	1.1E-04	.895
5	1.E-03	.00145	1.541	.397	4.8E+02	1.1E-04	.892

Pct Std Deviations            36.2    61.3    999.9    56.2    6.1

Correlation Matrix

	1.000				
	1.000	1.000			
	.998	.998	1.000		
	-.988	-.987	-.988	1.000	
	-.880	-.874	-.889	.923	1.000

Apparent Resistivity Measured at 1 Hz is 404.0  
 Apparent Resistivity Calculated from Inductive Coupling is 2162



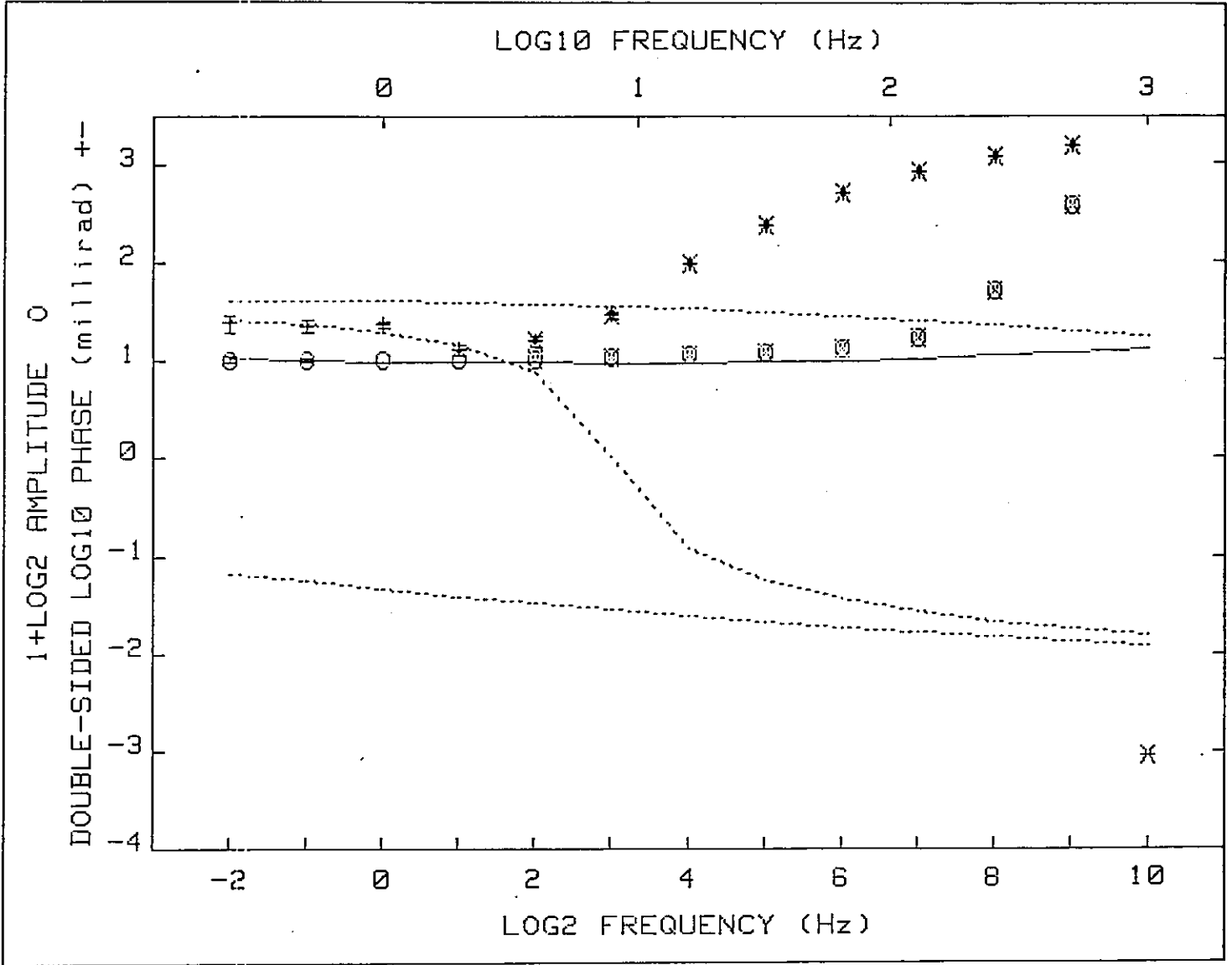
CRL: Number of dispersions= 2 Negative.  
 T1=1 C1=.25 M2=1 fixed

Iter	Lambda	Rchsqr	R0	M1	T2	C2
0	1.E-02	.02618	1.310	.457	5.9E-05	.454
1	1.E-02	.00206	1.169	.272	4.9E-05	.660
2	1.E-03	.00067	1.197	.267	3.9E-05	.651
3	1.E-04	.00062	1.203	.278	8.0E-06	.500
4	1.E-04	.00056	1.207	.287	3.6E-06	.432
5	1.E-04	.00055	1.210	.294	1.9E-06	.391
6	1.E-04	.00054	1.212	.299	1.2E-06	.361
7	1.E-04	.00053	1.214	.304	7.7E-07	.339
8	1.E-04	.00053	1.215	.309	5.5E-07	.320

Pct Std Deviations . . . . . 2.8 41.3 999.9 . 296.7

Correlation Matrix

	1.000			
	.954	1.000		
	-.916	-.985	1.000	
	-.928	-.992	.999	1.000



CRL: Number of dispersions= 2 Negative.  
 T1=1 C1=.25 M2=1 fixed

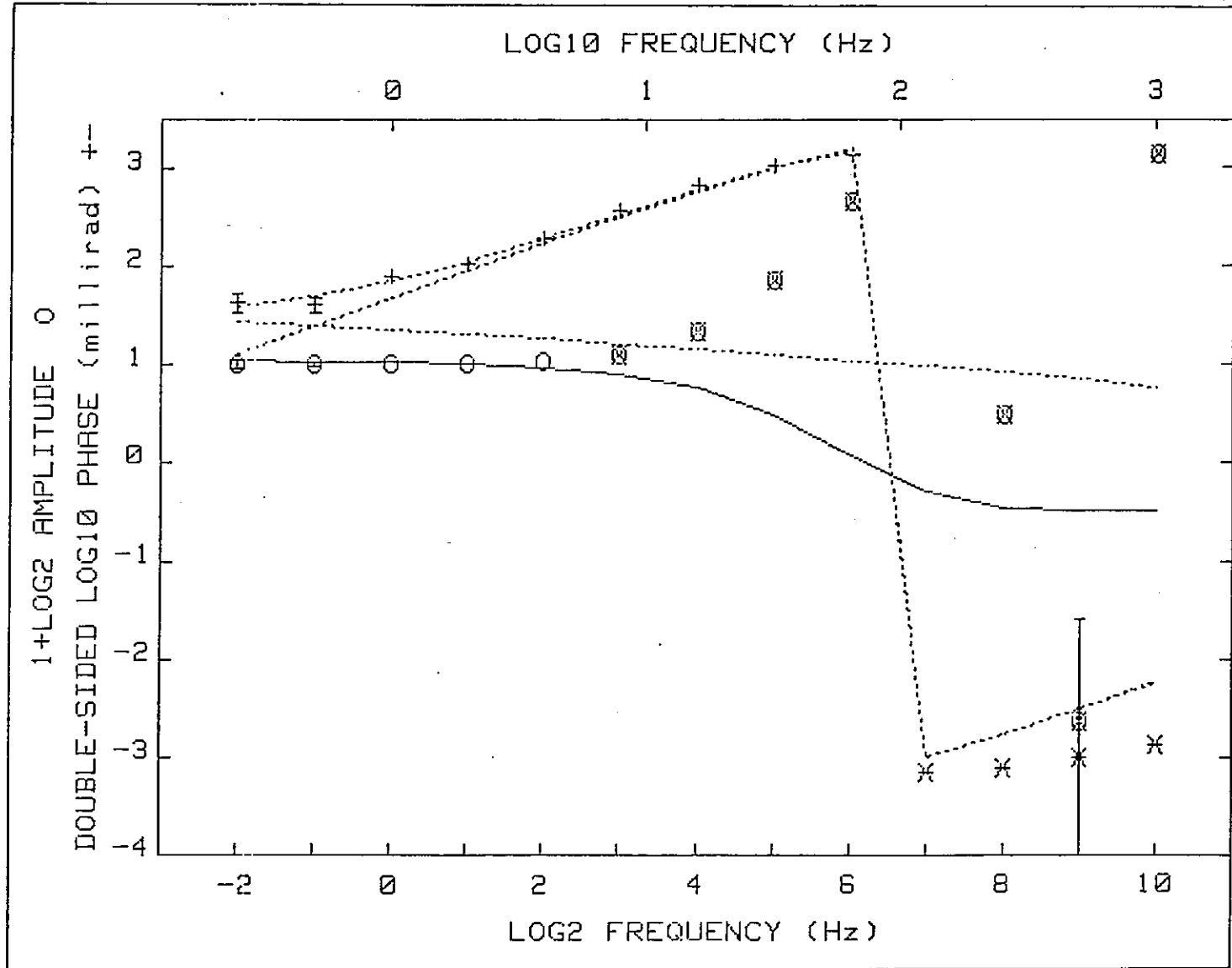
Iter	Lambda	Rchsq	R0	M1	T2	C2
0	1.E-02	.00944	1.210	.309	8.5E-07	.320
1	1.E-02	.00290	1.205	.332	8.8E-07	.260
2	1.E-03	.00232	1.211	.353	1.1E-06	.256
3	1.E-03	.00230	1.211	.345	1.3E-06	.267

Pct Std Deviations 6.6 956.7 999.9 999.9

Correlation Matrix

	1.000			
	.931	1.000		
	-.898	-.994	1.000	
	-.916	-.999	.998	1.000

Apparent Resistivity Measured at 1 Hz is 155.6  
 Apparent Resistivity Calculated from Inductive Coupling is 176.6



CRL: Number of dispersions= 2  
 C1=.25 M2=1 fixed

Iter	Lambda	Rchsq	R0	M1	T1	T2	C2
0	1.E-02	.00302	1.345	.281	4.1E+01	4.6E-03	.933
1	1.E+01	.00247	1.330	.282	4.1E+01	4.6E-03	.931
2	1.E+01	.00234	1.323	.283	4.1E+01	4.6E-03	.931
3	1.E+03	.00234	1.323	.283	4.1E+01	4.6E-03	.931

Pct Std Deviations            73.6   187.8   999.9   104.4   18.5

Correlation Matrix

	1.000				
	.997	1.000			
	.997	.991	1.000		
	-.994	-.994	-.990	1.000	
	-.960	-.944	-.971	.960	1.000

Apparent Resistivity Measured at 1 Hz is 87.69  
 Apparent Resistivity Calculated from Inductive Coupling is 190.5

As expected, the measurements with  $X = 1000'$  (305 meters) show a more clearly defined anomaly at depth. For these deep measurements the apparent IP effects within the anomaly at depth are greater than the high background IP effects to the north. In addition, there are a few, shallow, above background IP effects at the south edge of the measured data (50+00S to 60+00S).

No useful spectral interpretation can be made from the  $n = 5$  or  $n = 6$  measurements for either  $X = 500'$  or  $X = 1000'$ . As the enclosed examples show, the voltage levels were low in magnitude and the electrical noise levels were high in magnitude. The parameters for the IP dispersion for the deep anomalous source can not be determined from the present data.

The higher background IP effects to the north and the slightly anomalous IP effects at the south edge of the data (for  $X = 1000'$ ) were both detected for  $n = 1$  and  $n = 2$ . For these measurements, the measurements were more reliable; the standard deviations for the data were much smaller.

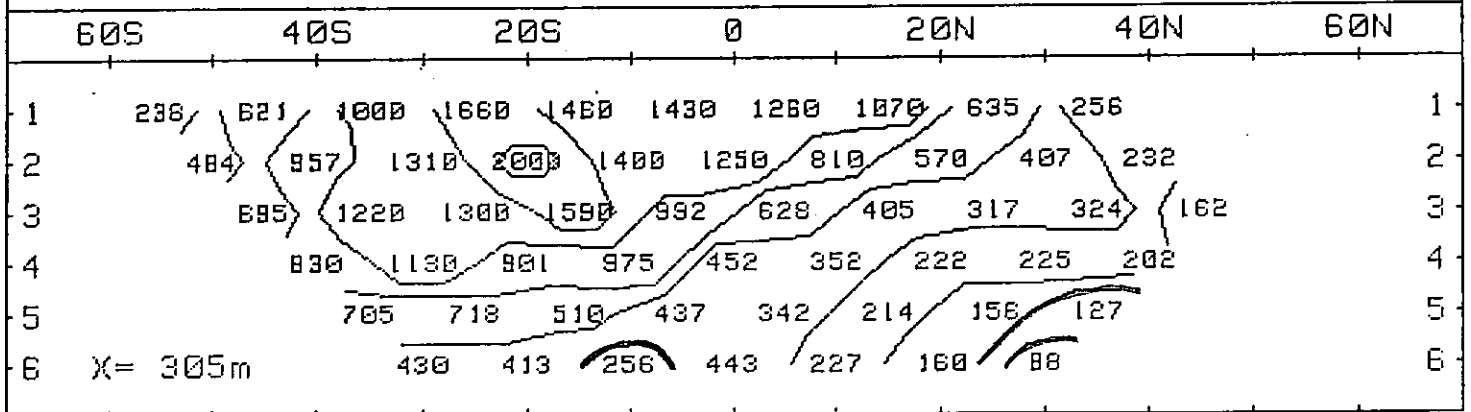
For both of these anomalies (D1, N1; D1, N2) and (D10, N1; D10, N2), the spectral IP data plots can be accurately inverted. The attached data plots show the accuracy with which the data can be approximated by two Cole-Cole Dispersions.

For both anomalies, the time-constants ( $T_1$ ) for the IP dispersion are in the range from ten seconds to one hundred seconds. Our field tests elsewhere in the Faro Area indicates that these large time-constant dispersions are due to large grain-sized sources, usually graphite.

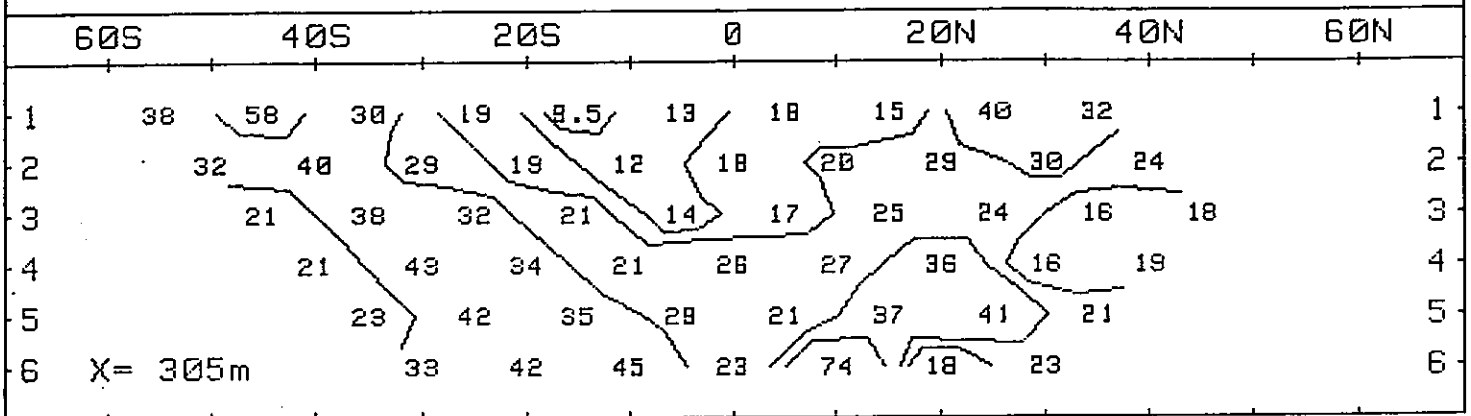
#### 4. CONCLUSIONS AND RECOMMENDATIONS

The Spectral IP Test Survey on Line 108W, over the Grum Deposit has shown that the deep source in the Faro Area can be detected if the

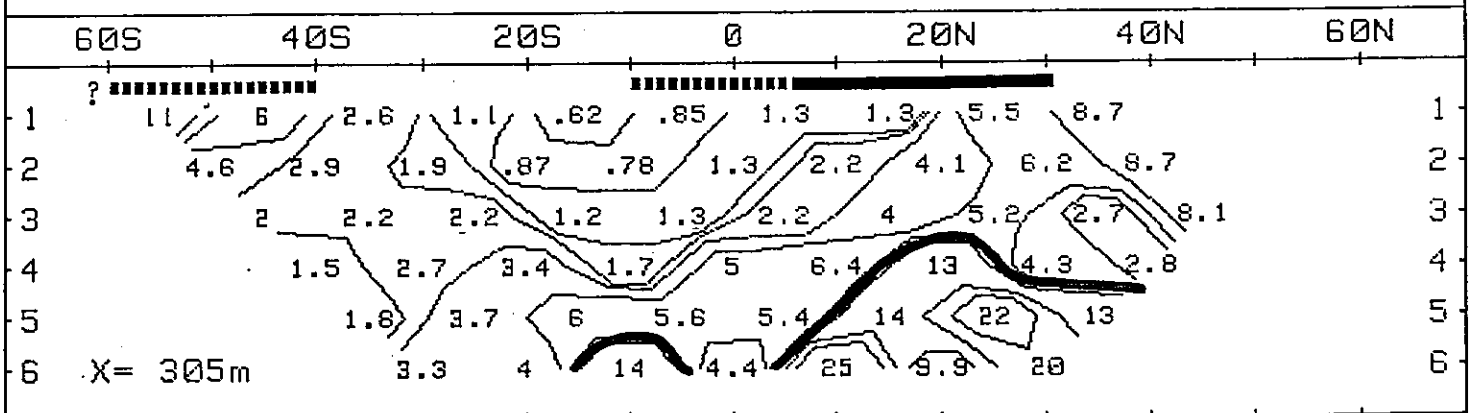
PHOENIX SPECTRAL IP OBS. RESISTIVITY (ohm-m) 1 Hz  
 CYPRUS ANVIL GRUM DEP. FARO Y.T. LINE 108W 16/9/81 L330



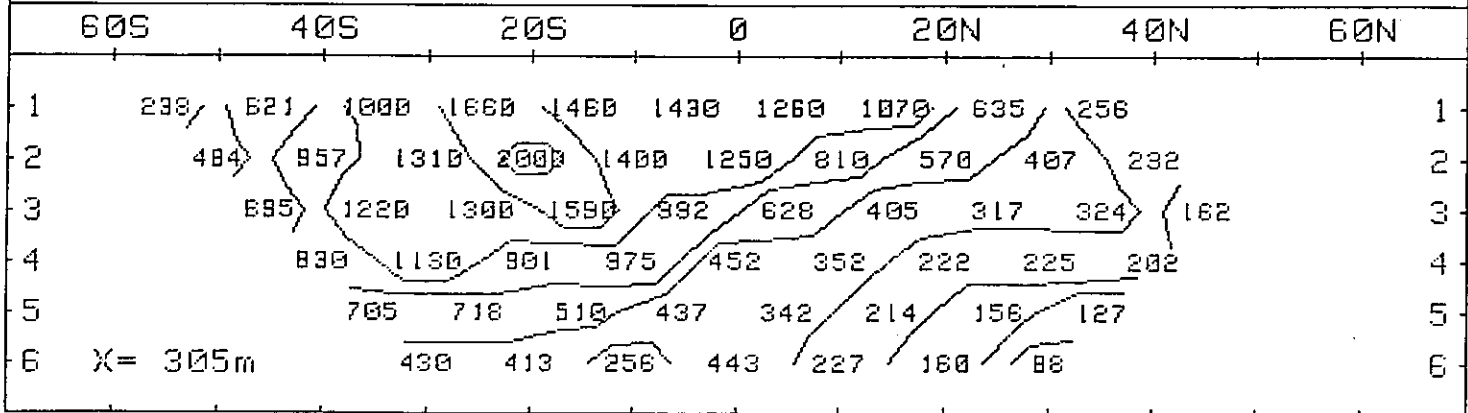
PHOENIX SPECTRAL IP DECOUPLED PHASE (mrad) 1 Hz  
 CYPRUS ANVIL GRUM DEP. FARO Y.T. LINE 108W 16/9/81 L330



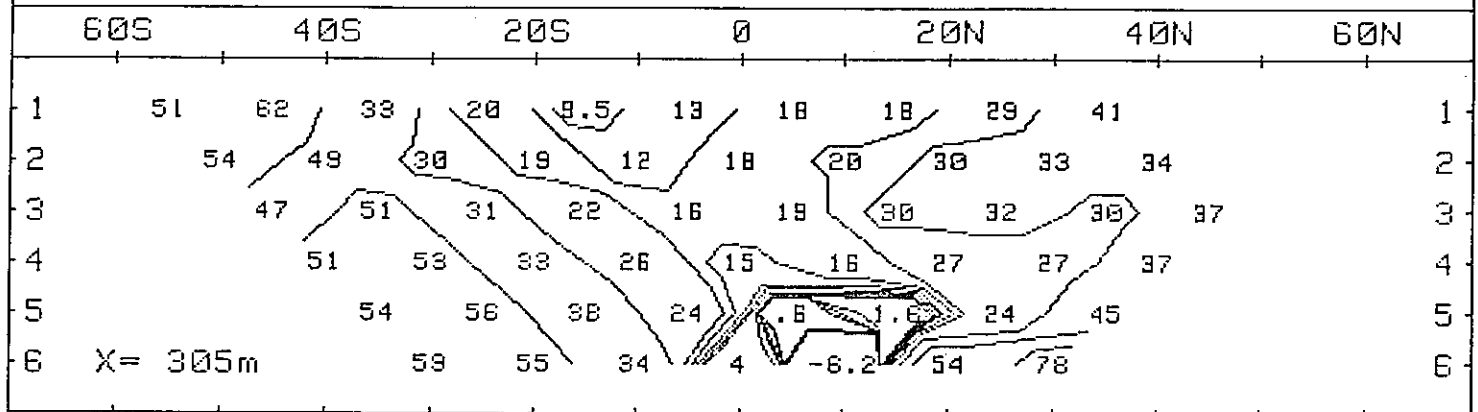
PHOENIX SPECTRAL IP DECOUPLED METAL FACTOR 1Hz  
 CYPRUS ANVIL GRUM DEP. FARO Y.T. LINE 108W 16/9/81 L330



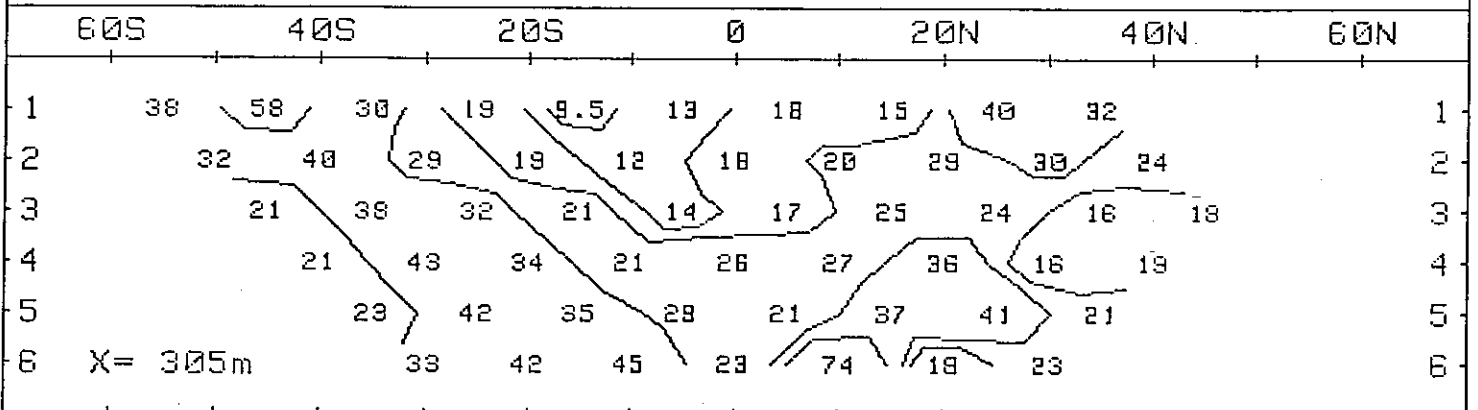
PHOENIX SPECTRAL IP OBS. RESISTIVITY (ohm-m) 1 Hz  
 CYPRUS ANVIL GRUM DEP. FARO Y.T. LINE 108W 16/9/81 L330

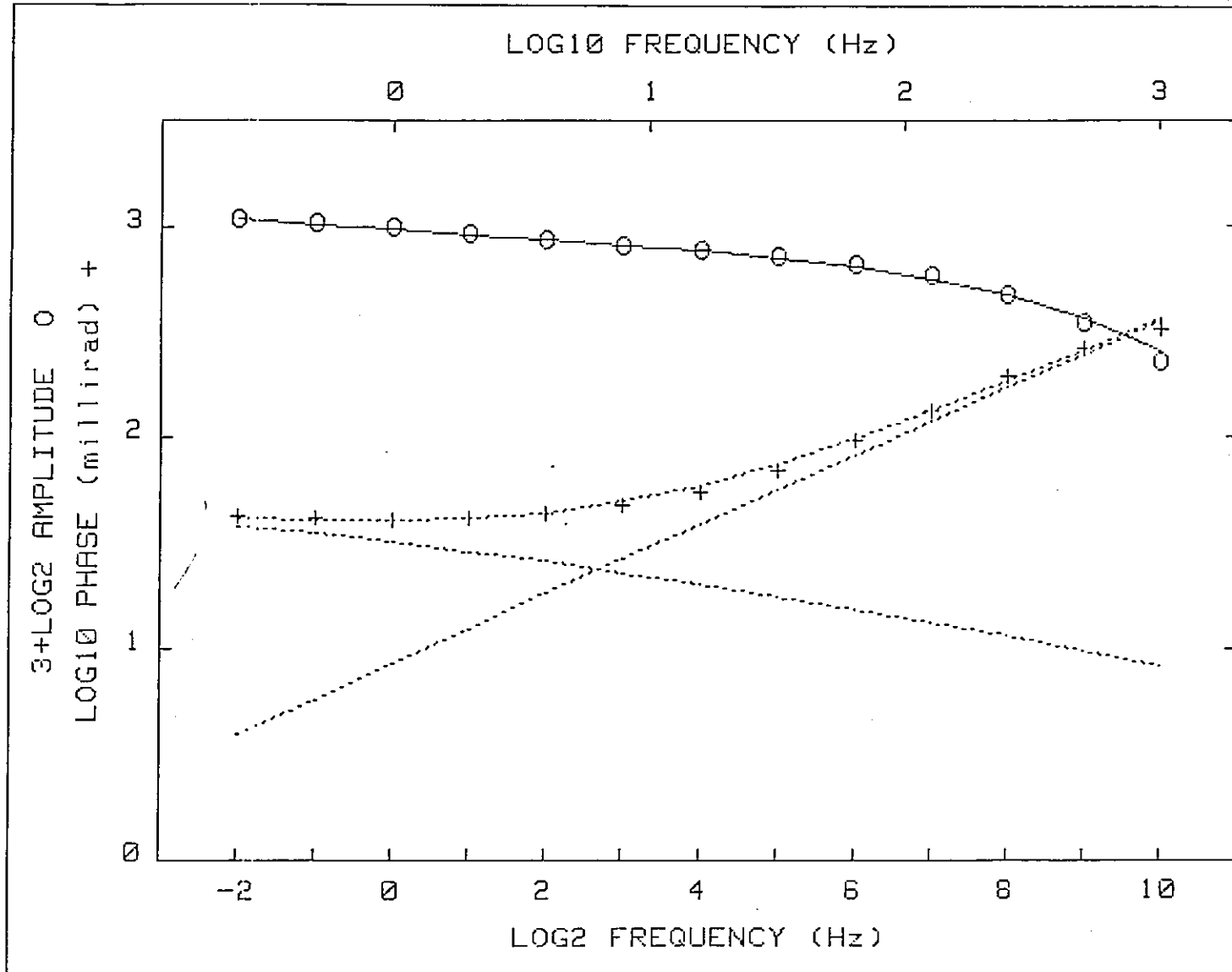


PHOENIX SPECTRAL IP PHASE (mrad) at 1 Hz  
 CYPRUS ANVIL GRUM DEP. FARO Y.T. LINE 108W 16/9/81 L330



PHOENIX SPECTRAL IP DECOUPLED PHASE (mrad) 1 Hz  
 CYPRUS ANVIL GRUM DEP. FARO Y.T. LINE 108W 16/9/81 L330





CRL: Number of dispersions= 2  
 C1=.25 M2=1 fixed

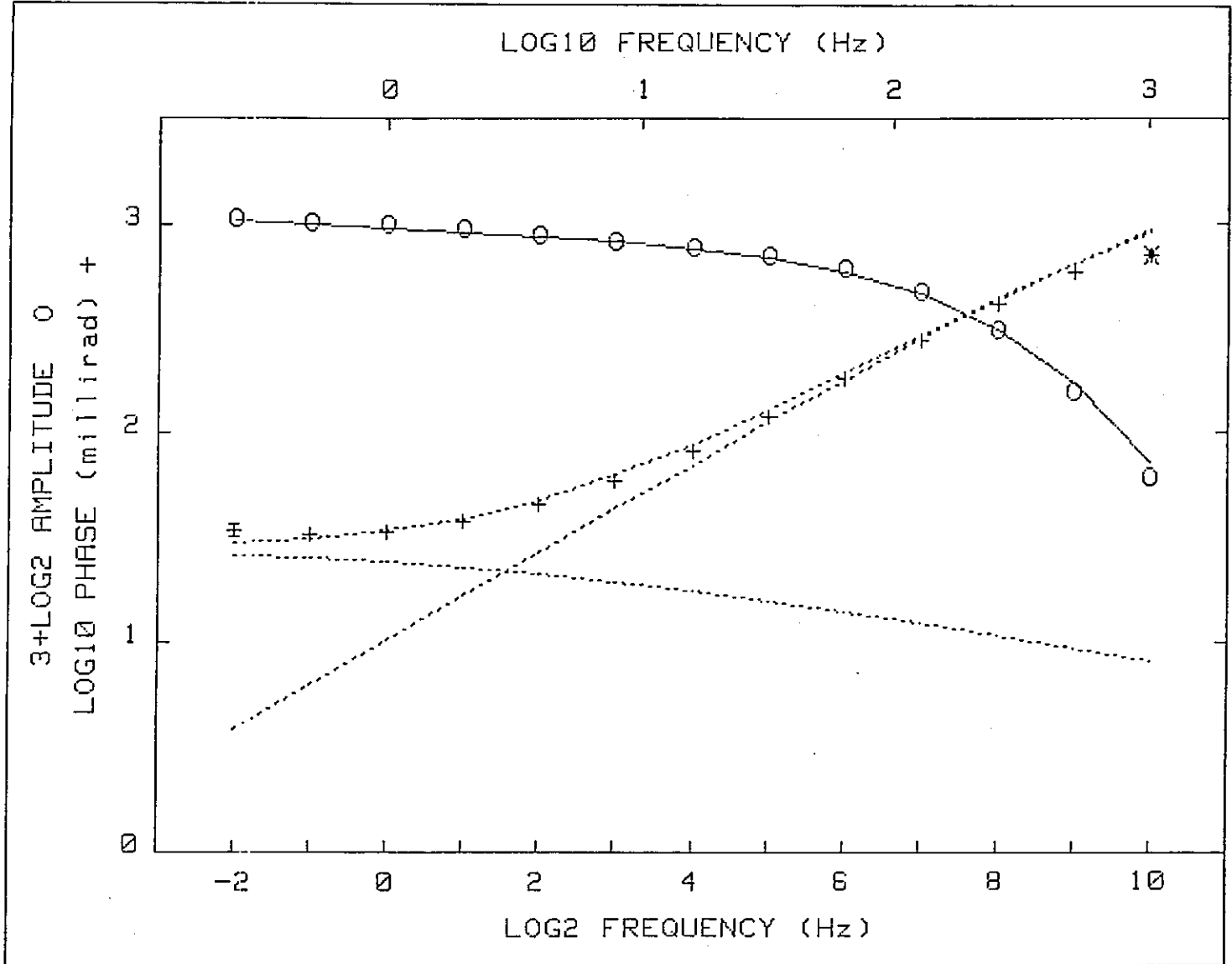
Iter	Lambda	Rchsq	R0	M1	T1	T2	C2
0	1.E-02	.00762	1.517	.411	4.5E+01	2.0E-05	.497
1	1.E-02	.00056	1.410	.368	4.6E+01	2.0E-05	.542
2	1.E-03	.00047	1.422	.368	4.8E+01	2.0E-05	.543
3	1.E-04	.00047	1.441	.375	6.1E+01	1.9E-05	.540

Pct Std Deviations            14.8    21.9    999.9    44.0    7.1

Correlation Matrix

	1.000					
	.997	1.000				
	.997	.990	1.000			
	-.979	-.972	-.980	1.000		
	-.982	-.876	-.920	.952	1.000	

Apparent Resistivity Measured at 1 Hz is 255.5  
 Apparent Resistivity Calculated from Inductive Coupling is 289.7



CRL: Number of dispersions= 2  
 C1=.25 M2=1 fixed

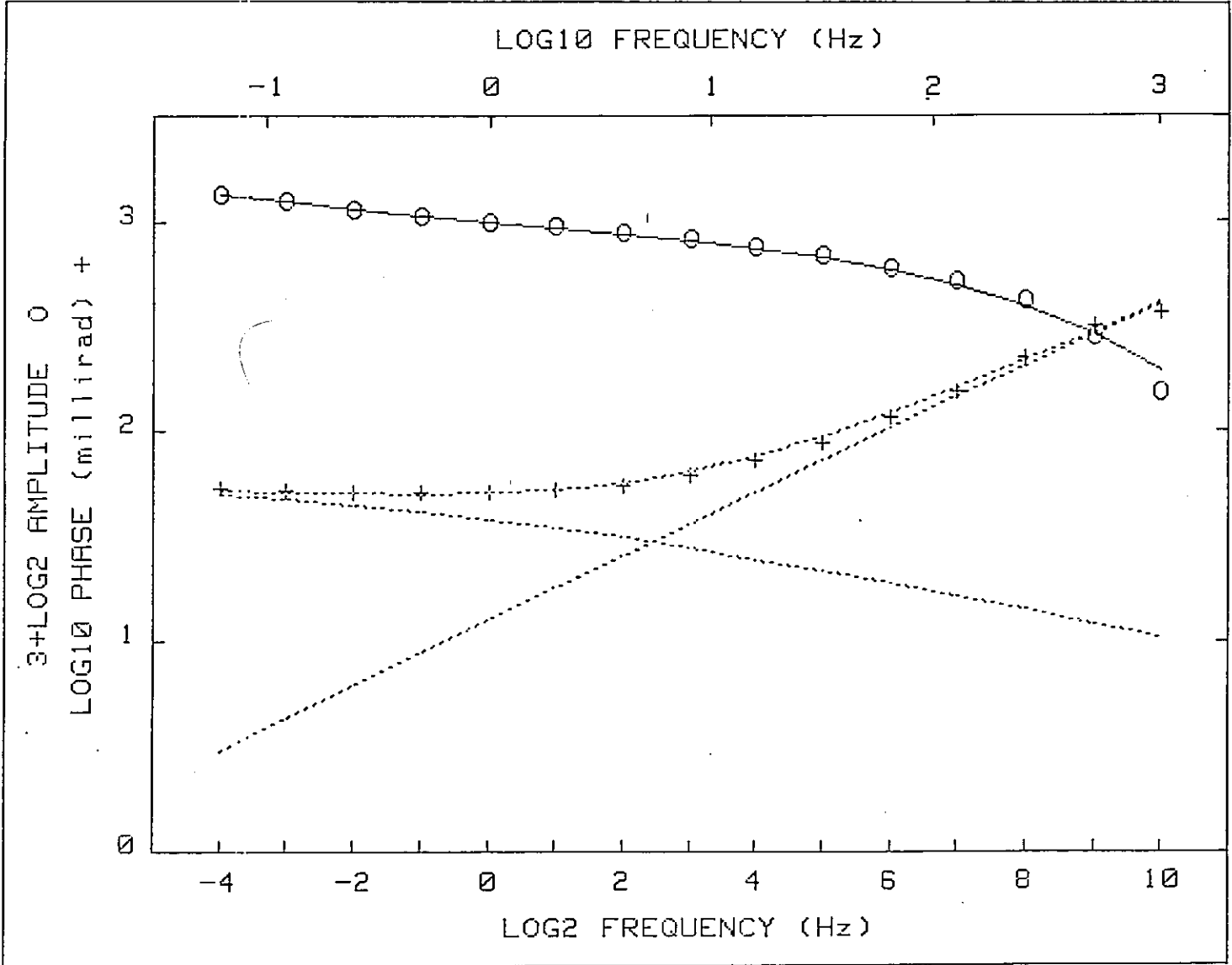
Iter	Lambda	Rchsq	R0	M1	T1	T2	C2
0	1.E-02	.00510	1.180	.225	1.0E+02	1.8E-04	.662
1	1.E-02	.00175	1.297	.293	7.8E+01	1.7E-04	.660
2	1.E-03	.00129	1.190	.235	9.6E+00	1.9E-04	.671
3	1.E-03	.00100	1.193	.239	4.7E+00	1.9E-04	.690
4	1.E-03	.00099	1.183	.235	3.4E+00	1.9E-04	.692

Pct Std Deviations                    6.9    12.8    999.9    9.8    5.1

Correlation Matrix

	1.000				
	.947	1.000			
	.987	.910	1.000		
	-.821	-.777	-.825	1.000	
	-.799	-.621	-.831	.862	1.000

Apparent Resistivity Measured at 1 Hz is 232.4  
 Apparent Resistivity Calculated from Inductive Coupling is 252.7



CRL: Number of dispersions= 2  
 C1=.25 M2=1 fixed

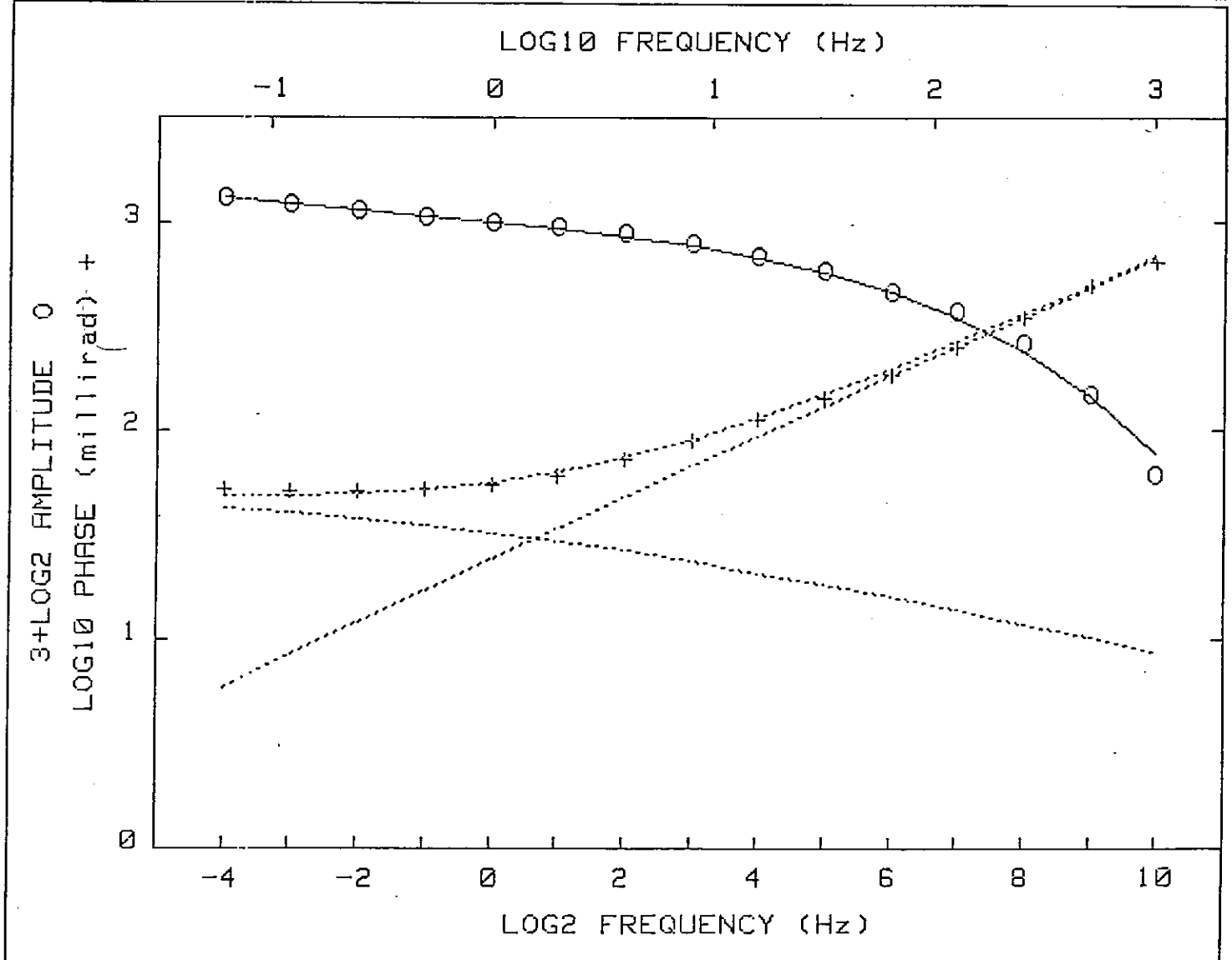
Iter	Lambda	Rchsq	R0	M1	T1	T2	C2
0	1.E-02	.02755	1.453	.365	4.2E+01	7.6E-05	.489
1	1.E-02	.00181	1.469	.388	3.5E+01	2.4E-05	.471
2	1.E-03	.00081	1.502	.409	3.6E+01	2.0E-05	.501
3	1.E-04	.00079	1.544	.418	6.4E+01	1.9E-05	.490
4	1.E-05	.00079	1.517	.411	4.5E+01	2.0E-05	.497

Pct Std Deviations            11.0    11.6    539.4    36.5    7.4

Correlation Matrix

1.000				
.988	1.000			
.992	.967	1.000		
-.952	-.929	-.951	1.000	
-.849	-.777	-.880	.927	1.000

Apparent Resistivity Measured at 1 Hz is 237.9  
 Apparent Resistivity Calculated from Inductive Coupling is 169.5



CRL: Number of dispersions= 2  
 C1=.25 M2=1 fixed

Iter	Lambda	Rchsq	R0	M1	T1	T2	C2
0	1.E-02	.01336	1.394	.311	1.5E+02	2.1E-04	.541
1	1.E-02	.00179	1.480	.362	1.1E+02	8.7E-05	.467
2	1.E-03	.00079	1.471	.370	5.5E+01	7.3E-05	.483
3	1.E-03	.00079	1.453	.365	4.2E+01	7.6E-05	.489

Pct Std Deviations            12.8    14.5    999.9    33.4    8.2

Correlation Matrix

	1.000				
	.985	1.000			
	.994	.966	1.000		
	-.976	-.972	-.966	1.000	
	-.930	-.868	-.945	.933	1.000

Apparent Resistivity Measured at 1 Hz is 483.9  
 Apparent Resistivity Calculated from Inductive Coupling is 105.9

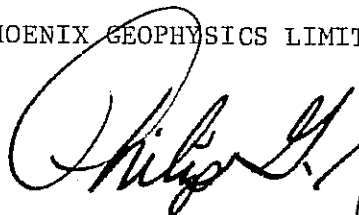
proper electrode interval is used. However, the measurements are difficult to make due to the low signal voltage levels and the high electrical noise levels.

These spectral plots are too unreliable to permit accurate inversions to be made to determine the electrical parameters of the IP dispersion. This is one of the more useful features of the Spectral IP Method and steps have been taken to assure better low frequency data. In further work in the Faro Area, a high-power IP transmitter will be used (15KVA to 30KVA). The higher current levels this power will permit will give higher voltage levels and smaller standard deviations.

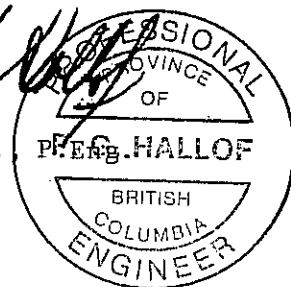
A further step that would aid with the interpretation of deep, indistinct anomalies, would be to double the density of the data on the pseudosection plotter. This can be accomplished by, for instance, repeating the X = 1000' measurements with the electrode positions shifted by 500'.

With these extra steps, the deep anomalies expected from the deep sources in the Faro Area should be more easily interpreted.

PHOENIX GEOPHYSICS LIMITED



Philip G. Hallof, Ph.D.  
Geophysicist



Expiry Date: February 25, 1982

Dated: February 2, 1982


CERTIFICATE

I, Philip George Hallof, of the City of Toronto, Province of Ontario, do hereby certify that:

1. I am a geophysicist residing at Suite 3505, 2045 Lakeshore Blvd. W., Toronto, Ontario.
2. I am a graduate of the Massachusetts Institute of Technology with a B.Sc. Degree (1952) in Geology and Geophysics, and a Ph.D. Degree (1957) in Geophysics.
3. I am a member of the Society of Exploration Geophysicists and the European Association of the Exploration Geophysicists.
4. I am a Professional Geophysicist, registered in the Province of Ontario, the Province of British Columbia and the State of Arizona.
5. I have no direct or indirect interest, nor do I expect to receive any interest directly or indirectly, in the property or securities of Cyprus-Anvil Mines Ltd., or any affiliate.
6. The statements made in this report are based on a study of published geological literature and unpublished private reports.
7. Permission is granted to use in whole or in part for assessment and qualification requirements but nor for advertising purposes.

Dated at Toronto

This 2nd day of February, 1982

  
Philip G. Hallof, Ph.D.



Expiry Date: February 25, 1982

RECENT AND FUTURE ADVANCES  
IN THE  
INDUCED POLARIZATION METHOD

by

Philip G. Hallof, Ph.D., and William H. Pelton, Ph.D.

Prepared for Presentation  
at the Prospectors and Developers Convention  
Royal York Hotel, Toronto  
March 8th to March 11th, 1981

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RECENT AND FUTURE ADVANCES  
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INTRODUCTION

The induced polarization (IP) method was introduced into Canadian exploration practice in the period from 1955 to 1960. Two measurement techniques were widely used, the pulse-transient method (time domain) and the variable frequency method (frequency domain). In the fifteen years following 1960, striking improvements were made in both frequency domain and time domain IP equipment, but only limited progress was made in a better understanding of the IP phenomenon, its source, and how to use it.

However, in the five year period since 1975, considerable progress has been made in our understanding of the IP method. Beginning with an understanding of the exact equivalence of the IP measurements in the phase-domain, the frequency-domain and the time-domain (see Figure 1), we have progressed to the study of the IP effect over the entire frequency range of interest. These studies have led to a greatly advanced understanding of the IP method in at least three important areas:-

- i) Our understanding of the IP phenomenon itself has been vastly increased. A greater knowledge of the source of the IP effect and its detailed behavior has suggested additional uses for the method, beyond the simple detection of anomalies.
- ii) A better idea of what we wanted to measure, and modern solid-state electronics, have resulted in greatly improved measurement techniques and instrumentation.

iii) The advent of smaller, less expensive digital computers has greatly helped with the interpretation of IP field data. Rapid forward problem solutions and the possibility of direct inversion of IP field data have made it possible for the geophysicists to give the exploration geologist a much better picture of just what the source of a particular IP anomaly might be.

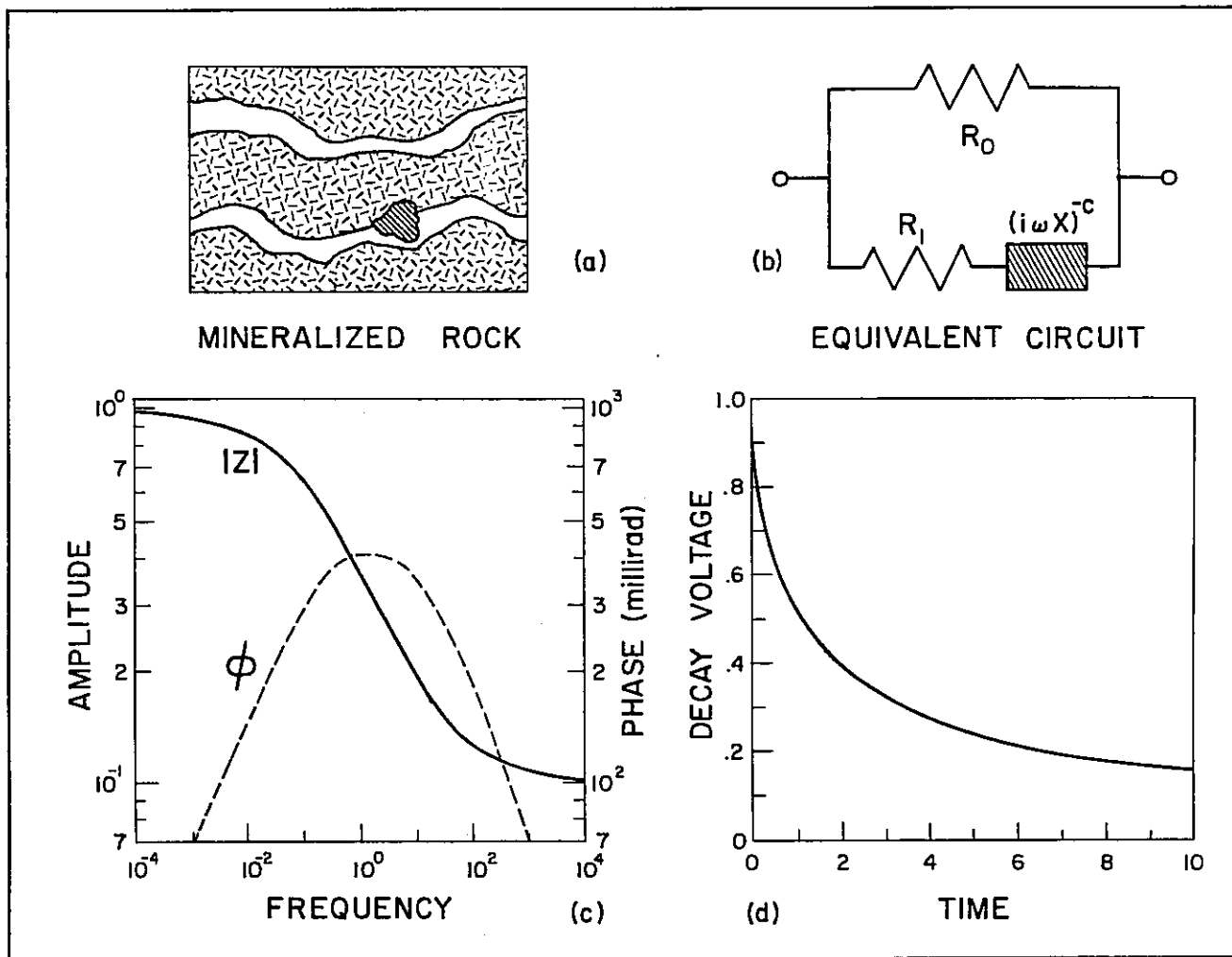


FIG. 1

WHAT IS THE NATURE OF THE IP EFFECT

Spectral induced polarization measurements (the measurement of the phase shifts over a wide frequency range) in open-pit mines, gave the first information concerning the detailed nature of the IP phenomenon. It became clear that four parameters, not just two (resistivity and IP

effect) were necessary to completely describe the IP effect.

The equation shown on Figure 2 and Figure 3 is formally known as the Cole-Cole Dispersion Equation; the four electrical parameters are:

- $R_0$  - the dc resistivity value
- $m$  - the IP effect (in non-dimensional form)
- $\tau$  - the time constant of the IP effect
- $c$  - the exponent of the frequency ( $\omega$ )

Since the very beginning of our experience with spectral IP measurements, we have found that these four parameters, and the Cole-Cole equation will adequately describe any IP effect that has been measured. The first measurements were made using one meter electrode intervals, within open-pit mines. The circles on Figure 4 and Figure 5 show the data points measured over massive graphite and massive sulphides at the Anvil Mine. As the solid line curves on the drawings show, the measured data points can be almost exactly replaced using a Cole-Cole Dispersion.

The four parameters obtained by the computer inversion of the graphite data are typical of those that have been measured over all graphite deposits. The IP effect is very large and the time constant ( $\tau$ ) is much greater than 1.0 seconds.

Large IP effects are also measured over massive sulphide sources. However, invariably the time constant value for ( $\tau$ ) is less than 1.0 seconds. The measurements shown on Figure 6 are for a massive sulphide source at the Kidd Creek Mine; again there is a large IP effect, and a small time constant.

The critical frequency ( $F_c$ ) for a Cole-Cole Dispersion is defined as that frequency for which the maximum phase-shift is measured. An examination of the equation and curves in Figure 2 shows that as the time-constant ( $\tau$ ) is increased by a factor of ten, the critical frequency ( $F_c$ ) will decrease by a factor of ten. Therefore, we can expect that for the spectral IP response of polarizable sources ( $F_c$ ) will be inversely proportional to ( $\tau$ ).

An understanding of the physical parameters that affect the time-constant ( $\tau$ ) of a spectral IP response has come from two areas of research. One is the study of the spectral IP response from artificial rocks; these samples are prepared using metallic particles of known material and size (see Figure 8 for example). The second line of attack has been to formulate the mathematical expressions that might describe the macro-characteristics of conduction in a mineralized rock, with certain simplifying assumptions. Initially spherical particles were assumed for this work and an example of the predicted spectral results is shown on Figure 7.

We have shown these two different sets of data, because they are very similar in character. Further, they show the same result that all of

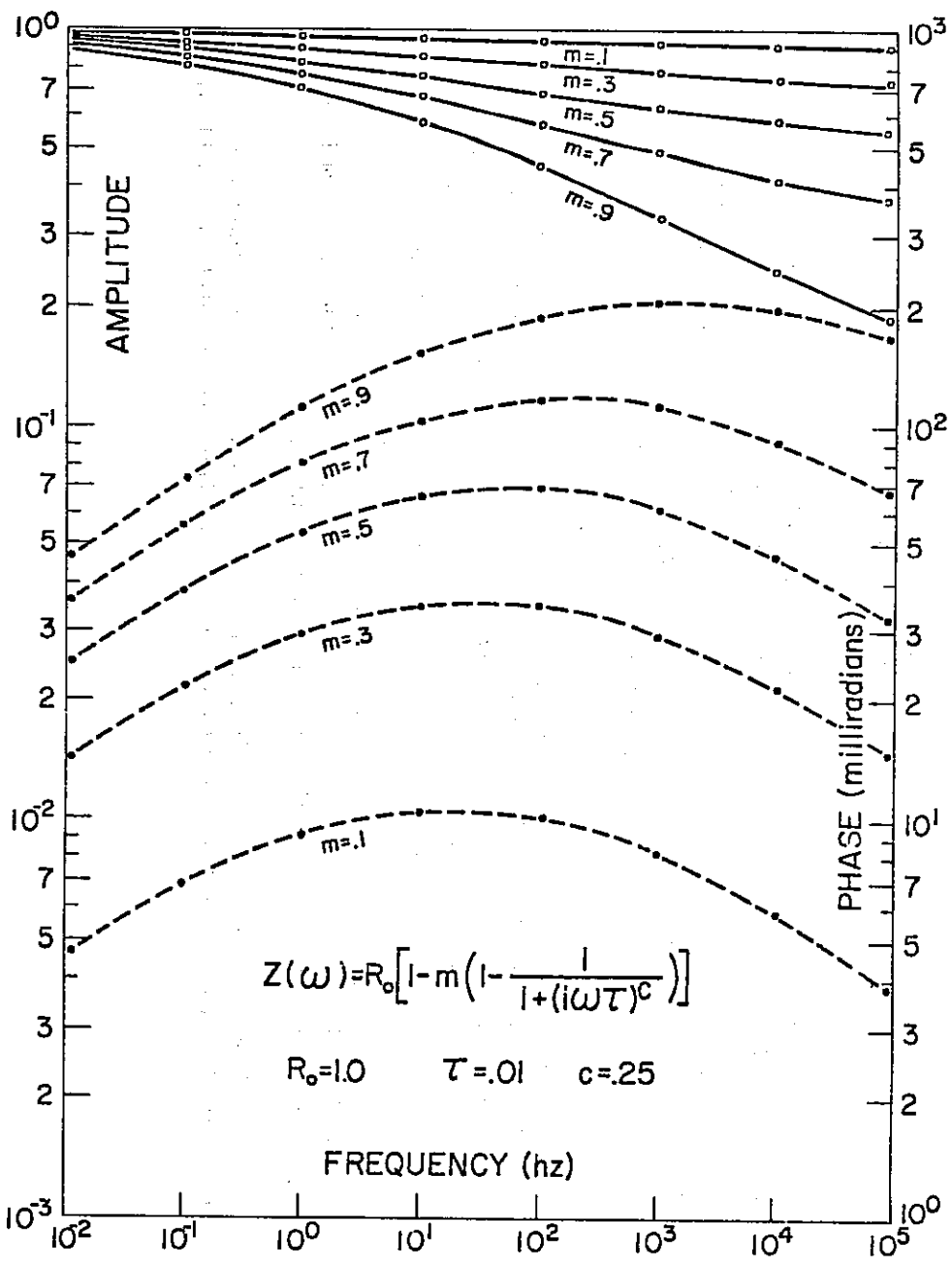


FIG. 2

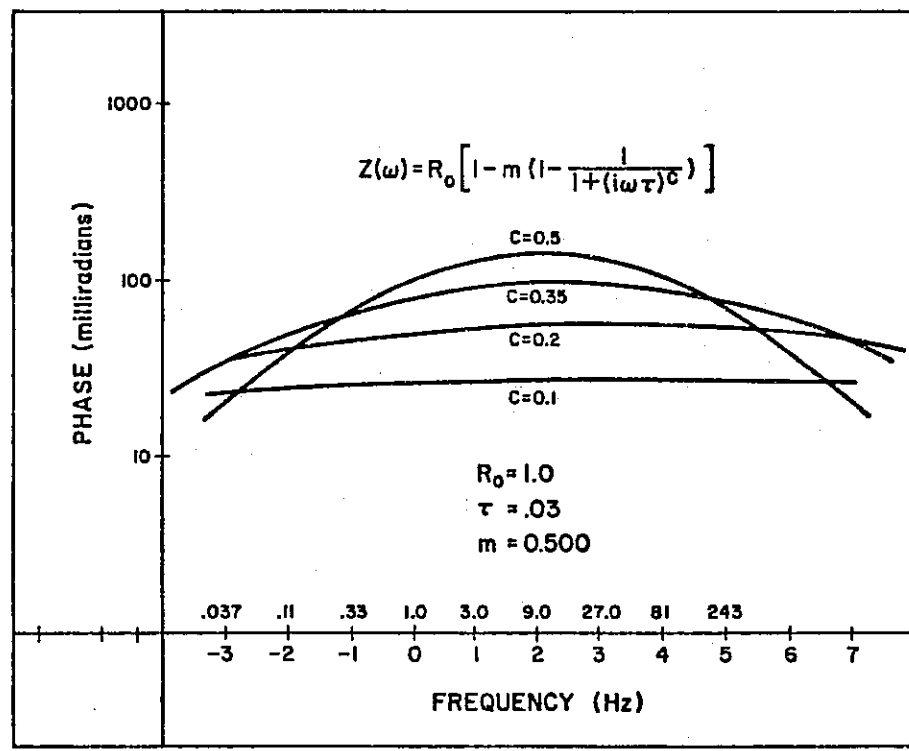


FIG. 3

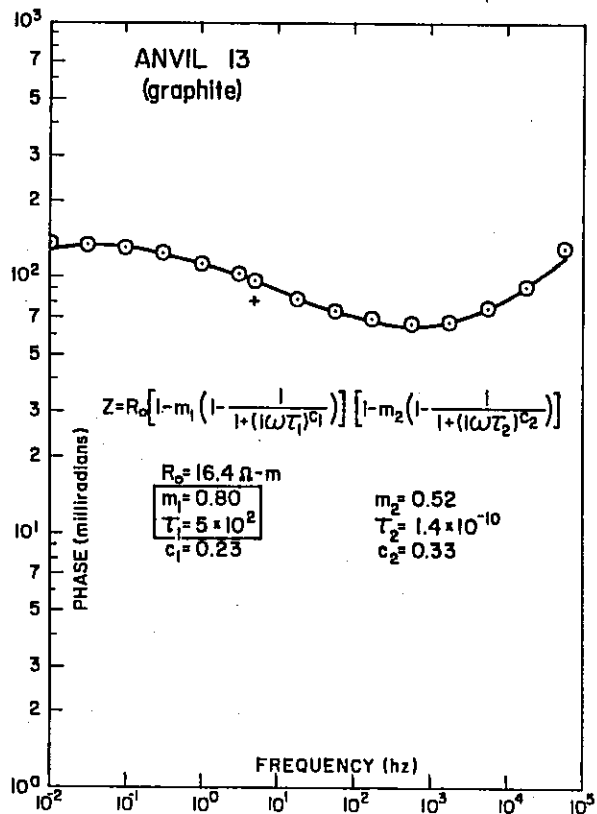


FIG. 4

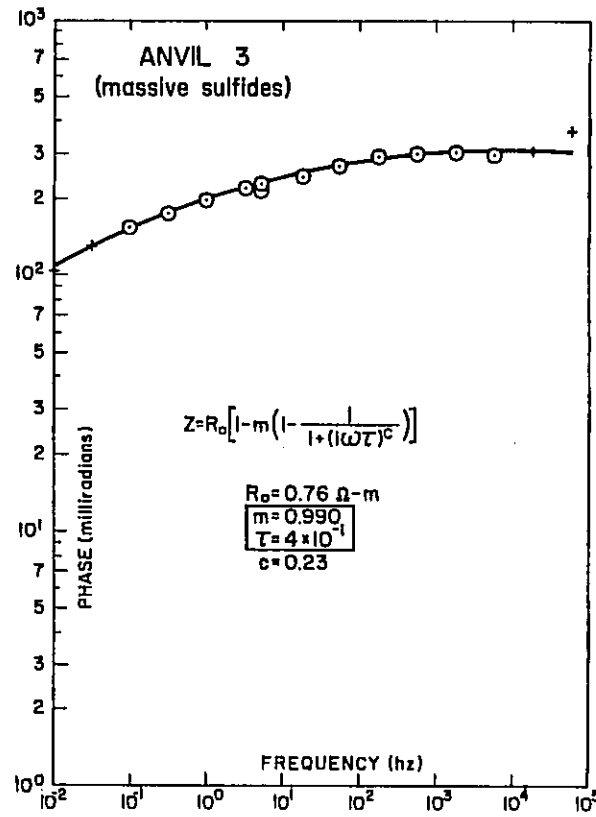


FIG. 5

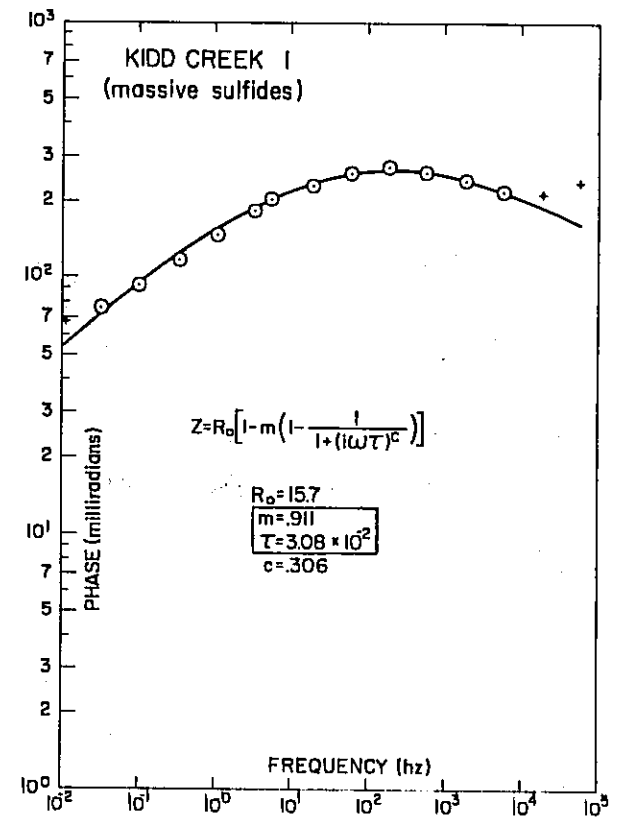
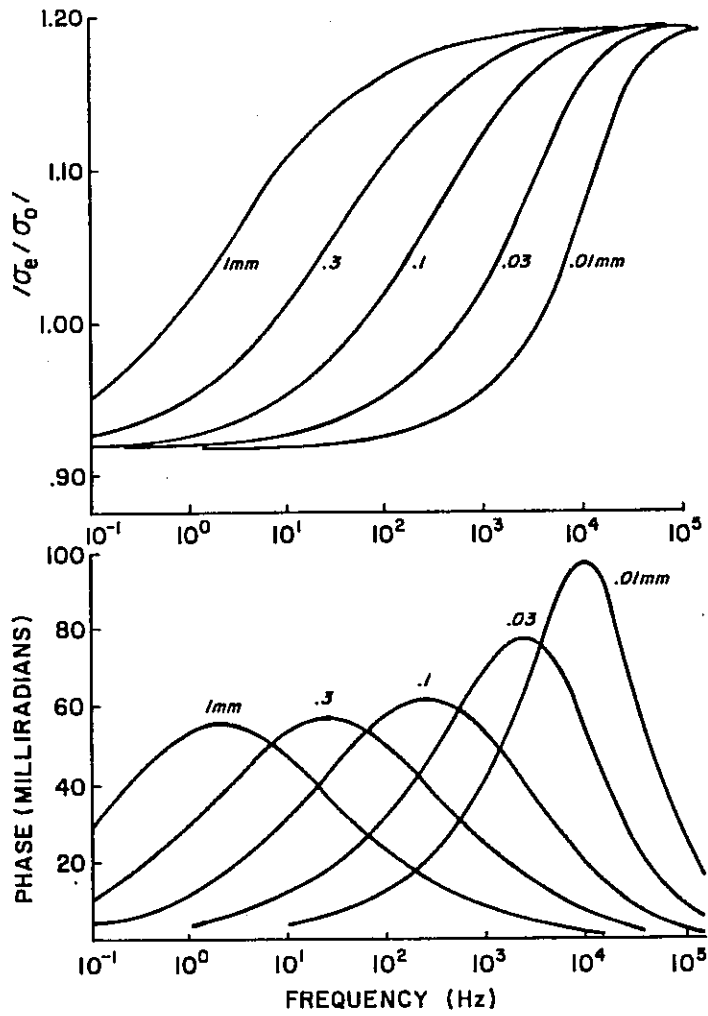


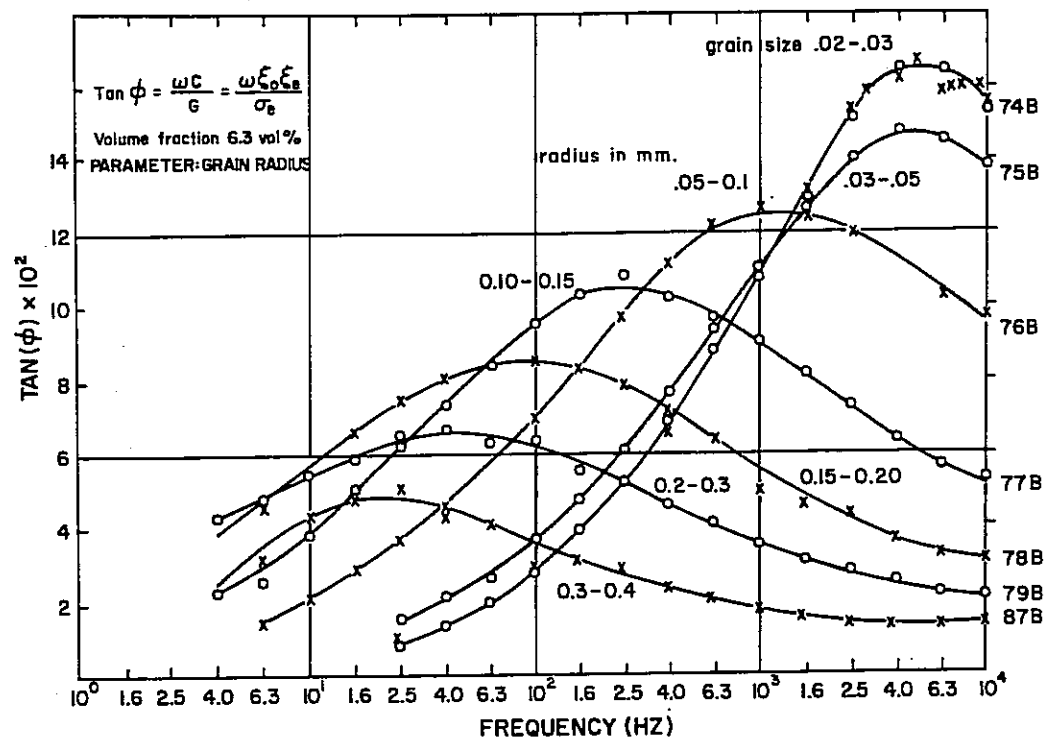
FIG. 6

### THEORETICAL RESULTS FOR VARIATION IN GRAIN SIZE OF POLARIZABLE PARTICLE



from Wong

FIG. 7



from Wong  
data by Grisse mann

FIG. 8

our research has shown. The time-constant ( $\tau$ ) of a spectral IP dispersion curve is directly related to the physical size of the metallic particles that are the source of the IP effect. The critical ( $f_c$ ) is therefore inversely related to the grain-size. Thus, for all of the massive sulphide deposits for which we have spectral measurements, the effective grain-size of the polarizable particles is appreciably smaller than that for graphite sources.

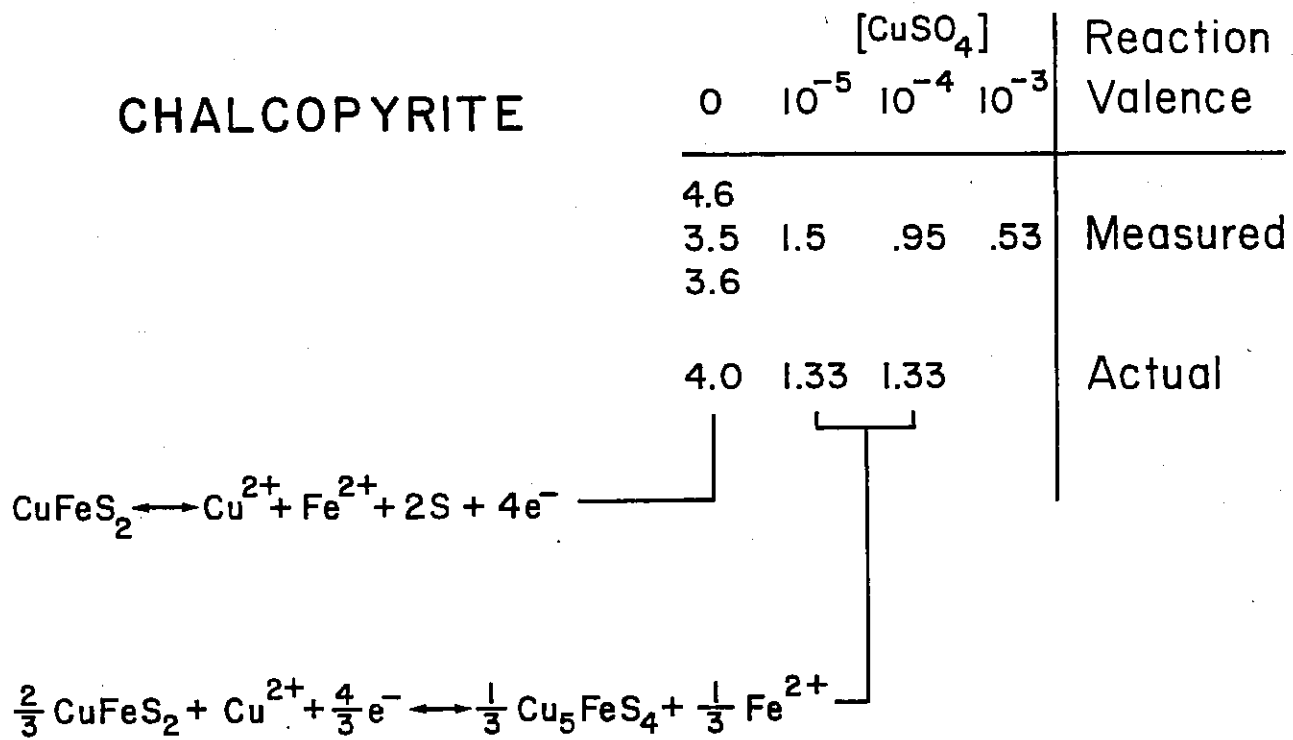
There is one additional test we have devised to confirm the validity of the electrochemical model we have developed to describe the detailed nature of the IP phenomenon. The spectral IP results for a single mineral electrode (pyrite, chalcopyrite, bornite, galena, etc.) immersed in distilled water can be measured, and then inverted using the computer. This inversion determines the various electrical parameters that describe the electrical transfer function for that particular mineral. These electrical parameters, and the analytically determined concentrations of the metal cation in the solution after steady-state current flow has been achieved, can be used to calculate the reaction valence (total number of electrons involved) of the chemical reaction at the interface. This measured value of the reaction valence can be compared to the various values that would be possible from various chemical reactions.

In past research, these measured numbers have not agreed within an order of magnitude with those predicted by the commonly accepted dissolution reaction equations for the various sulphide minerals. The results we have achieved for galena (Figure 10) and chalcopyrite (Figure 9) are typical of the work that has been done recently. The measured values of the reaction valence give much closer agreement to the predicted valence than was previously the case. For chalcopyrite, the predicted reaction varies with the copper ion concentration, and the agreement is still quite acceptable.

With the success of the above described research, has come the implication that we may now understand the IP phenomenon well enough to use some of our knowledge in a predictive way. The resistivity and phase-shift data plotted in the pseudo-section format on Figure 11 are from the Kennedy Property near Winnemucca, Nevada. The results suggest a broad, weak IP anomaly that might be due to a "porphyry copper" type source containing a small concentration of metallic mineralization. However, one notices immediately that the phase-shifts measured at 9.0 Hz are the same magnitude as those measured at 0.11 Hz (a factor of  $3^{-4}$  lower in frequency). This is the same Spectral IP characteristic that was previously measured at the Brenda Mine in British Columbia.

This unusual circumstance is confirmed by the spectral plots for two typical dipole pairs, shown in Figure 12 and Figure 13. With the inductive coupling effects removed by computer inversion, the phase-shift vs. frequency curve is very flat over a frequency range of  $3^8$  Hz. This requires a very low ( $c_1$ ) value of 0.125 (see Figure 3); more importantly, a value of IP effect ( $m$ ) in excess of 0.500 is necessary. Therefore, although the measured phase-shifts are low in magnitude, the true IP effect within the source material must be large. This should not be considered to be a weak anomaly.

(8)



(Wadsworth, 1980)

FIG. 9

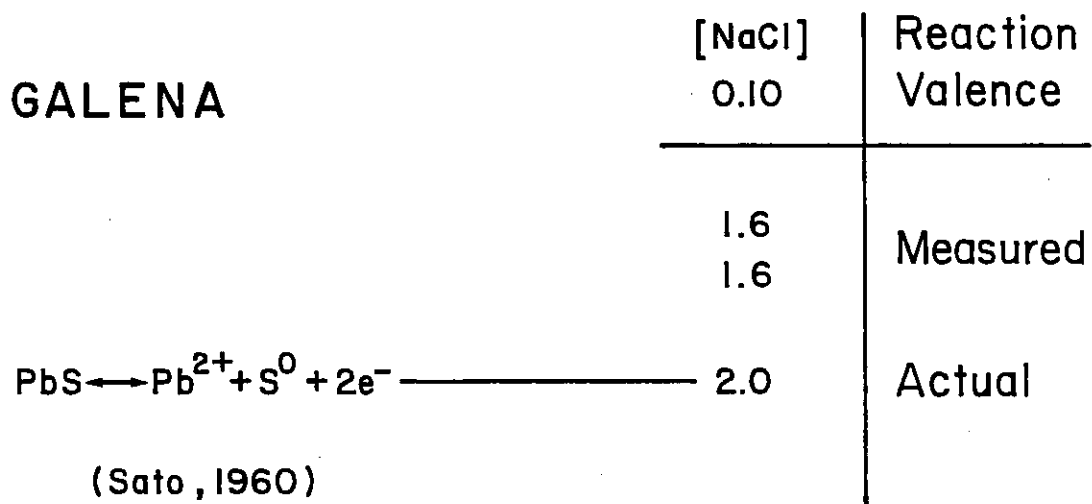
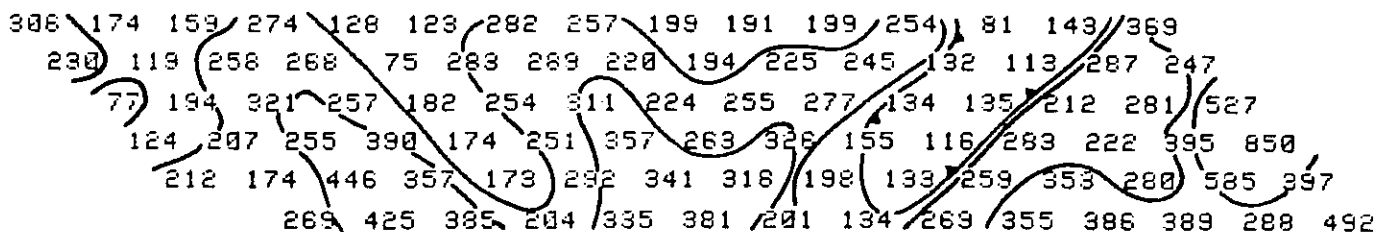


FIG. 10

F 37

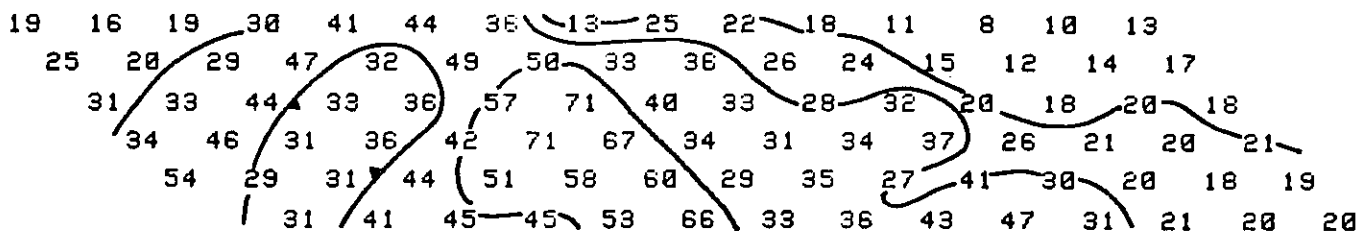
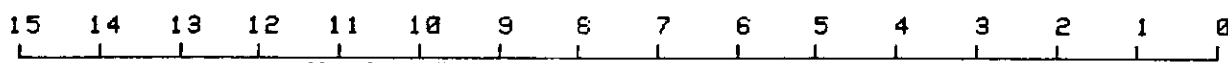
W A BOWES KENNEDY AREA WINNEMUCCA NEVADA X=200 JUNE 1979

Apparent Resistivity(ohm-m)



Decoupled Phase(mrad)

9 Hz



Decoupled Phase(mrad)

.1111 Hz

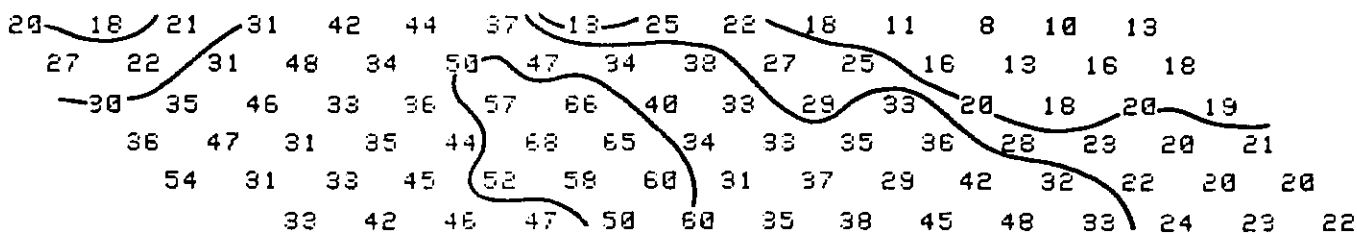
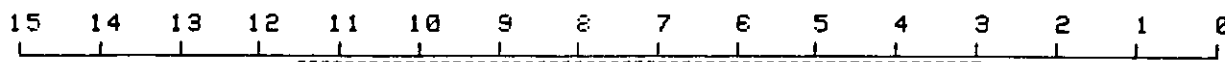
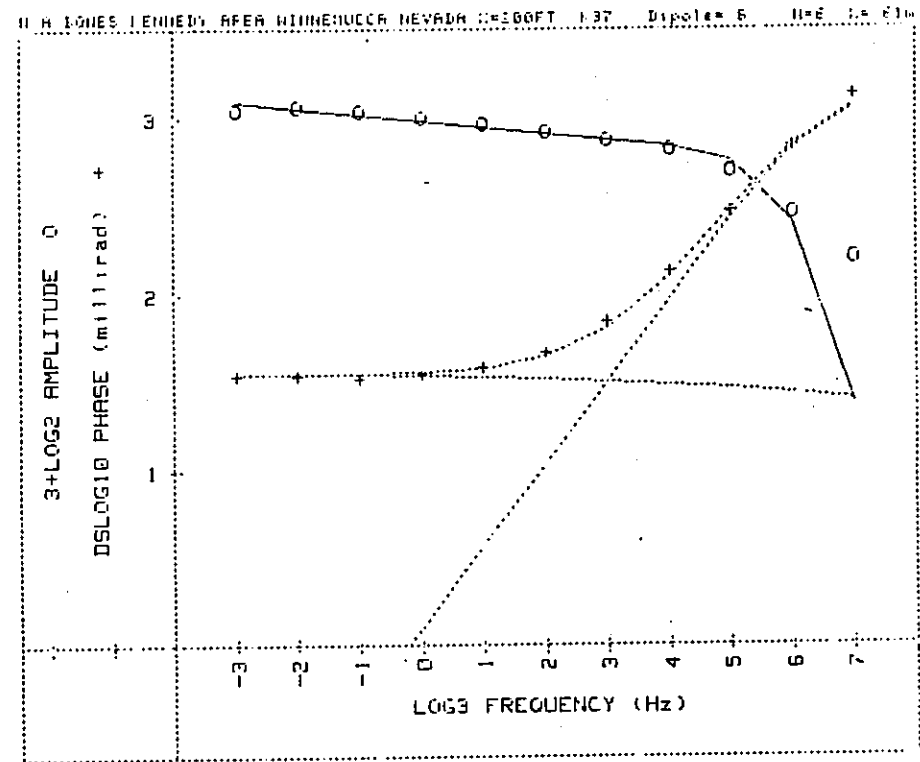
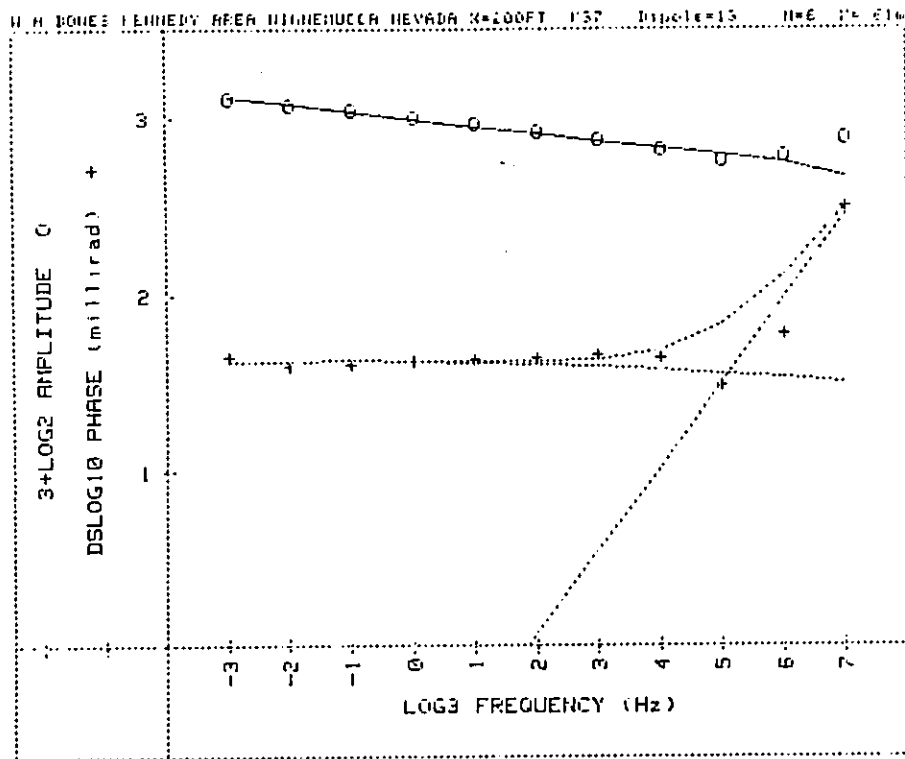


FIG. 11



Iter	Lambda	Rchsq	R0	M1	T1	C1	M2	T2	C2
0	1.E-02	.00068	1.599	.567	2.0E+01	.125	1.000	2.0E-05	.992
1	1.E-01	.00064	1.599	.566	2.0E+01	.125	1.000	2.1E-05	1.000
2	1.E+00	.00063	1.596	.564	2.0E+01	.125	1.000	2.1E-05	1.000
3	1.E+00	.00063	1.595	.563	2.0E+01	.125	1.000	2.1E-05	1.000

Pct Std Deviations 5.4 1.6 499.4 0.0 0.0 10.5 9.8

Correlation Matrix

1.0000									
.5235	1.0000								
.9782	.3667	1.0000							
0.0000	0.0000	0.0000	0.0000						
0.0000	0.0000	0.0000	0.0000	0.0000					
-.3604	.0864	-.3990	0.0000	0.0000	1.0000				
-.6447	.0524	-.7125	0.0000	0.0000	.7305	1.0000			

Apparent Resistivity Measured at 1 Hz is 425

Apparent Resistivity Calculated from Inductive Coupling is 2117

FIG. 12

Iter	Lambda	Rchsq	R0	M1	T1	C1	M2	T2	C2
0	1.E-02	.00102	1.482	.517	2.0E+01	.125	1.000	1.8E-04	.977
1	1.E-02	.00102	1.481	.518	1.9E+01	.125	1.000	1.8E-04	.978

Pct Std Deviations 6.7 2.6 1243.0 0.0 0.0 5.4 2.8

Correlation Matrix

1.0000									
.6234	1.0000								
.9788	.4657	1.0000							
0.0000	0.0000	0.0000	0.0000						
0.0000	0.0000	0.0000	0.0000	0.0000					
.1989	.0726	-.2067	0.0000	0.0000	1.0000				
-.4224	.0478	-.4586	0.0000	0.0000	.4189	1.0000			

Apparent Resistivity Measured at 1 Hz is 201

Apparent Resistivity Calculated from Inductive Coupling is 235

FIG. 13

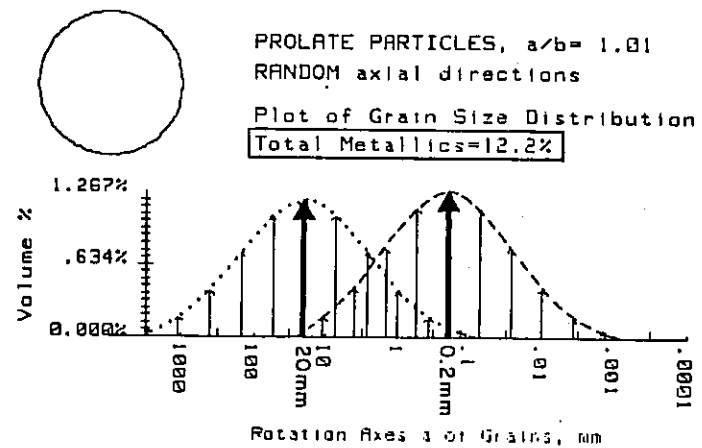
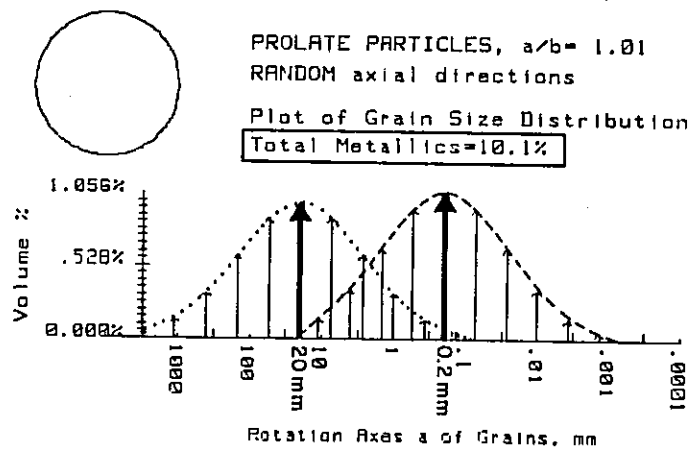
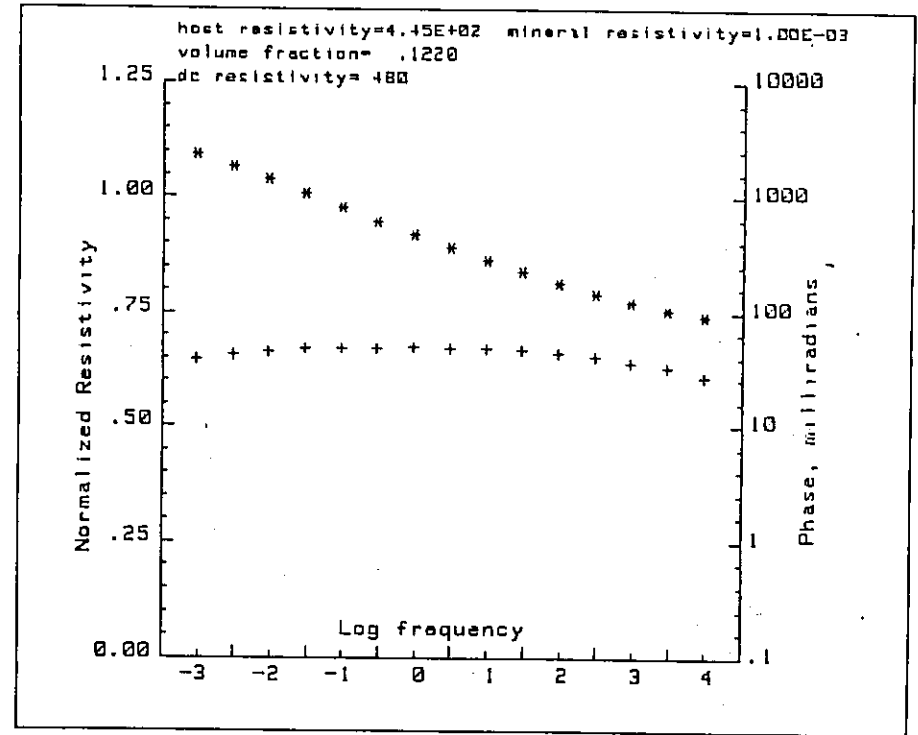
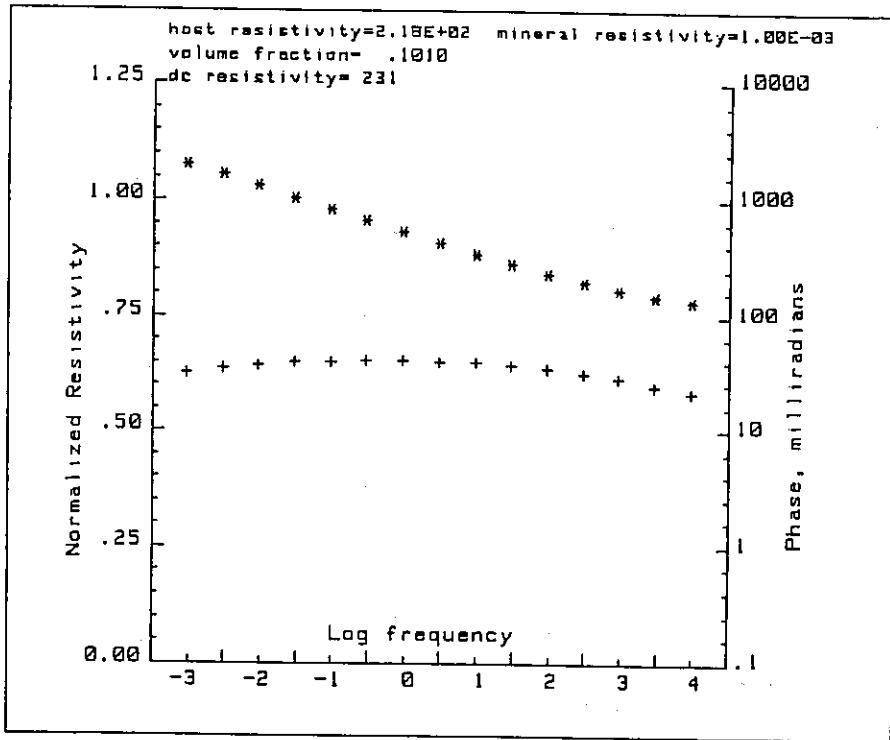


FIG. 14

FIG. 15

We used these results in our first attempt to predict the nature of the mineralization that is the source of the "weak" anomaly shown in Figure 11. The mathematical formula used to predict the spectra shown in Figure 7 has been expanded to include prolate and oblate spheroids, as well as spheres, of metallic material. By a trial and error process, the spectra shown in Figure 14 and Figure 15 were generated as being a "good fit" to the measured data.

The chosen parameters for the sources are probably not unique; however, some valid points can be drawn from the calculation presented.

- i) the flat spectrum (small  $c$  value) cannot be obtained from a single grain-size (time-constant) population. There must be at least two different grain-size populations present.
- ii) Although the two grain-size distributions, with mean values of 20mm and 0.2mm, are probably not exactly correct, a difference of at least two orders of magnitude is necessary to produce the value  $c = 0.125$ .
- iii) Despite the low magnitude for the IP field anomaly, the true IP effect within the source is large ( $m_1 > 0.500$ ) and this results in a fairly high concentration of metallic mineralization predicted for the source (10 - 12% total metallic mineralization).

The first drill hole at the Kennedy Project is now nearing completion; it was collared at approximately station No. 5 along the line of data shown in Figure 11. There was metallic mineralization throughout the drill hole, with an increase below about 75-80 meters of depth.

Throughout the entire length of the drill hole (in excess of 250 meters), there were two different types of mineralization present. There were veinlets (large grain-size) that might, or might not, contain quartz. All of the veinlets contained chalcopyrite or molybdenite or chalcopyrite with molybdenite; some pyrite was also present. There was also fine-grained disseminated metallic magnetite throughout the hole (the small grain-sized source). The magnetite contains appreciable gold value; there is enough gold present to give detectable assays throughout the length of the hole.

We do not have enough information to determine if the total concentration of metallic mineralization approaches 10 to 12 percent. However, in other respects the source of the IP anomaly does seem to approximate that predicted by the spectral IP analysis.

#### IMPROVED MEASUREMENT TECHNIQUES AND INSTRUMENTATION

The much publicized "explosion" in solid-state, semi-conductor

electronic components has produced as many changes in the geophysical instrumentation industry as it has in other phases of modern life. Beginning with diodes and transistors and now including complex microprocessor chips, there has been a bewildering number of improvements in this industry. They have resulted in a major reduction in the size of, and the power requirements of (smaller battery packs), geophysical field instruments; at the same time the reliability, digital data storage capability and analysis capability have been almost infinitely increased.

The Phoenix IPV-2, Prospecting Phase IP system is an example of one of the least complicated units that can be engineered with modern electronics. The types of features that can now be made available are:

- i) Microprocessor-controlled signal stacking of each measurement to give almost infinite noise rejection through coherent detection.
- ii) Matched, heated crystal clocks give phase stability between current source and voltmeter without connecting cable.
- iii) Dual-channel electronics permits two separate voltages to be measured simultaneously.
- iv) Digital readout of magnitude and phase-shift of measured voltage at either channel.
- v) Five frequency capability (0.11, 0.33, 1.0, 3.0, 9.0 Hz) permits limited analysis of spectral character of any anomaly.
- vi) Appreciable increase in survey speed over either frequency-domain or time-domain IP systems.

By completing a reconnaissance survey with a single frequency (typically 0.33 Hz or 1.0 Hz) it is possible to achieve the lowest cost possible for IP field surveys. However, once an anomaly is detected, the IPV-2 can be used to make detailed measurements that give some feeling for the character of the spectral response of the source.

If detailed measurements are made at only two frequencies, it is possible to make some statement about the critical frequency ( $F_c$ ) of the IP spectrum from the source. On Figure 17 and Figure 18 are shown the detailed phase-IP measurements made using  $X = 100'$  and 0.33 Hz over two anomalies previously located by a reconnaissance IP survey using larger electrode intervals. On both Line 18E and Line 15E the anomalous pattern indicates a relatively narrow source, with some depth to the top.

During the detailed survey with  $X = 100'$ , several of the anomalous values in each anomalous pattern were measured at both 0.33 Hz and 3.0 Hz.

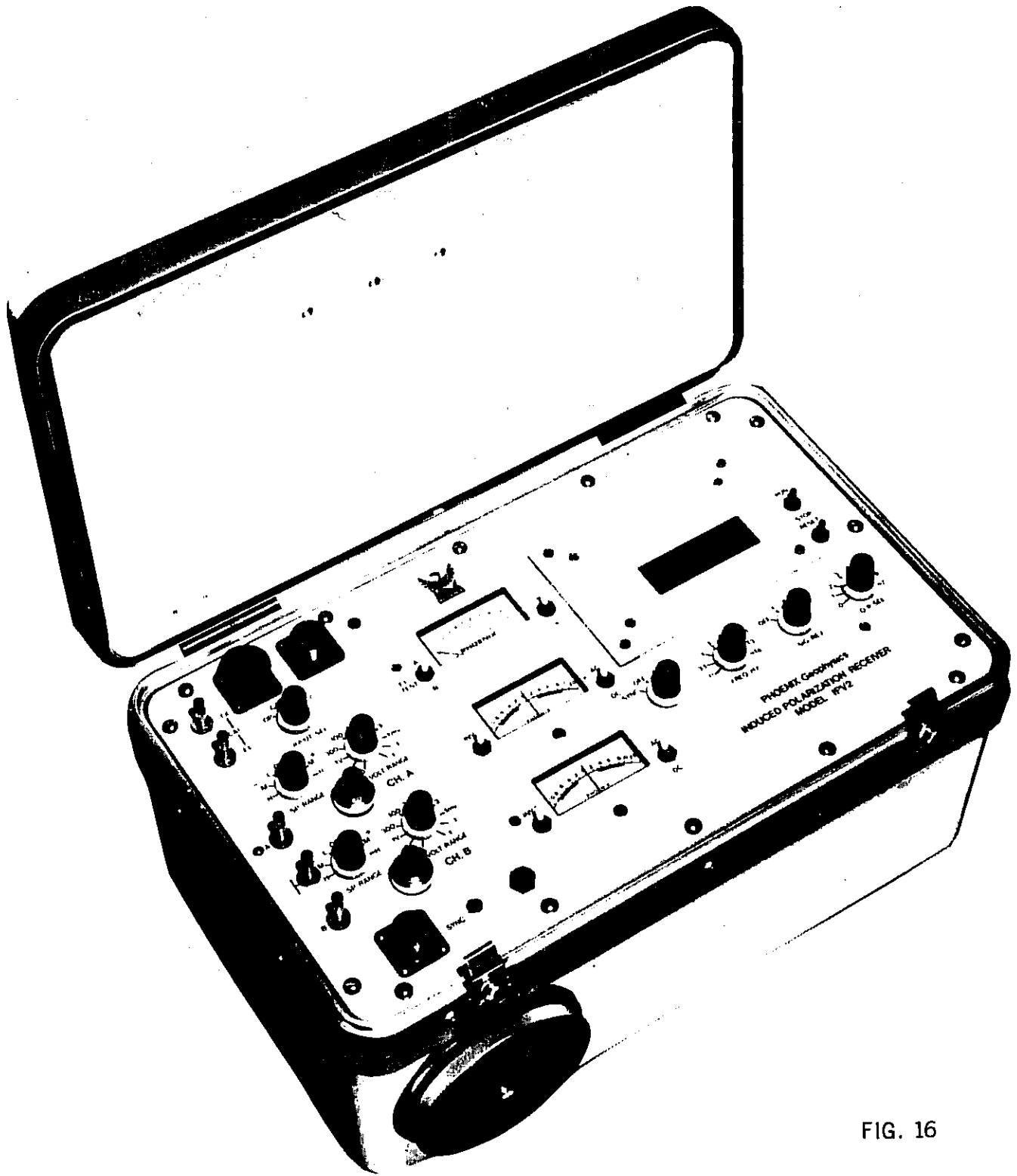


FIG. 16

PRICE COMPANY BLOCK NO.8 PROPERTY  
IPV-2 PHASE IP TEST

LINE-18E

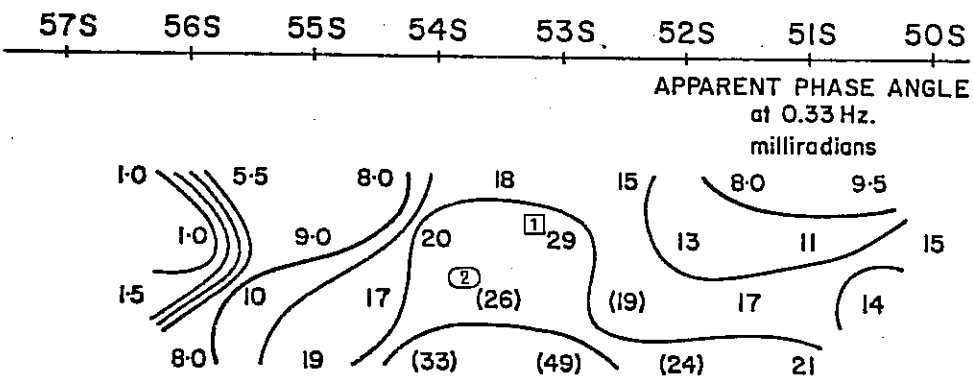
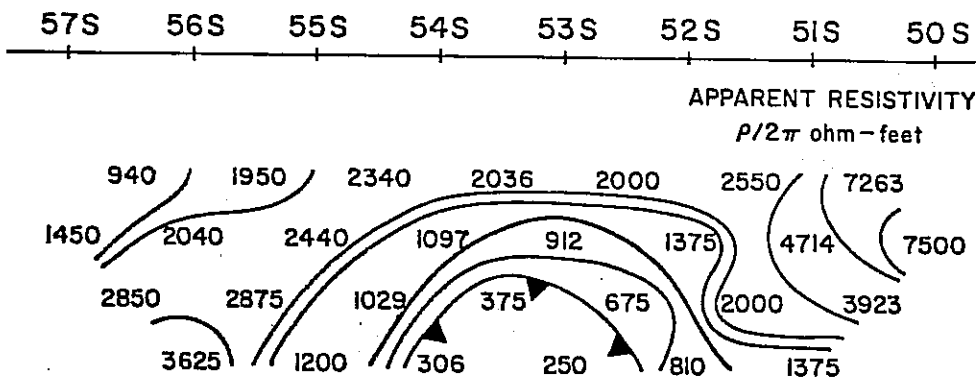


FIG. 17

PRICE COMPANY BLOCK NO.8 PROPERTY  
IPV-2 PHASE IP TEST

LINE-15E

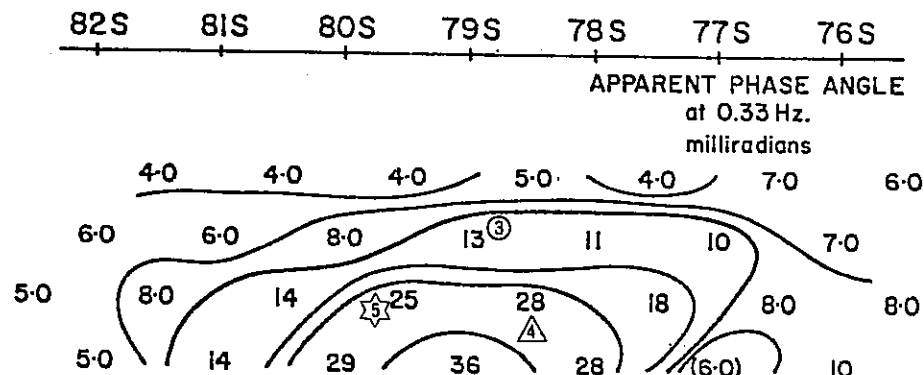
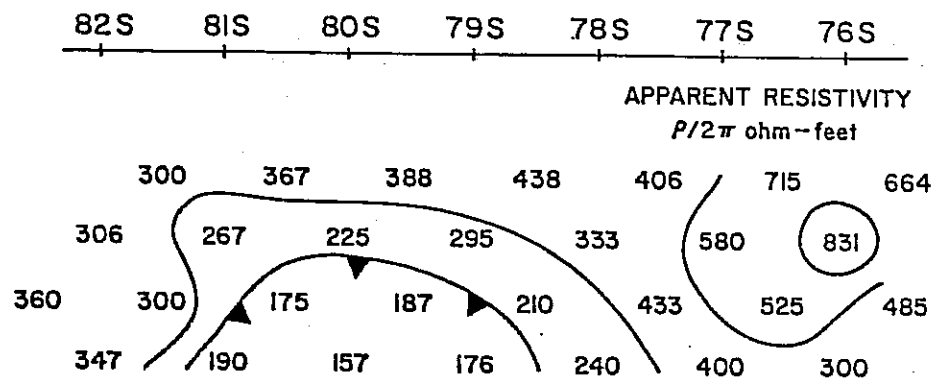


FIG. 18

The results are shown on Figure 19 and Figure 20. On Line 18E, the apparent phase values measured increase with frequency, indicating a value for  $F_c$  above 3.0 Hz. On Line 15E the apparent phase values were decreasing with frequency, indicating a value for  $F_c$  that is below 0.33 Hz.

The relationship we have established between the critical frequency ( $F_c$ ) - the time constant ( $\tau$ ) - and the grain-size of the polarizable particles, coupled with our experience over known zones suggests that the grain-size of the source of the anomaly on Line 18E is appreciably smaller than the grain-size of the source of the anomaly on Line 15E.

Our expectation that the large grain-size source is more likely to be graphite is confirmed by the drilling of the two anomalies. The anomaly on Line 18E was found to be due to massive sulphide mineralization; the source of the anomaly located on Line 15E was a graphitic tuff zone, with small concentrations of pyrite.

If field time is taken to make measurements at all five frequencies, it is possible to make a more formal presentation of the spectral characteristics of two anomalies. On Figure 21, is shown a computer-plotted pseudo-section of a single line of phase IP data. The plot shows the apparent resistivity value and the apparent phase value (1.0 Hz) measured for each dipole pair. At the lower part of the data plot is shown the short (five frequency) spectrum (normalized to the value at 1.0 Hz) for each dipole pair.

In this case, the two anomalous zones had a considerable strike length, in an area where the surface was largely covered. The electrical parameters of the two sources were much the same; however, as shown on Figure 21, the source of Zone I is clearly different from the source of Zone II. From drilling, it is known that Zone I is due to a broad, irregular zone of pyrite and sphalerite (with appreciable silver) in a tuffaceous rock unit. The source of Zone II is pyrrhotite in an adjoining andesite rock unit.

Another example of a difference in spectral character is shown in Figure 22. In this location, the source of Zone A had been traced for a considerable distance along strike. Zone B was detected only on two adjacent lines. The difference in spectral character strongly suggests that Zone B is not merely a greater width to the mineralized zone that is the source of Zone A.

The Phoenix IPV-3, Spectral IP Receiver is the most advanced induced polarization receiver ever put into field operation (Figure 23). It completely replaces a truck-mounted system that was assembled largely from off-the-shelf computers, voltmeters, oscilloscopes, etc. The entire system is controlled by the Phoenix CP-1, geophysical computer, that was designed especially for this function. Through an FM-modulated link, the CP-1 also controls the spectral IP current source. Some of the more important features of the IPV-3 are:

- i) Six simultaneous voltage channels plus current monitor.

LINE-18E  $R\phi_{0.3-3.0}$  TEST

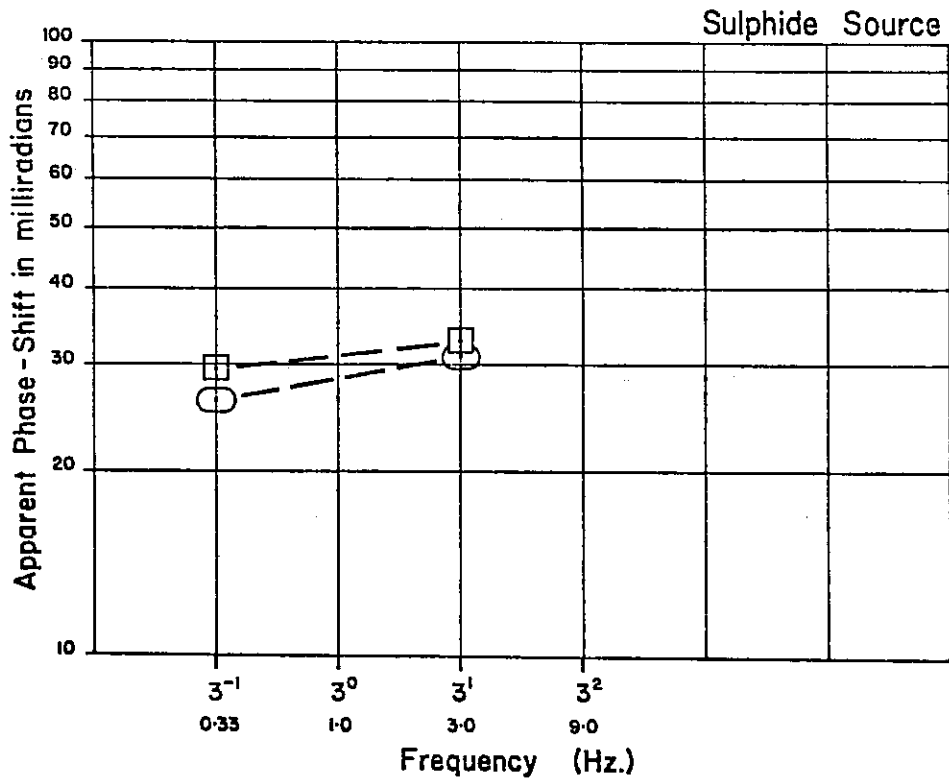


FIG. 19

LINE-15E  $R\phi_{0.3-3.0}$  TEST

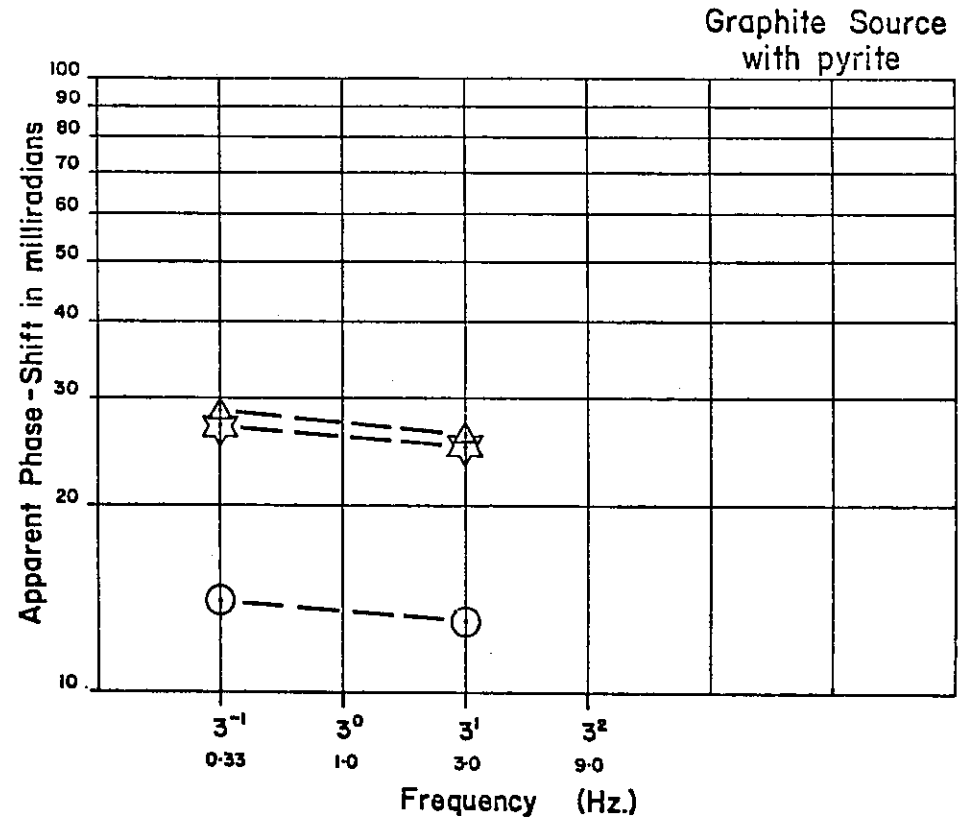


FIG. 20

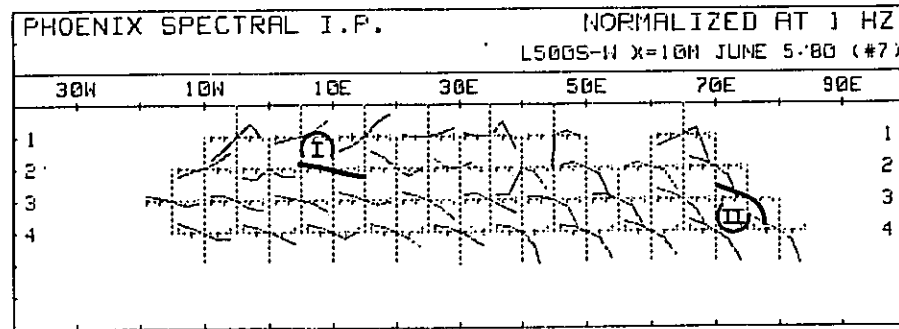
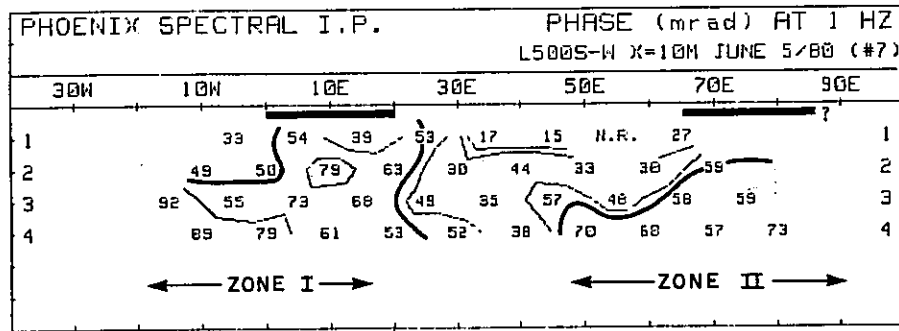
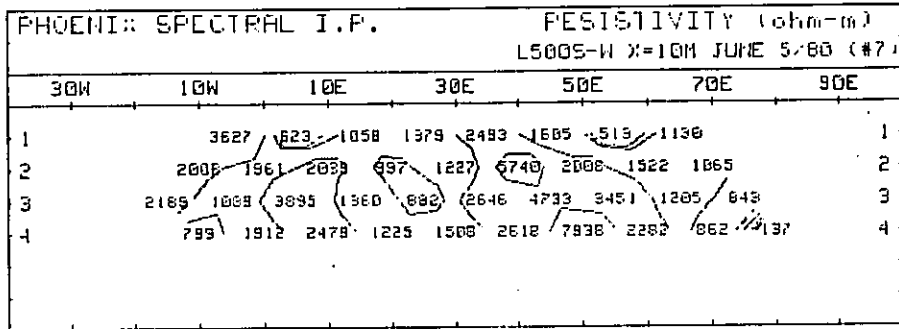


FIG. 21

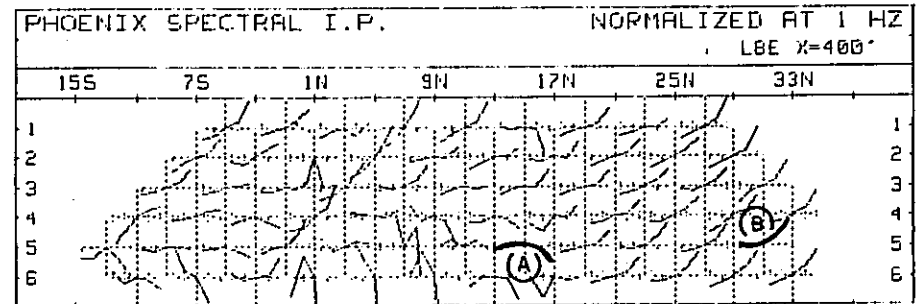
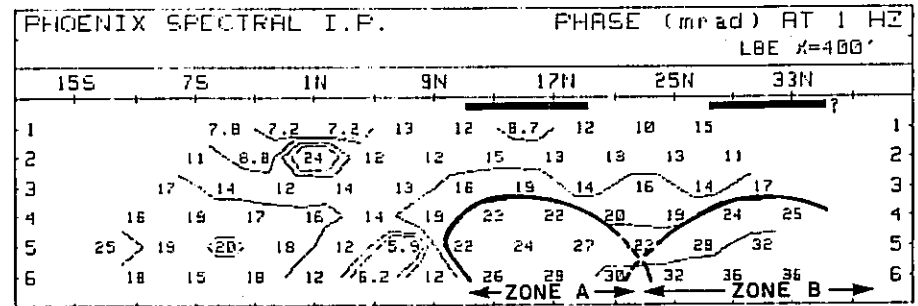
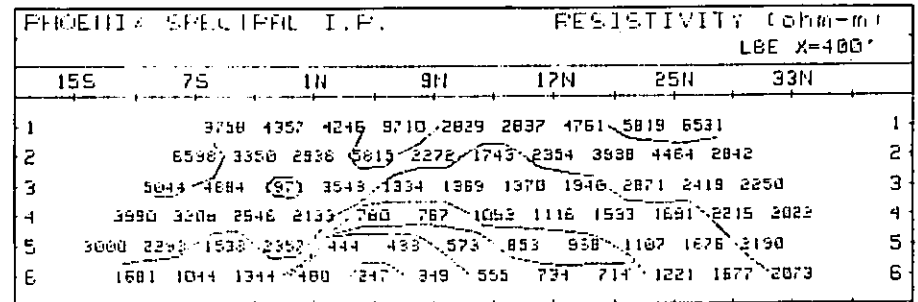


FIG. 22

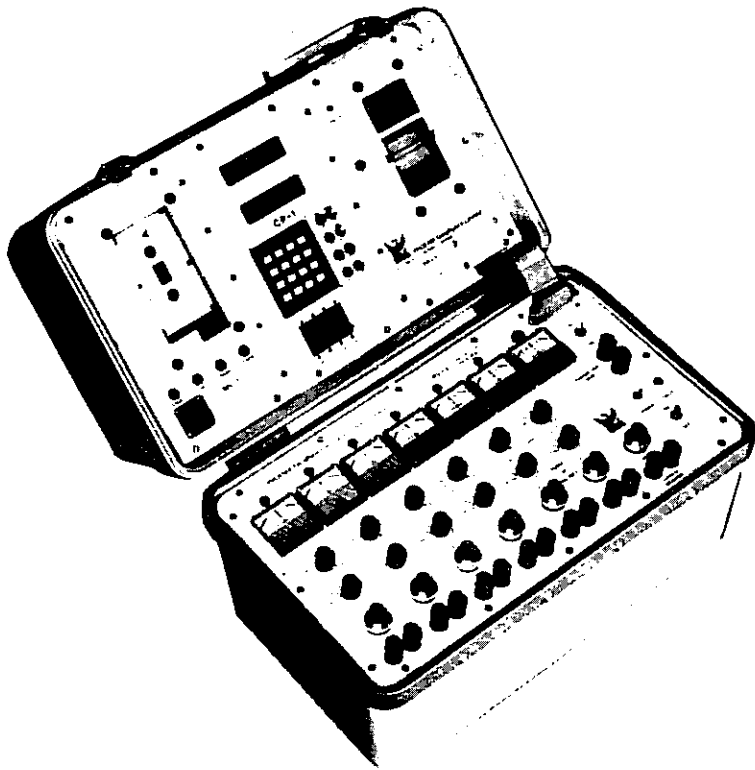


FIG. 23

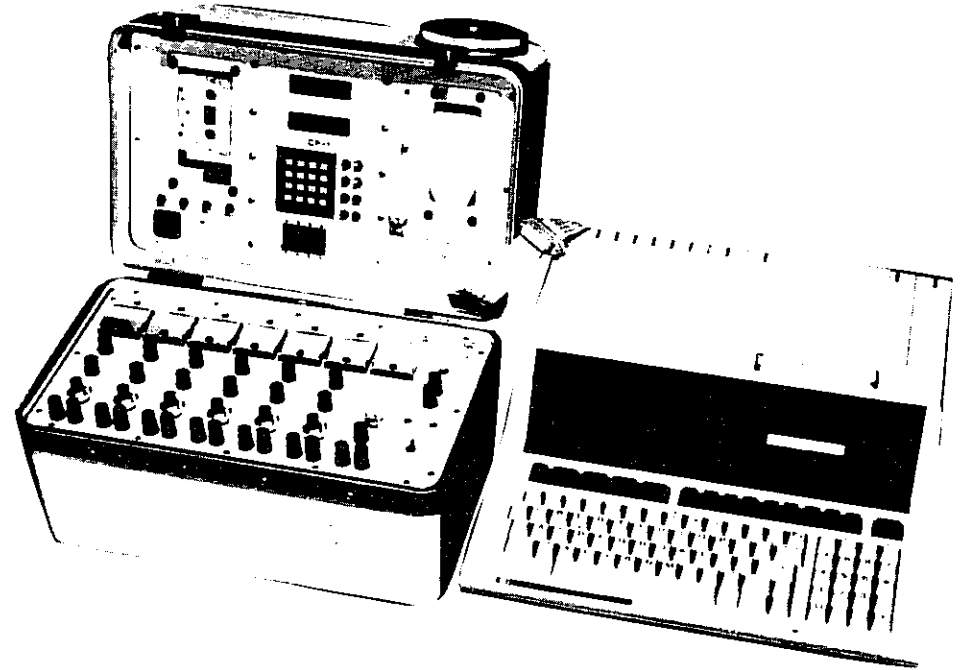


FIG. 24

- ii) Frequency range from  $2^{-14}$  Hz (0.00006 Hz) to  $2^{+16}$  Hz (65,536 Hz) with automatic scanning programmable through CP-1.
- iii) Real time de-convolution and coherent filtering via use of CP-1.
- iv) Field Data available through digital read-out and recorded on paper tape and tape cassette.
- v) Field results directly transferrable from tape cassette through CP-1 and interface to computer for analysis and interpretation. (Figure 24).

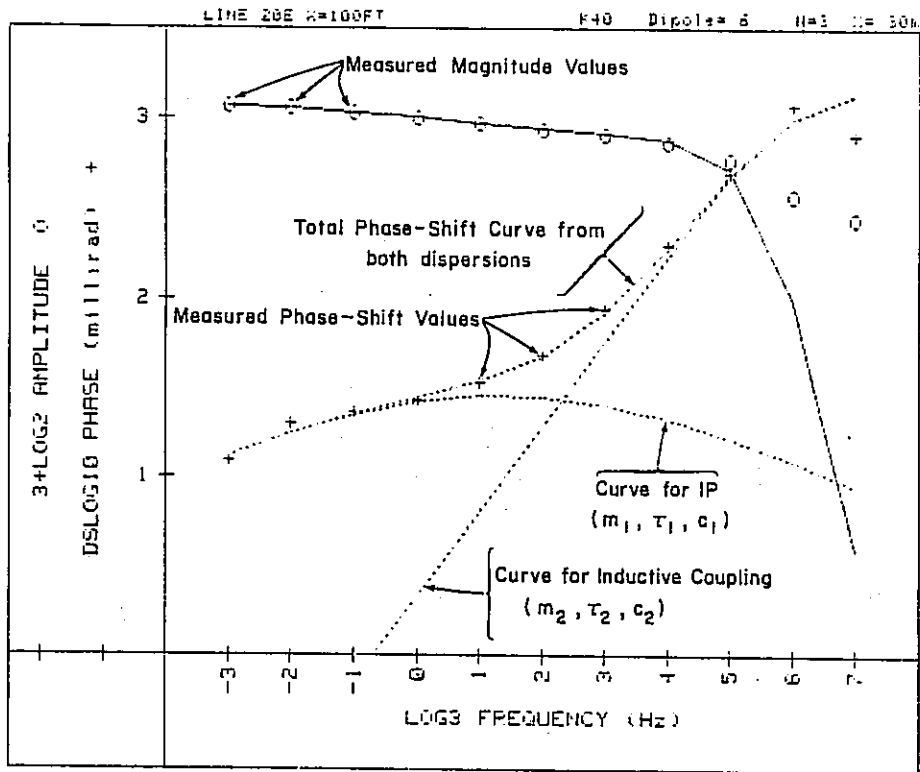
By using the IPV-3 and the HP-85 computer in the field, the geophysicist has at his disposal all of the interpretational power of the spectral IP method. As shown on Figure 25 it is possible to remove the inductive coupling effects from the broad-band spectral IP results. When this information is removed, it is possible to invert the remaining, decoupled IP spectra and determine the apparent values of  $R_0$ ,  $m$ ,  $\tau$ ,  $c$  for that particular dipole-pair measurement.

For the spectra shown in Figure 25, the values are  $m = 0.172$ ;  $c = 0.376$ ;  $\tau = 0.044$  sec.; these parameters suggest a small grain-size, and in fact the source is a massive sulphide zone in New Brunswick. The parameter values obtained by inverting the spectral curve from all of the anomalous dipole pairs were essentially the same; we have found this to be universally true if the anomaly is due to a single, more-or-less uniform source.

For the spectral curve shown in Figure 26, the parameter values of the source are  $m = 0.509$ ;  $c = 0.180$ ;  $\tau = 540$  seconds. The large time-constant confirms the large grain-size for the graphite zone that is the source of the anomaly.

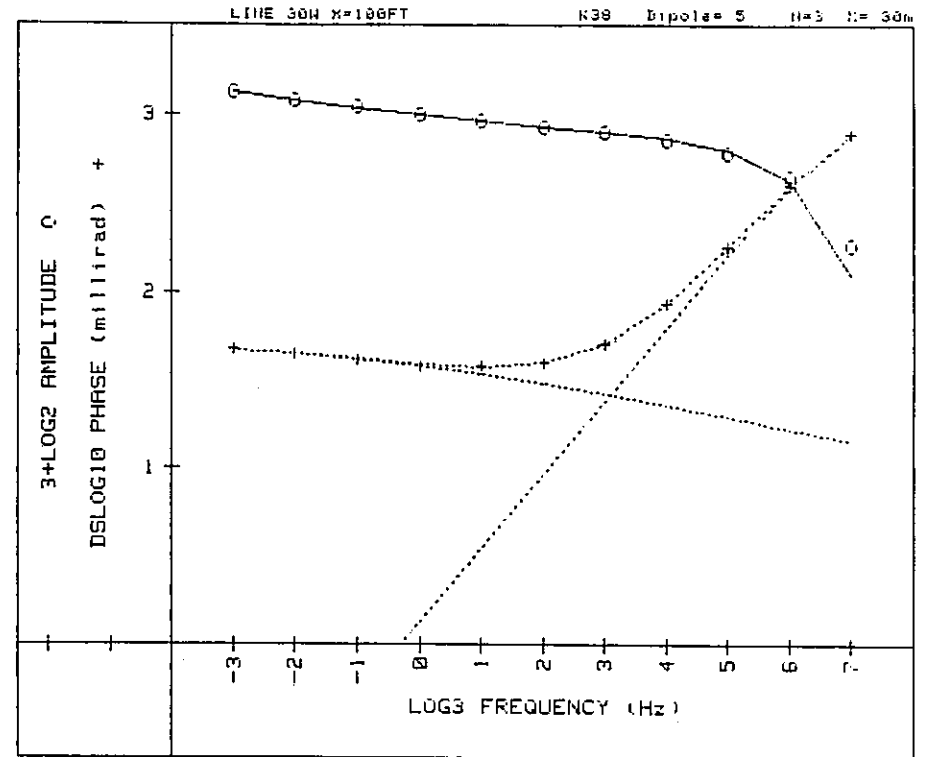
For IP measurements with large electrode intervals in areas of low resistivity (high conductivity) inductive coupling effects will completely distort IP field data. This can happen even for low frequency measurements (in the variable frequency method) or long time delay measurements (in the time domain technique). When the inductive coupling effects are removed from the spectral curve for each dipole pair (as in Figure 12 and Figure 25) the pseudo-section can be reconstructed using the decoupled data.

The data pseudo-section shown in Figure 27 is from the Elura Deposit near Cobar, New South Wales. The resistivity results show a deep, variable high conductivity surface layer. The phase IP measurements at 3.0 Hz with  $X = 100$  meters are completely dominated by inductive coupling effects that essentially depend only upon the electrode separation ( $n$ ). However, when the inductive effects are removed from the data for each dipole pair, the pseudo-section clearly shows the presence of the massive sulphide zone, at depth, at approximately station 2500E.



Iter	Lambda	Rchsq	F0	IP EFFECT			INDUCTIVE COUPLING EFFECT			
				M1	T1	C1	M2	T2	C2	
0	1.E-02	.05406	1.063	.221	4.0E-02	.323	1.000	1.1E-04	1.000	
1	1.E-02	.00168	1.063	.196	2.5E-02	.373	1.000	3.6E-04	1.000	
2	1.E-03	.00053	1.072	.186	3.1E-02	.355	1.000	3.3E-04	1.000	
3	1.E-04	.00067	1.071	.174	4.3E-02	.373	1.000	3.4E-04	1.000	
4	1.E-05	.00066	1.071	.172	4.4E-02	.376	1.000	3.4E-04	1.000	
Pct Std Deviations				.3	16.4	84.2	13.5	0.0	5.5	0.0
Correlation Matrix				1.0000						
				.5173	1.0000					
				-.4239	-.9715	1.0000				
				-.5302	-.9438	.3971	1.0000			
				0.0000	0.0000	0.0000	0.0000	0.0000		
				-.2093	-.6536	.6657	.5343	0.0000	1.0000	
				0.0000	0.0000	0.0000	0.0000	0.0000	0.0000	
Apparent Resistivity Measured at 1 Hz is				78						
Apparent Resistivity Calculated from Inductive Coupling is				10						

FIG. 25



Iter	Lambda	Rchsq	R0	M1	T1	C1	M2	T2	C2	
										0
1	1.E-02	.00033	1.569	.483	1.3E+02	.180	1.000	9.3E-05	.893	
2	1.E-03	.00009	1.515	.487	1.9E+02	.180	1.000	8.6E-05	.891	
3	1.E-04	.00006	1.577	.502	4.0E+02	.180	1.000	8.7E-05	.882	
4	1.E-05	.00005	1.704	.508	5.3E+02	.180	1.000	8.7E-05	.879	
5	1.E-06	.00005	1.707	.509	5.4E+02	.180	1.000	8.7E-05	.879	
Pct Std Deviations				2.4	2.0	53.1	0.0	0.0	2.0	1.0
Correlation Matrix				1.0000						
				.9916	1.0000					
				.5085	.9710	1.0000				
				0.0000	0.0000	0.0000	0.0000			
				0.0000	0.0000	0.0000	0.0000	0.0000		
				-.1166	-.0989	-.1282	0.0000	0.0000	1.0000	
				-.4909	-.4268	-.5273	0.0000	0.0000	.7128	1.0000
Apparent Resistivity Measured at 1 Hz is				231						
Apparent Resistivity Calculated from Inductive Coupling is				26						

FIG. 26

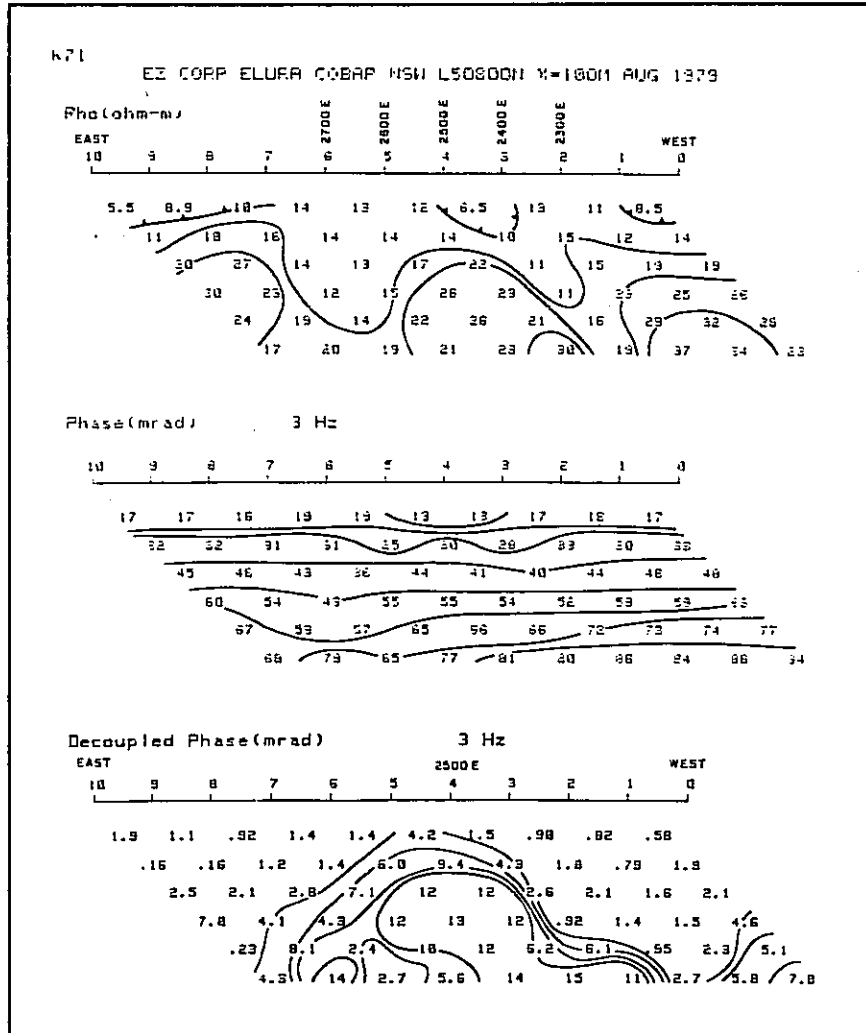


FIG. 27

COMPUTER ASSISTED INTERPRETATION OF IP DATA

The increased size, storage capability and speed of the modern digital computer has been a big boon to the exploration geophysicist. Forward problem solutions for potential methods (gravity, magnetics, etc.) have been available for some time. Recently, numerical modelling methods and large, very fast computers have made similar forward problem solutions available for applied field geophysics such as resistivity and induced polarization.

The forward problem solution shown in Figure 28 is for a fairly normal, dipping source. (The geometry is Two-Dimensional). The patterns of the induced polarization and resistivity pseudo-sections are similar to those we have previously seen from field data over dipping sources.



Recently however, we had occasion to try to interpret a field anomaly that was very similar to that shown in Figure 29. The presence of the shallow, limited depth extent source was known from previous drilling. By the trial and error method, we were able to show that the pattern (which is very similar to that in Figure 28) need not have originated from a dipping source.

By limiting the solutions somewhat, it is possible to create computer programs to do the inverse problem. That is, the computer accepts the field data as input and as output gives the vertical, tabular source that gives the best "least square fit" to that data. For comparison, the computer also outputs the pseudo-section data that would be measured for the source it has chosen.

The simplest problem is for resistivity data alone. The results shown in Figure 30 are from a geothermal exploration program in which the resistivity survey was completed using  $X = 250$  metres. The large, deep resistivity low was measured on the side of an extinct volcano, under recent lava flows. The computer inversion gives a major, poorly determined, conductor (with a high conductivity times width product) at a depth of 350 to 400 metres. This data was used to show the geologist that the source of the resistivity anomaly could be a good geothermal target.

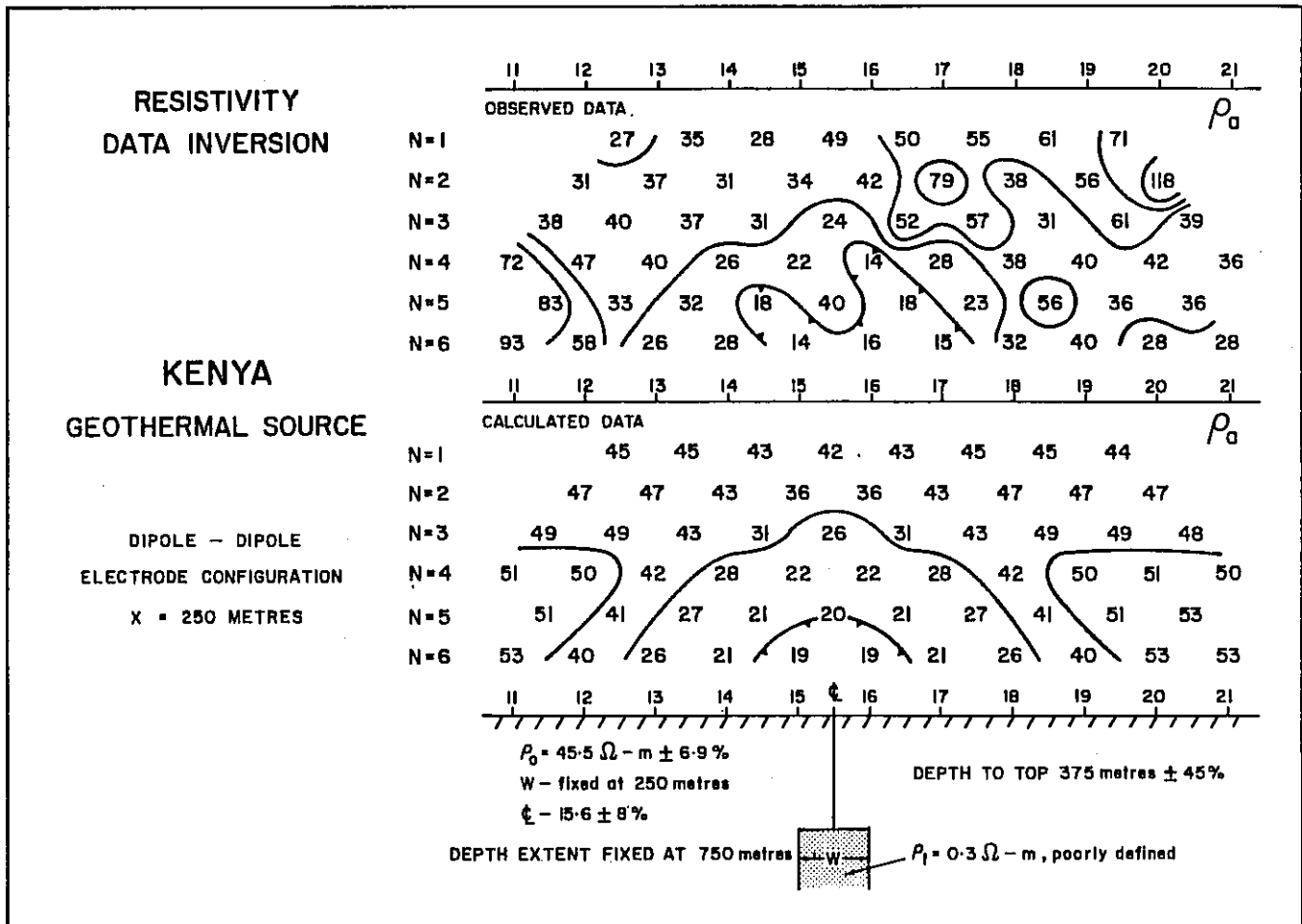


FIG. 30

The example shown in Figure 31 and Figure 32 is from the Pine Point Area. The computer inversion shows that in this case the source has a poor resistivity contrast with the surrounding limestones. However, the IP effect is quite anomalous.

If there is conductive overburden present, and the source itself is conductive, the inversions will be more complex. The data shown on Figure 33 and Figure 34 were measured with  $X = 100$  feet, in an area where the depth of overburden was determined to be about fifty feet. The calculated results give a "good fit" to the measured data, and the source described by this inversion is a typical Canadian Shield, massive sulphide deposit.

The inversion program can also be used to show that two dissimilar anomalies can in fact be due to the same type of source. The inversions of the IP data shown in Figure 35 and Figure 36 are those for the field data previously shown in Figure 17 and Figure 18. The background resistivities are very different, and the depths to the top of the sources are different by a factor of three. However, the indicated source parameters are much the same. As we saw previously, the spectral characteristics of the IP anomaly from the two sources are quite different.

The use of the computer inversion techniques make it possible to obtain an estimate of the source parameters, even in situations in which an unusual geologic environment would make it difficult to interpret the data. The field results shown in Figure 37 are from a deeply buried source in a highly conductive environment. The rocks in which the survey was completed were porous, and filled with saline water. However, the inversion output on Figure 38 shows the source is not unlike that determined in Figure 34.

#### WHAT OF THE FUTURE?

There seems to be almost no limit to the future development of the induced polarization method. As the orebodies for which we search become deeper, and more difficult to detect, we will need all of the help we can get.

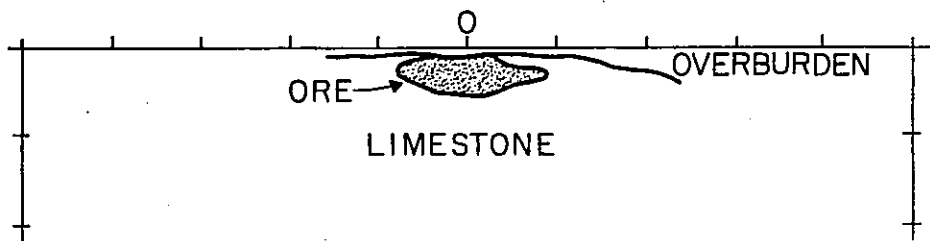
A simple extension of what we are now doing should lead to the following examples in the next twelve to twenty-four months.

- a) The computer can be used to store parameters obtained by the computer inversions of spectral IP data. Therefore, it will be possible to correlate the electrical parameters ( $m$ ,  $\tau$ ,  $c$ ) of all sources known in a given geographical area or a certain geologic environment. For a new, unknown source, the computer could be asked to output the probability that the unknown source will be due to metallic mineralization of a certain type.

The data shown in Figure 39 was a first attempt at such a correlation. On the plot are shown the IP effect vs. Time-Constant ( $m$ , vs.  $\tau$ ) coordinates for a few massive sulphide

### PINE POINT, N-38A, X=200FT

568 310 660 350 324 980 1290 820  
 500 600 400 230 511 380 724 565 530  
 427 550 437 296 390 462 380 329 431 NR



### RESISTIVITY MODEL

539 548 538 319 333 497 552 539  
 541 553 540 362 537 396 480 560 541  
 535 549 544 371 554 554 416 435 554 535

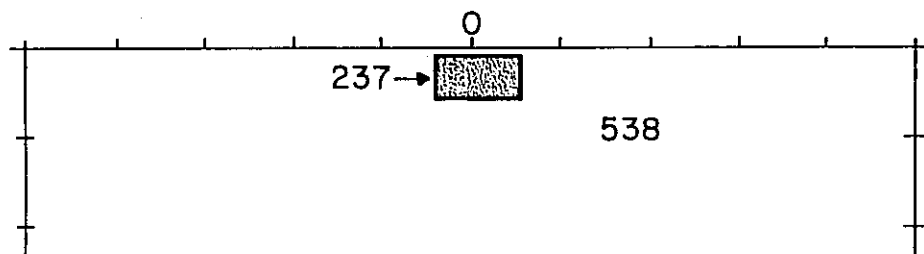
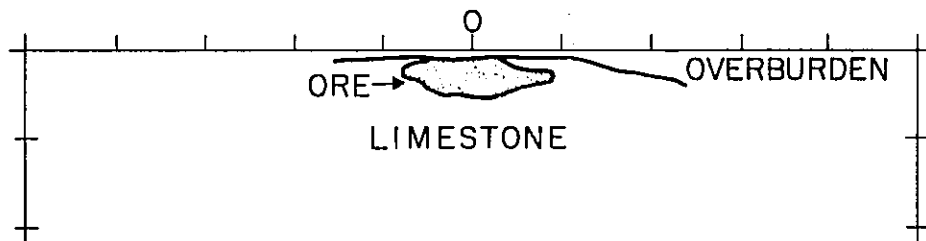


FIG. 31

### PINE POINT, N-38A, X=200 FT

1.6 1.5 1.1 12.1 8.8 2.7 1.8 2.6  
 1.7 1.1 1.8 8.1 2.0 5.3 2.7 2.8 3.8  
 1.2 1.2 1.4 6.8 0.8 0.8 3.5 4.6 1.4 NR



### IP MODEL

1.8 1.6 1.5 10.4 9.2 3.6 1.4 1.8  
 1.8 1.4 1.5 8.6 1.7 6.5 3.9 1.2 1.7  
 1.9 1.5 1.7 8.0 1.4 1.3 5.5 4.3 1.2 1.9

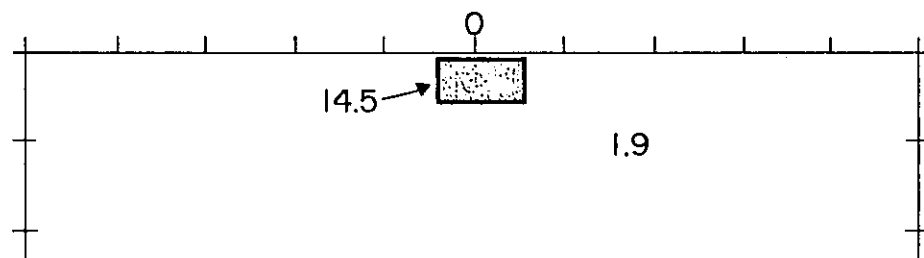


FIG. 32

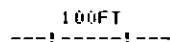
COMPUTER INTERPRETATION  
OF GEOPHYSICAL DATA

INDUCED POLARIZATION AND RESISTIVITY

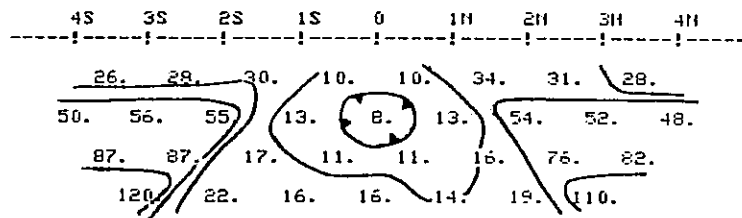
CANADIAN SHIELD MASSIVE SULPHIDE DEPOSIT, OVERLAIN BY CONDUCTIVE OVERBURDEN  
ISD - COPPERFIELDS LINE 0+00

DETERMINE: HORIZONTAL POSITION  
DEPTH TO TOP  
WIDTH  
RES OF SOURCE  
RES OF OVERBURDEN  
IP OF SOURCE

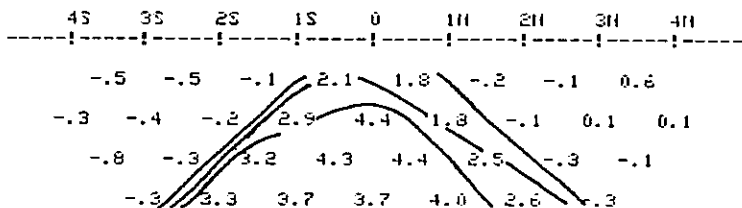
DIPOLE-DIPOLE ARRAY X = 100 FEET



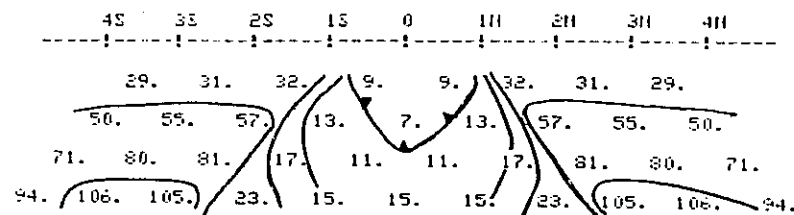
RESISTIVITY DATA - OBSERVED ( OHM-METERS )



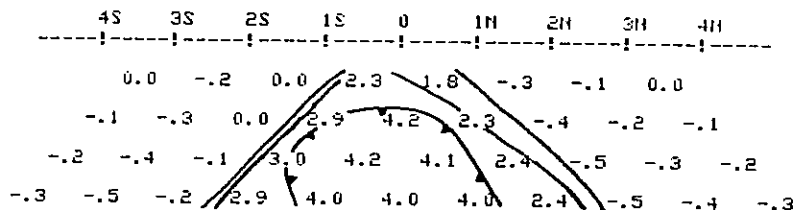
INDUCED POLARIZATION DATA - OBSERVED ( PFE )



RESISTIVITY DATA - CALCULATED ( OHM-METERS )



INDUCED POLARIZATION DATA - CALCULATED ( PFE )



PARAMETERS OF THE CALCULATED SOURCE

	4S	3S	2S	1S	0	1N	2N	3N	4N
CENTER AT STN 0+13S									
DEPTH OF OVERBURDEN									
WIDTH									
RES OF SOURCE									
RES OF OVERBURDEN									
IP OF SOURCE									
DEPTH EXTENT GRTR THAN 200FT									

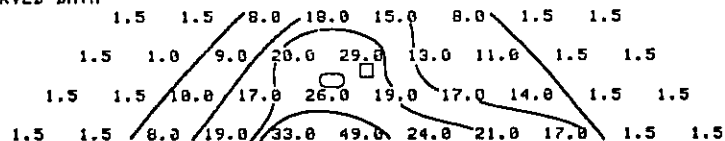
FIG. 33

FIG. 34

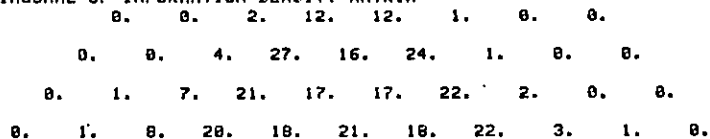
### LINE-18E

FINAL MODEL FOR DATA SET 1

OBSERVED DATA



100\*DIAGONAL OF INFORMATION DENSITY MATRIX



CALCULATED DATA

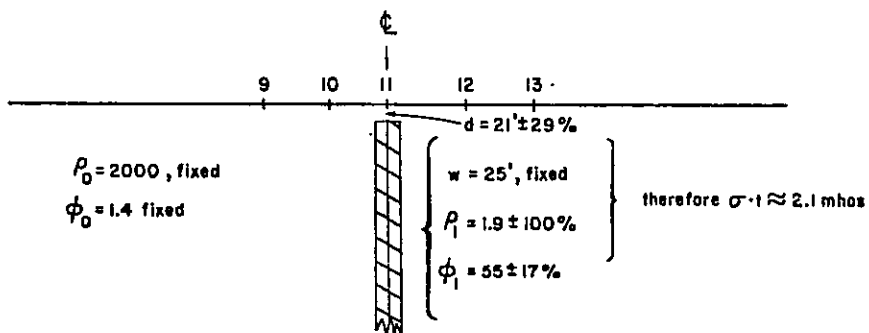
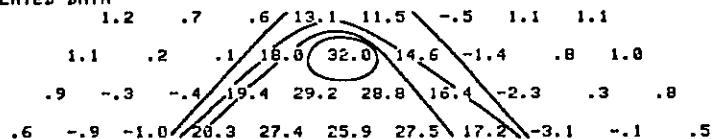


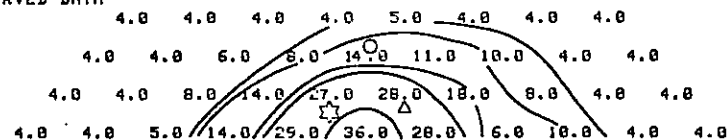
FIG. 35

### LINE-15E graphite source, with some pyrite

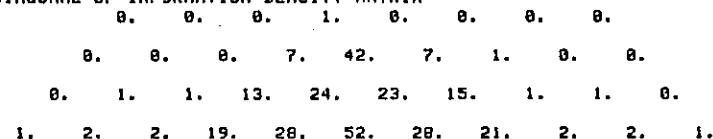
### LINE-15E phase at 0.3 Hz.

FINAL MODEL FOR DATA SET 1

OBSERVED DATA



100\*DIAGONAL OF INFORMATION DENSITY MATRIX



CALCULATED DATA

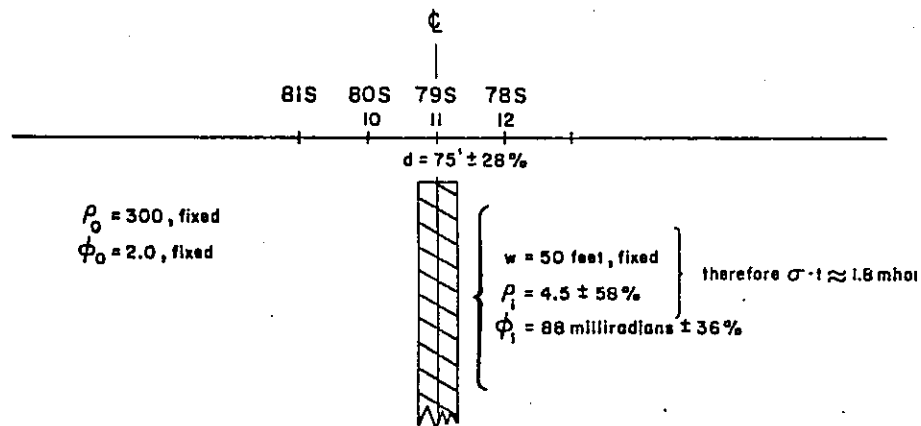
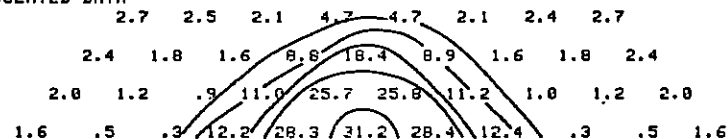


FIG. 36

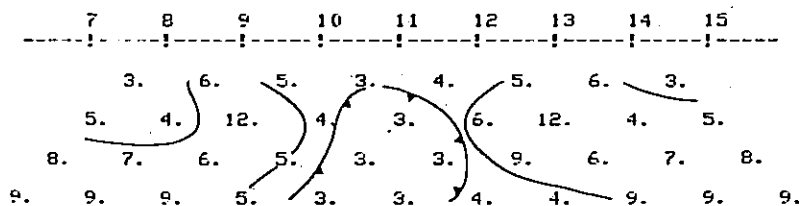
COMPUTER INTERPRETATION  
OF GEOPHYSICAL DATA

INDUCED POLARIZATION AND RESISTIVITY

DEEPLY BURIED SOURCE IN VERY CONDUCTIVE ENVIRONMENT - DETERMINE POSITION,  
DEPTH, AND PHYSICAL  
PROPERTIES OF SOURCE

DIPOLE-DIPOLE APPHY X = 25 METERS 25M.  
-----

RESISTIVITY DATA - OBSERVED ( OHM-METERS )



INDUCED POLARIZATION DATA - OBSERVED ( PFE )

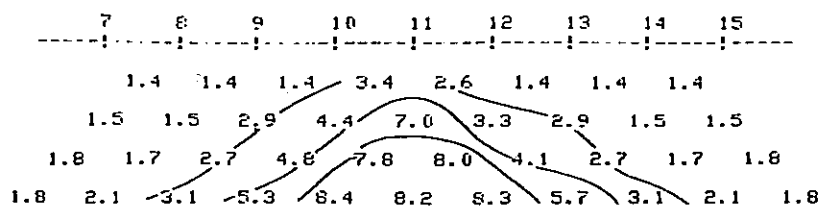
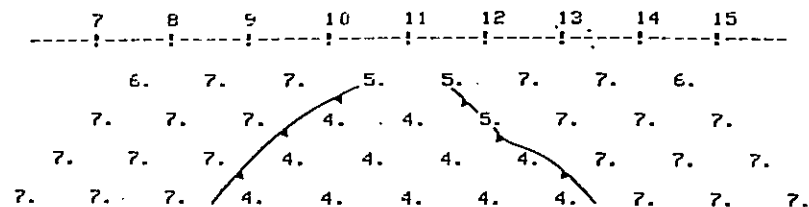
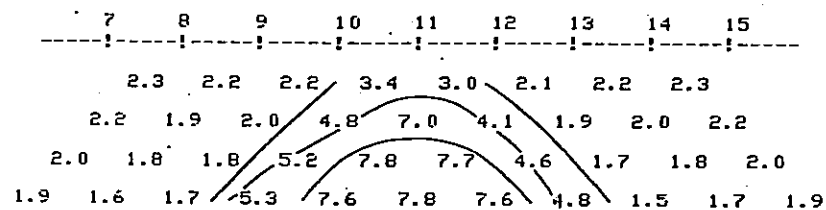


FIG. 37

RESISTIVITY DATA - CALCULATED ( OHM-METERS )



INDUCED POLARIZATION DATA - CALCULATED ( PFE )



PARAMETERS OF THE CALCULATED SOURCE

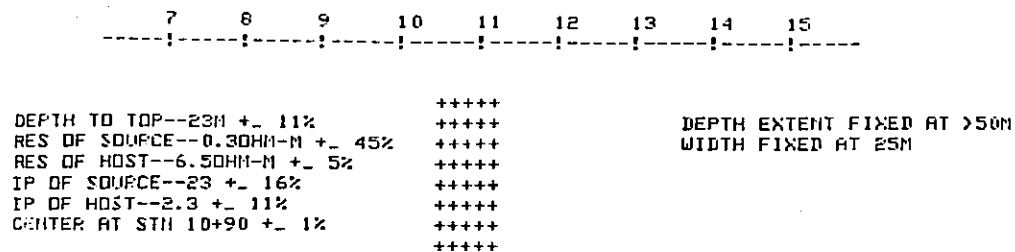


FIG. 38

sources and massive graphite sources in Canada. More recent determinations have not been added to the plot, although additional field work in North America and Australia has shown the same grouping.

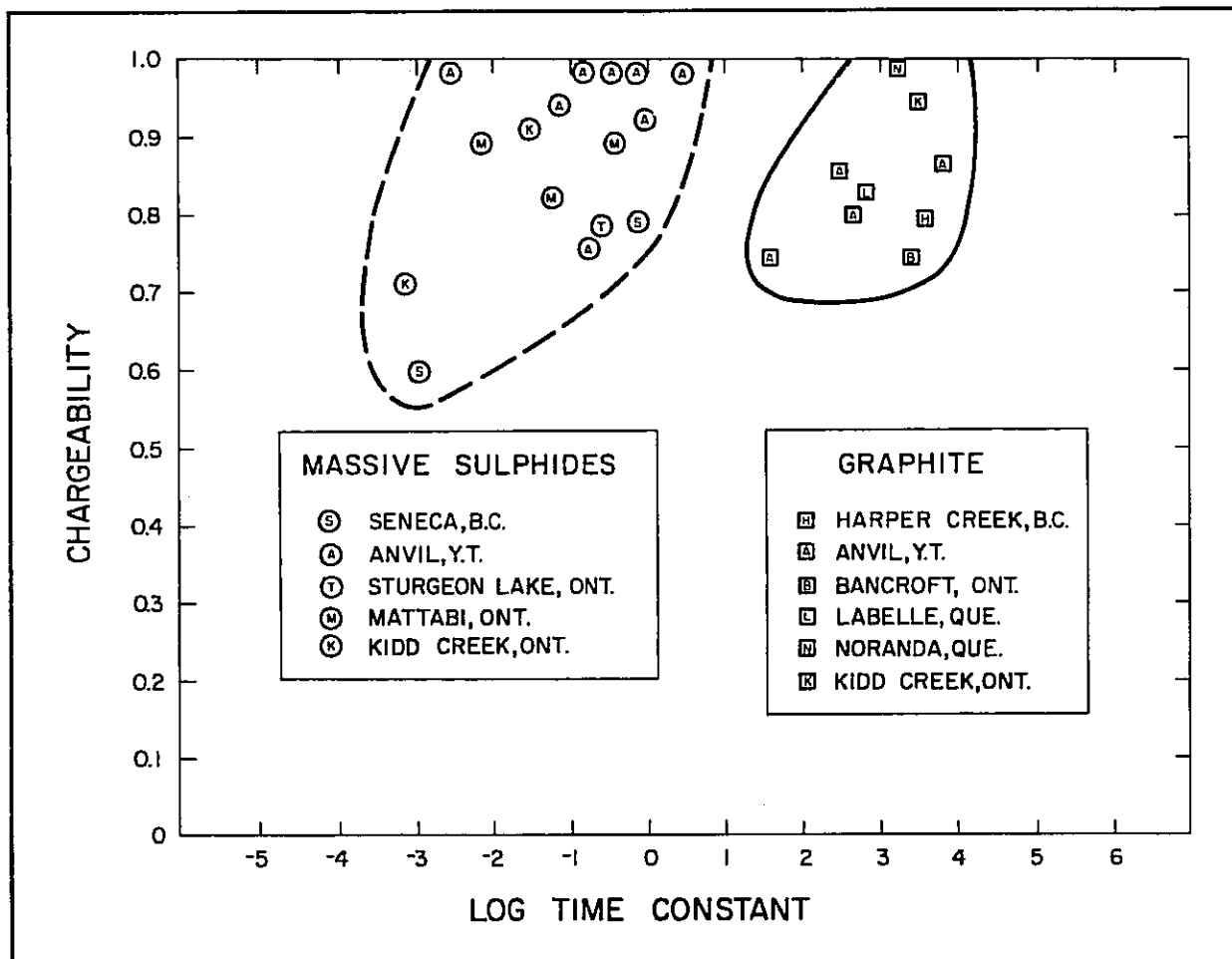


FIG. 39

b) The frequency range of spectral IP measurements with the IPV-3 System extends well up into the range that is normally used for electromagnetic prospecting. Indeed, as we have shown previously, the inductive coupling (EM) effects must be removed from the spectral IP plot before the electrical parameters of the source can be determined. In the near future, it should be possible to use this electromagnetic coupling information to gain further insight about the source.

The induced polarization and resistivity results in Figure 40 are from an area in South Australia. There is a known

pyrrhotite zone at Station 4 to Station 3, and a narrow Pb-Zn-Ag zone at about Station 0. The resistivity data shows a low over the conductive pyrrhotite zone. The raw phase data at 9.0 Hz and 0.11 Hz shows that IP effects are present, but the measurements (particularly at 9.0 Hz) are distorted by inductive coupling (EM) effects. When the inductive effects are removed, there is a definite IP anomaly from the pyrrhotite zone and a weak IP anomaly from the Pb-Zn-Ag zone. (see Figure 41).

The pseudo-sections shown in Figure 42 are a first attempt to use the inductive effects that have been removed from the dipole-dipole IP data. The M3 parameter is the electromagnetic component that has been removed from the measurement with each dipole pair.

The source centered at Station 4 to Station 3 gives a resistivity low, an IP high, and also a very large electromagnetic component. The source centered at Station 0 gives a weak, but definite IP anomaly. This source is not large enough to give a resistivity low using  $X = 25$  meters. However, the electromagnetic coupling between the IP survey wires and the source is different from the resistive coupling between the electrodes and the source. There definitely is an electromagnetic component due to the Pb-Zn-Ag source.

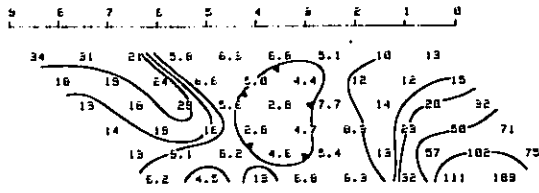
c) With small desk-top computers, the calculation of a single forward IP problem can require five to ten hours. However, recent advances in small computers and new, interaction interfaces between CRT Displays and light-stylus hardware, will soon give the exploration geophysicist almost instantaneous interpretation capability.

For instance, the geometry shown on Figure 43 could be the first guess at the source of an IP anomaly. The geometry would be drawn upon the CRT Display by the geophysicist using an active light-stylus. At the push of a key, the computer would almost instantly display the pseudo-section shown. The computer would then print a "hard copy" of the field data, the calculated data and the geometry and parameters of the source.

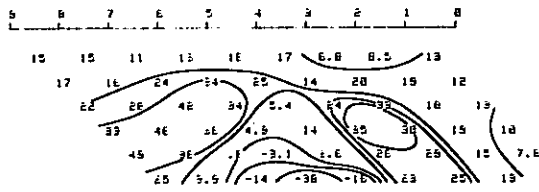
The geophysicist could then use the light-stylus to alter the geometry to that shown in Figure 44. A push of the "operate" key would produce the pseudo-section shown in Figure 44. The pattern is only slightly different, but if it has greater similarity to the field data pattern, the geophysicist knows that he is on the right track.

161 SOUTH AUSTRALIA LINE 8005 R-2511 AUG 1979

Apparent Resistivity(ohm-m)



Phase(mrad) 9 Hz



Phase(mrad) 1 Hz

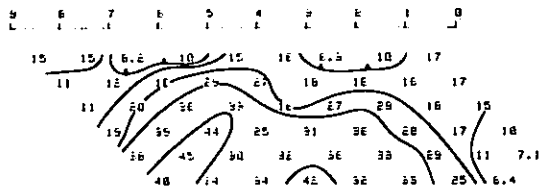
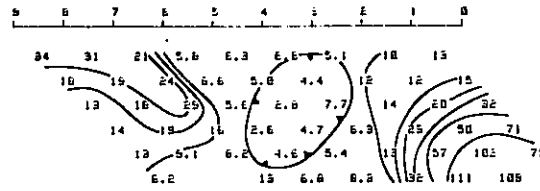


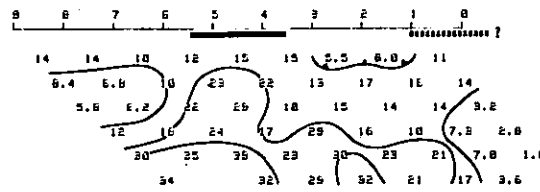
FIG. 40

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Apparent Resistivity(ohm-m)



Decoupled Phase (mrad) 9 Hz



Decoupled Phase (mrad) .111 Hz

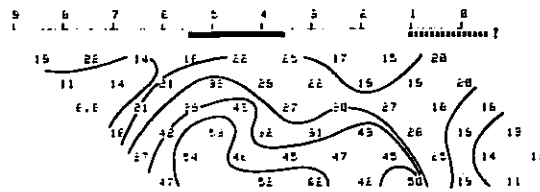
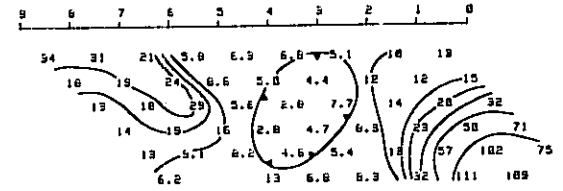


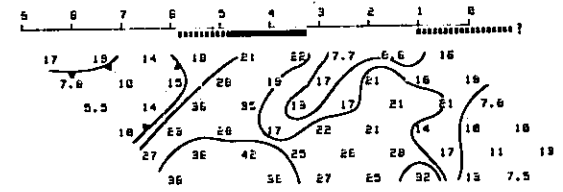
FIG. 41

SOUTH AUSTRALIA LINE 8005 R-2511 AUG 1979

Apparent Resistivity(ohm-m)



Apparent Chargeability(MI in %)



M3 Apparent Electromagnetic Component in %

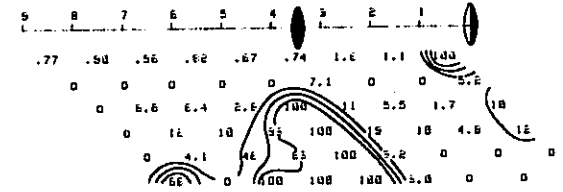


FIG. 42

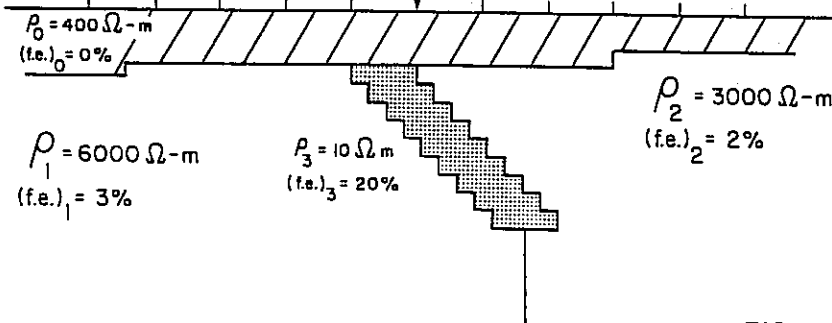
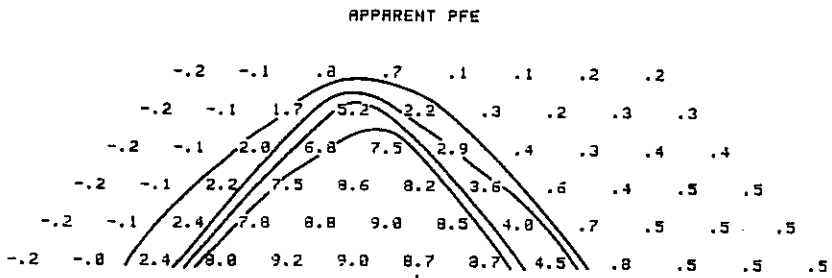
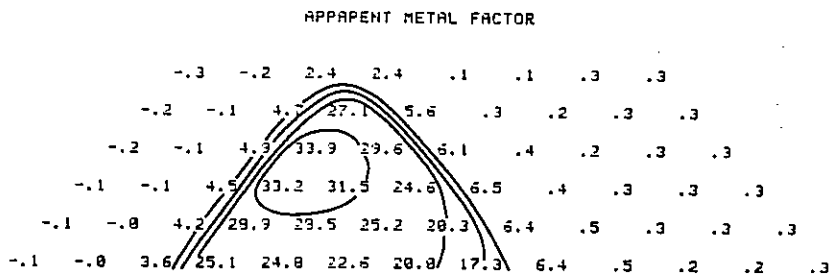
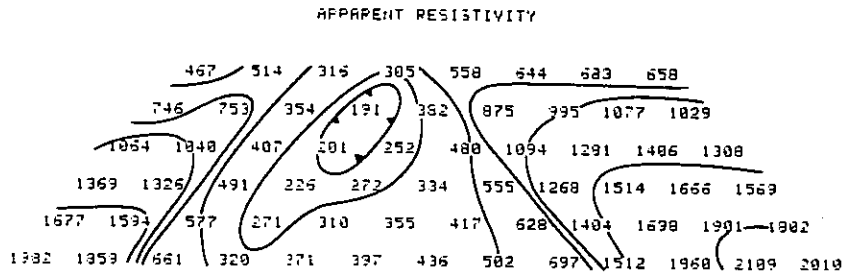


FIG. 43

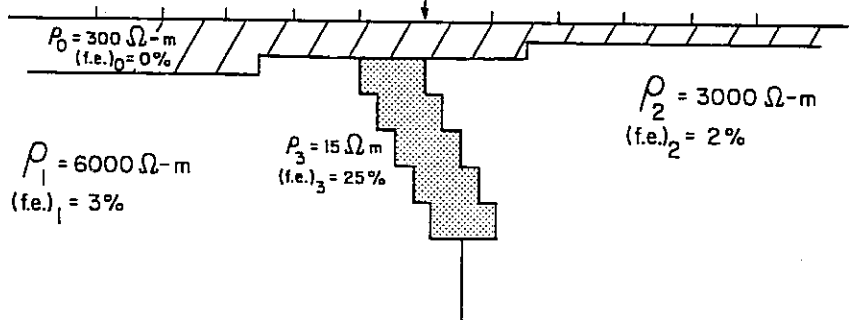
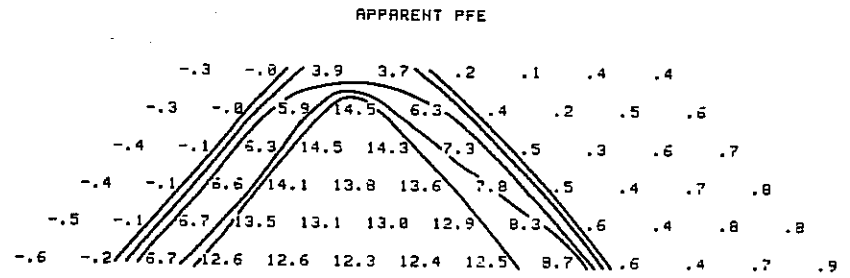
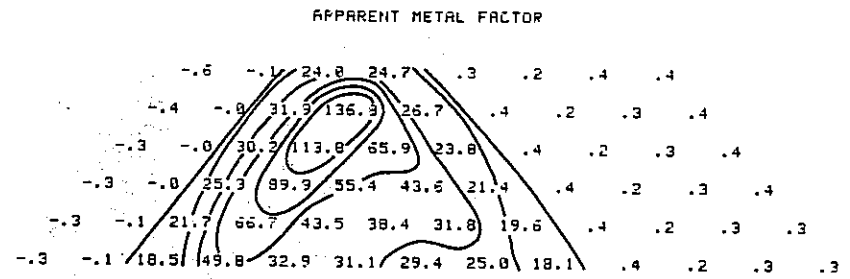
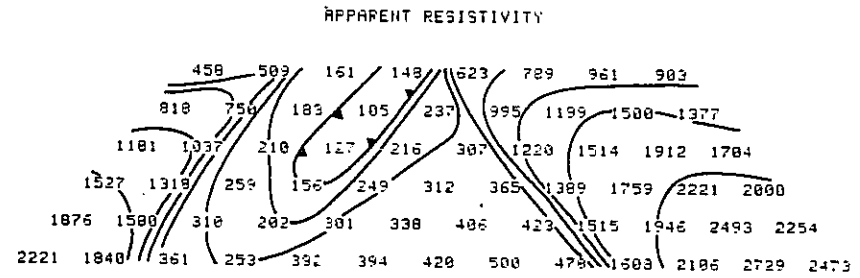


FIG. 44



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